

4 摂動宇宙の観測量

4.1 宇宙背景放射のゆらぎ

輝度関数 (brightness function)

$$\Theta(k, \mu, \tau) = \sum_l (-i)^l (2l+1) P_l(\mu) \Theta_l(k, \tau) \quad (4.1)$$

ボルツマン方程式

$$\Theta' + ik\mu\Theta + ik\mu\Theta + \Psi' = \tau_c^{-1} \left[\Theta_0 - \Theta + ik\mu v_b - \frac{1}{2} P_2(\mu) \Theta_2 \right] \quad (4.2)$$

光学的厚み (optical depth)

$$\kappa(\tau) \equiv \int_{t(\tau)}^{t_0} dt n_e \sigma_{TC} = \int_{\tau}^{\tau_0} d\tau a n_e \sigma_T = \int_{\tau}^{\tau_0} \frac{d\tau}{\tau_c} \quad (4.3)$$

ボルツマン方程式の形式解

$$\Theta(k, \mu, \tau_0) = - \int_0^{\tau_0} d\tau e^{ik\mu(\tau-\tau_0)-\kappa} \left\{ \Psi' + ik\mu\Phi + \kappa' \left[\Theta_0 + ik\mu v_b - \frac{1}{2} P_2(\mu) \Theta_2 \right] \right\} \quad (4.4)$$

$$= \int_0^{\tau_0} d\tau e^{ik\mu(\tau-\tau_0)} \mathcal{S}(k, \tau) \quad (4.5)$$

源関数 (source function)

$$\mathcal{S}(k, \tau) = e^{-\kappa} (\Phi' - \Psi') + g(\tau) \left(\Theta_0 + \frac{1}{4} \Theta_2 + \Phi \right) + (g v_b)' - \frac{3}{4k^2} (g \Theta_2)'' \quad (4.6)$$

視程関数 (visibility function)

$$g(\tau) = -\kappa' e^{-\kappa} \quad (4.7)$$

形式解のルジャンドル表示

$$\Theta_l(k, \tau_0) = \int_0^{\tau_0} \mathcal{S}(\tau, k) j_l[k(\tau_0 - \tau)] \quad (4.8)$$

近似解 ($g(\tau) \rightarrow \delta(\tau - \tau_{\text{dec}})$, $\Theta_2 \simeq 0$)

$$\begin{aligned} \Theta_l(k, \tau_0) &\simeq [\Theta_0(k, \tau_{\text{dec}}) + \Phi(k, \tau_{\text{dec}})] j_l[k(\tau_0 - \tau_{\text{dec}})] \\ &\quad - 3\Theta_1(k, \tau_{\text{dec}}) \left\{ j_{l-1}[k(\tau_0 - \tau_{\text{dec}})] - \frac{l+1}{k(\tau_0 - \tau_{\text{dec}})} j_l[k(\tau_0 - \tau_{\text{dec}})] \right\} \\ &\quad + \int_0^{\tau_0} d\tau e^{-\kappa} [\Phi'(k, \tau) - \Psi'(k, \tau)] j_l[k(\tau_0 - \tau)] \end{aligned} \quad (4.9)$$

温度ゆらぎ

$$\frac{\delta T}{T}(\theta, \phi) = \Theta(\mathbf{x} = \mathbf{0}, \mathbf{n}, \tau_0) = \sum_{l=0}^{\infty} \sum_{m=-l}^l a_{lm} Y_l^m(\theta, \phi) \quad (4.10)$$

$$a_{lm} = \int_0^{\pi} \sin \theta d\theta \int_0^{2\pi} d\phi Y_l^{m*}(\theta, \phi) \frac{\delta T}{T}(\theta, \phi) \quad (4.11)$$

温度ゆらぎのパワースペクトル

$$C_l = \frac{1}{2l+1} \sum_{m=-l}^l \langle |a_{lm}|^2 \rangle = 4\pi \int \frac{k^2 dk}{2\pi^2} P_{\Phi_{m0}}(k) |\hat{\Theta}_l(k)|^2 \quad (4.12)$$

ここで、

$$\Theta_l(\mathbf{k}, \tau_0) = \hat{\Theta}_l(k) \Phi_{m0}(\mathbf{k}) \quad (4.13)$$

ザックス・ウォルフエ効果

$$C_l \simeq \frac{4\pi}{9} \int \frac{k^2 dk}{2\pi^2} P_{\Phi_{m0}}(k) [j_l(k\tau_0)]^2 \quad (4.14)$$

$$= \frac{1}{2\pi} \left(\frac{H_0^2 \Omega_{m0}}{D(\tau_0)} \right)^2 \int \frac{dk}{k^2} P(k) [j_l(k\tau_0)]^2 \quad (4.15)$$

べき則パワースペクトル $P(k) = Ak^n$ に対して

$$C_l \simeq \frac{A\tau_0^{1-n}}{8\sqrt{\pi}} \left(\frac{H_0^2 \Omega_{m0}}{D(\tau_0)} \right)^2 \frac{\Gamma\left(\frac{3-n}{2}\right) \Gamma\left(l + \frac{n-1}{2}\right)}{\Gamma\left(2 - \frac{n}{2}\right) \Gamma\left(l + \frac{5-n}{2}\right)} \quad (4.16)$$

音響振動

$$C_l \sim \int \frac{dk}{k^2} P(k) |\cos[kr_s(\tau_{\text{dec}})]|^2 \{j_l[k(\tau_0 - \tau_{\text{dec}})]\}^2 \quad (4.17)$$

音響ピーク位置

$$l^{(m)} \sim \frac{m\pi(1+z_{\text{dec}})d_A(z_{\text{dec}})}{r_s(\tau_{\text{dec}})} \quad (4.18)$$