

Stability constants of lanthanide(III)-EDTA complex formation and
Gd-break with tetrad effect in their series variation

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Abstract

The logarithmic stability constants of (1:1) lanthanide(III)-EDTA complex formation, $\log K(\text{Ln-EDTA})$, have been reported by polarography, radiotracer-cation exchange, and potentiometry with different ion-selective electrodes. The polarographic data (Schwarzenbach et al., 1954) with uncertainties less than 0.1 for all Ln but Ho are the experimental basis to put forward the Gd-break and convex tetrad effect of $\log K(\text{Ln-EDTA})$. Nevertheless, the data for each $\log K(\text{Ln-EDTA})$ by the different methods are variable in a range of ± 0.5 or more, although their series changes are fairly parallel. Apparently the polarographic data can never be assessed by the other methods, but we found that all the reported data except those by the radiotracer-cation exchange need further corrections when the standardized condition of $T=298.15\text{ K}$ (25°C) and ionic strength $\mu=0.1$ is chosen with assuming that $\log K(\text{Cu-EDTA})=18.80$ (20°C) and $\log K(\text{Pb-EDTA})=18.04$ (20°C). After the corrections and the subsequent normalization of $\log K(\text{Y-EDTA})=17.62$, all the five data sets of $\log K(\text{Ln-EDTA})$ by the different methods converge within $2\sigma(\text{mean})$ errors less than 0.9 % for all Ln. Here are accepted the converged values and their $2\sigma(\text{mean})$ errors as the experimental $\log K(\text{Ln-EDTA})$ ones and uncertainties. The Gd-break and convex tetrad effect certainly exist in $\log K(\text{Ln-EDTA})$ beyond the experimental uncertainty, but the observed convex tetrad effect in the light Ln series appears less obvious than in the heavy Ln one. The slightly asymmetrical tetrad effect is possibly due to the hydration changes of Ln-EDTA and high $\text{Ln}^{3+}(\text{aq})$ series.

Introduction

It is well known that the trivalent lanthanide (Ln) ions in aqueous solution form stable (1:1) complexes with the chelating agent of ethylenediaminetetraacetic acid (EDTA). The stability constants of Ln(III)-EDTA complex formation have been reported by several different methods including polarography, radiotracer-cation exchange, and potentiometry with various ion-selective electrodes (Wheelwright et al., 1953; Schwarzenbach et al., 1954; Betts and Dahlinger, 1959; and Suzuki et al., 1980), which are summarized in Table 1. Gritmon et al. (1977) also reported their results of stability constants of Ln(III)-EDTA complex formation at $\mu=0.5$ M (NaClO₄) and 25°C, but the μ value of ionic strength is greater than $\mu=0.1$ or 0.09 M (KNO₃ or KCl) in the other cases. Therefore, the results by Gritmon et al. (1977) are not considered here. The reported results are plotted against the atomic number of Ln series in Fig. 1, together with their logK(Y-EDTA) data. The break is observed at the position of Gd in the series change of logK(Ln-EDTA). Wheelwright et al. (1953) interpreted that EDTA may act as a hexadentate group for Ln³⁺ from La to Eu but a pentadentate one from Gd to Lu, while Schwarzenbach and Gut (1956) emphasized that the Gd-break may be related to the half-filled effect at Gd³⁺ with the electronic configuration of [Xe](4f⁷).

Fidelis and Siekierski (1966) reported that the separation factor for successive Ln members, $\beta=K_{(Z+1)}/K_{(Z)}$ in reversed phase partition chromatography show regular series variations in the Ln(III)-complex systems with the two phosphonic acid derivatives (HEH ϕ P and HDEHP). Peppard et al. (1969 and 1970) demonstrated that the

regularity of β corresponds to the lanthanide tetrad effect in $\log K$. Fidelis and Siekierski (1966) also mentioned that β 's for $K(\text{Ln-EDTA})$ and other Ln(III) complex systems show such regularities. However, Rowlands (1967) re-examined the regularity using stability constant data for Ln(III) with fourteen complexing ligands including $\text{HEH}\phi\text{P}$, HDEHP , and EDTA , and concluded that the proposal by Fidelis and Siekierski (1966) was based on insufficient data and left doubt as to the general validity, but he accepted that the regularity exists in the two Ln(III) -complex systems with $\text{HEH}\phi\text{P}$ and HDEHP .

The plot of all the reported experimental data of $\log K(\text{Ln-EDTA})$ listed in Table 1 is obviously indicating a large variability of 0.5 or more for each Ln (Fig. 1). The reported experimental errors by Schwarzenbach et al. (1954) are not greater than 0.1, though $\log K(\text{Ho-EDTA})$ is not reported, so that the Gd-break and small tetrad-like variation is plausible as long as the data set is accepted. However, if the large variability of Fig. 1 is regarded as the experimental uncertainty assigned to each $\log K(\text{Ln-EDTA})$ from the different experimental methods, it is impossible to make significant discussion on the Gd-break and convex tetrad effect as argued by Rowlands (1967). The polarographic data by Schwarzenbach et al. (1954) suggesting the Gd break and small tetrad effect, can never be crosschecked by the different methods.

Concerning the Gd-break and tetrad effect of Ln(III) -complexing behavior in natural water systems, REE geochemists are being split into the three: the first is accepting the discussion by Rowlands (1967) and addressing its unimportance (Wood, 1990; Byrne and Sholkovitz, 1996), the second is apparently uninterested in it, and the third is

putting forward the tetrad effect variations of $\log K(\text{LnCO}_3^+)$ and $\log K(\text{Ln}(\text{CO}_3)_2^-)$ and the importance in seawater solutions (Kawabe, 1999a; Ohta and Kawabe, 2000a, b). Hence we are concerned with the proposal by Fidelis and Siekierski (1966) and the conclusion by Rowlands (1967).

Apart from the variability of $\log K$ for each Ln-EDTA in Fig. 1, the parallelism of reported data appears fairly good. Hence it seems important to compare all the data in Table 1 at the standardized experimental condition, because the experimental temperatures are not exactly the same in Table 1. The reference values of $\log K(\text{Cu-EDTA})$ at 20°C are also different between Schwarzenbach et al. (1954) and Wheelwright et al. (1953).

The purpose of this study is to show that the apparent variability of $\log K$ for each Ln in Fig. 1 is too large to be accepted as the uncertainty evaluated from the different experimental methods. Indeed, the most of the experimental data in Table 1 involve uncorrected differences from those at the standardized condition at 25°C and $\mu=0.1$. The acceptable values of $\log K(\text{Ln-EDTA})$ and their experimental uncertainties will be evaluated by making the corrections for the standardized condition and the normalization with $\log K(\text{Y-EDTA})$ values, and then the Gd-break and convex tetrad effect will be assessed solely on the basis of experimental data of $\log K(\text{Ln-EDTA})$.

Results and Discussion

1) Necessary corrections for the standardized experimental condition

When the standardized experimental condition is chosen as $T= 298.15 \text{ K}$ (25°C) and

$\mu=0.1$, all the reported data in Table 1 except those by Betts and Dahlinger (1959) are to be corrected with respect to the temperature difference between 20° and 25°C. Suzuki et al. (1980) reported the experimental results at 25°C, but they used the reference values of $\log K(\text{Cu-EDTA})=18.80$ and $\log K(\text{Pb-EDTA})=18.04$ at 20°C without temperature corrections. Hence their data are not free from such corrections. Wheelwright et al. (1953) accepted $\log K(\text{Cu-EDTA})=18.38$ (20°C) in order to calculate their results from the potentiometric data with glass electrode at 20°C, whereas Schwarzenbach et al. (1954) reported their data at 20°C by using the revised value of $\log K(\text{Cu-EDTA})=18.80$ (20°C).

The temperature correction for $\log K$ is based on the assumption as to approximate constant values of ΔH and ΔS in ΔG for a small temperature range:

$$\Delta G \approx -2.303RT \log K(T) \approx \Delta \bar{H} - T\Delta \bar{S}, \quad (1)$$

where we neglect the subtle difference between ΔG and $-RT \ln K(T)$ at $\mu = 0.1$ as the first-order approximation. For two slightly different temperatures of T_0 and T , it follows that

$$\log K(T) \approx (T_0/T) \cdot \log K(T_0) + (T - T_0) \cdot \Delta \bar{S} / (2.303RT). \quad (2)$$

Note that the ΔS value is necessary for the temperature correction of (2), and that the ΔS value is usually calculated from experimental values of ΔH and $\Delta G = -2.303RT \log K$. Hence, even when the experimental value of ΔH is known, the ΔS value cannot be calculated until $\log K(T)$ value is known. Nevertheless, it is only necessary to use an approximate value of $\Delta \bar{S}$ in (2), because the second term of (2) is a small correction insensitive to the choice of $\log K(T)$ value. The ΔS values for Ln-EDTA formation

have been reported by Mackey et al. (1962), in which $\log K(\text{Ln-EDTA})$ at 25°C by Betts and Dahlinger (1959) and $\log K(\text{Y-EDTA})$ by Wheelwright et al. (1953) were accepted, together with their experimental ΔH values. The ΔH values for Cu(II)- and Pb(II)-EDTA formation at 25°C and $\mu=0.1$ (KNO₃) are -36.8 (kJ/mol) and -54.7 (kJ/mol), respectively (Handbook of Chemistry, 5th ed., 2004). The approximate values of $\Delta \bar{S}$ in (2) for Cu(II)- and Pb(II)-EDTA formation are +236 (J/mol/K) and +162 (J/mol/K), respectively.

Here is accept $\log K(\text{Cu-EDTA})=18.80$ (20°C) as in Schwarzenbach et al. (1954), and then the data by Wheelwright et al. (1953) are corrected for the difference between $\log K(\text{Cu-EDTA})=18.38$ and 18.80 at 20°C. The corrected data of Wheelwright et al. (1953) are further adjusted to the condition of 25°C using eq. (2), along with the reported data by Schwarzenbach et al. (1954). The corrected values to 25°C are listed in Table 2.

In the case of the two data sets with Pb- and Cu electrodes of Suzuki et al. (1980), the temperature corrections were applied to their reference values of $\log K(\text{Cu-EDTA})=18.80$ and $\log K(\text{Pb-EDTA})=18.04$ at 20°C by using eq. (2). They are reduced to $\log K(\text{Cu-EDTA})=18.69$ and $\log K(\text{Pb-EDTA})=17.88$ at 25°C. The respective differences between the two temperatures are the corrections to the two data sets with Pb- and Cu electrodes. The resultant values are listed in Table 2. All the values of Table 2 are for the standardized experimental condition of $T= 298.15$ K (25°C) and $\mu=0.1$ with assuming that $\log K(\text{Cu-EDTA})=18.80$ and $\log K(\text{Pb-EDTA})=18.04$ at 20°C. They are plotted against the atomic number of Ln in

Fig. 2-A.

It is apparent in Table 2 and Fig. 2-A that the polarographic data by Schwarzenbach et al. (1954) are almost in accordance with the data with Cu-electrode by Suzuki et al. (1980), while the potentiometric ones by Wheelwright et al. (1953) are close to the data with Pb-electrode by Suzuki et al. (1980). The reported data by the radiotracer-cation exchange technique (Betts and Dahlinger, 1959) at 25°C and $\mu=0.1(\text{KCl})$ in Table 1, which are free from the assumptions as to $\log K(\text{Cu-EDTA})$ and $\log K(\text{Pb-EDTA})$, are closer to the data sets by potentiometry with Pb-electrode (Suzuki et al., 1980) and glass electrode (Wheelwright et al., 1953).

2) Acceptable values of $\log K(\text{Ln-EDTA})$ and their uncertainties

The corrected data in Table 2 and those by Betts and Dahlinger (1959) in Table 1 reveal that the experimental data for $\log K(\text{Ln-EDTA})$ by five different methods are fairly consistent, but there still remain some discrepancies possibly due to systematic experimental differences (Fig. 2-A). In order to find the acceptable data set and their uncertainties we have normalized the corrected data in Table 2 to $\log K(\text{Y-EDTA})=17.62$. The normalizing value of $\log K(\text{Y-EDTA})$ is that by Pb-electrode method of Suzuki et al. (1980), which is closer to the corrected potentiometric value of 17.37 ± 0.15 by Wheelwright et al. (1953). In addition, although Betts and Dahlinger (1959) did not report the $\log K(\text{Y-EDTA})$ value, their data for $\log K(\text{Ln-EDTA})$ are closer to those by Pb-electrode method of Suzuki et al. (1980) and by the glass electrode of Wheelwright et al. (1953). If we normalize the data to

$\log K(\text{Y-EDTA})=18.09\pm 0.04$ of the corrected polarographic value by Schwarzenbach et al. (1954) or 18.12 by Cu-electrode method of Suzuki et al. (1980), the systematic discrepancy of about 0.5 is left between those by Betts and Dahlinger (1959) and the others. This is the reason why we prefer the normalization by $\log K(\text{Y-EDTA})=17.62$ to those by other values.

The normalized values plotted in Fig. 2-B are listed in Table 3. The comparison of Fig. 1 with Fig. 2-B is illustrating the importance not only of the corrections for the standardized experimental condition but also of the subsequent normalization by $\log K(\text{Y-EDTA})=17.62$. The mean value for each $\log K(\text{Ln-EDTA})$ and the two-sigma error for the mean have been calculated from the four data sets by potentiometry and polarography and from the five data sets including those by the radiotracer-cation exchange (Betts and Dahlinger, 1959). The means and the two- σ (mean) errors for the experimental data with and without those by the radiotracer-cation exchange are listed in the last two columns of Table 3. There are no significant differences between the means and the two- σ (mean) errors for $n=4$ and 5, suggesting that the reported values by the radiotracer-cation exchange (Betts and Dahlinger, 1959) are consistent with those by potentiometry and polarography corrected to the standardized condition and then normalized to $\log K(\text{Y-EDTA})=17.62$.

The two- σ (mean) errors for the respective means of $\log K(\text{Ln-EDTA})$ for the five data sets are less than 0.9 % for all Ln. Although two- σ (mean) errors for the $\log K(\text{Ln-EDTA})$ values with $\text{Ln}=\text{La}$, Ho , and Er are equal to 0.1 or slightly beyond 0.1, those errors for the others are less than 0.1. We have accepted the means and

two- σ (mean) errors in the last column of Table 3 as the most probable values of $\log K(\text{Ln-EDTA})$ and their experimental uncertainties. They are based on the results by five different methods for the standardized condition that $T=298.15\text{ K}$ (25°C), $\mu=0.1$, $\log K(\text{Cu-EDTA})=18.80(20^\circ\text{C})$, and $\log K(\text{Pb-EDTA})=18.04$ (20°C), together with the normalization condition of $\log K(\text{Y-EDTA})=17.62$ (25°C).

3) Gd-break and tetrad effect existing in the series change of $\log K(\text{Ln-EDTA})$

The $\log K(\text{Ln-EDTA})$ values accepted by us give a smooth curve for each tetrad of Ln similar to Peppard et al. (1969), which is shown in Fig. 3. The smooth tetrad curves have been determined by the least-squares method assuming second-order polynomials of the atomic number of Ln for the respective tetrads, in which the following constraints were imposed: the cusps between the first and second tetrad curves and between the third and fourth ones are located at the mid-points between Nd and Pm and between Ho and Er, respectively. The least-squares method with such constraints has applied to the problem of lanthanide tetrad effect by Minami and Masuda (1997), in which the details of the fitting method are described. The second-order polynomials was used to describe the smooth tetrad curves from experimental data by Masuda et al. (1994). The characteristic variations exhibit the Gd-break and the tetrad effect beyond the experimental errors of two- σ (mean). For most of the accepted values of $\log \log K(\text{Ln-EDTA})$, their two- σ (mean) errors are less than 0.1.

The Gd-break and convex tetrad effect of $\log K(\text{Ln-EDTA})$ have confirmed solely on the basis of experimental data. It seems interesting that the tetrad effect variation of

Fig. 3 is not quite symmetrical between the light and heavy Ln series: The convex tetrad effect in the light Ln series appears less obvious than in the heavy Ln one. The slightly asymmetrical tetrad effect as the series change property of $\log K(\text{Ln-EDTA})$ is possibly related to the hydration changes of light $\text{Ln}^{3+}(\text{aq})$ and $\text{Ln-EDTA}^-(\text{aq})$ series (Rizkalla and Choppin, 1991; Kawabe, 1999a and b). This invites further studies of the Ln-EDTA formation and the other (1:1) Ln-chelate complex formation from the viewpoints of thermodynamics, spectroscopy, and the improved refined spin-pairing energy theory for thermodynamic data (Kawabe, 1992 and 1999a, b; Kawabe and Masuda, 2001). They will be presented elsewhere, together with further implications to REE geochemistry.

Conclusions

The polarographic data of $\log K(\text{Ln-EDTA})$ with uncertainties less than 0.1 for all Ln but Ho (Schwarzenbach et al., 1954) have been the experimental basis to put forward the Gd-break and convex tetrad effect of $\log K(\text{Ln-EDTA})$. Nevertheless, $\log K(\text{Ln-EDTA})$ data by polarography, radiotracer-cation exchange, and potentiometry with various ion-selective electrodes, exhibit a large variability of ± 0.5 or more for each Ln, although their series changes appear to be fairly parallel. Apparently it is impossible to assess the polarographic data from those by the other methods, but we found that the reported data, except those by the radiotracer-cation exchange, need further corrections for the standardized condition of 25°C and $\mu=0.1$ with assuming $\log K(\text{Cu-EDTA})=18.80$ (20°C) and $\log K(\text{Pb-EDTA})=18.04$ (20°C). The large

variability of ± 0.5 or more is too large to be accepted as the experimental uncertainty for $\log K(\text{Ln-EDTA})$. The conclusions in this study are as follows:

1) All the four data sets of $\log K(\text{Ln-EDTA})$ by polarography and potentiometry, when corrected properly to the standardized condition and normalized to $\log K(\text{Y-EDTA})=17.62$, are combined with the data by the radiotracer-cation exchange which are not reporting the $\log K(\text{Y-EDTA})$ value. The five data sets converge within $2\sigma(\text{mean})$ errors less than 0.9 % for all Ln.

2) The means of five experimental values and $2\sigma(\text{mean})$ errors for $\log K(\text{Ln-EDTA})$ have been accepted. The polarographic data by Schwarzenbach et al. (1954) are justified by the potentiometric data with different ion-selective electrodes and those by the radiotracer-cation exchange.

3) The Gd-break and convex tetrad effect certainly exist in the series change of $\log K(\text{Ln-EDTA})$ beyond the experimental uncertainty. The convex tetrad effect in the light Ln series appears less obvious than in the heavy Ln one. The faintly asymmetrical tetrad effect is possibly related to the hydration changes of light $\text{Ln}^{3+}(\text{aq})$ and Ln-EDTA series.

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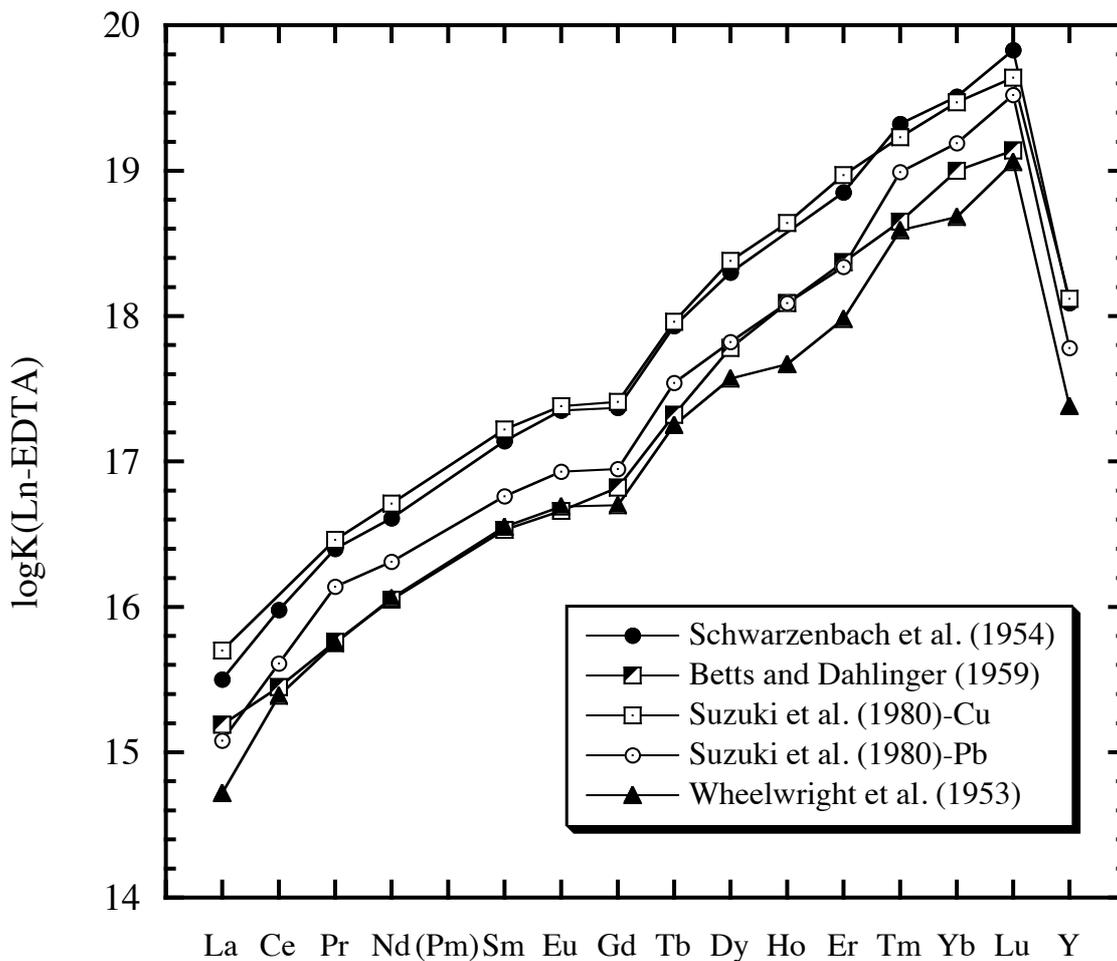


Fig.1. Reported data of $\log K$ for Ln-EDTA complex formation (Table 1) are plotted against the atomic number of Ln. The apparent variability of $\log K$ for each Ln exceeds ± 0.5 , though the series changes are fairly parallel among the five data sets. The data of $\log(\text{Y-EDTA})$ are plotted on the right of Lu.

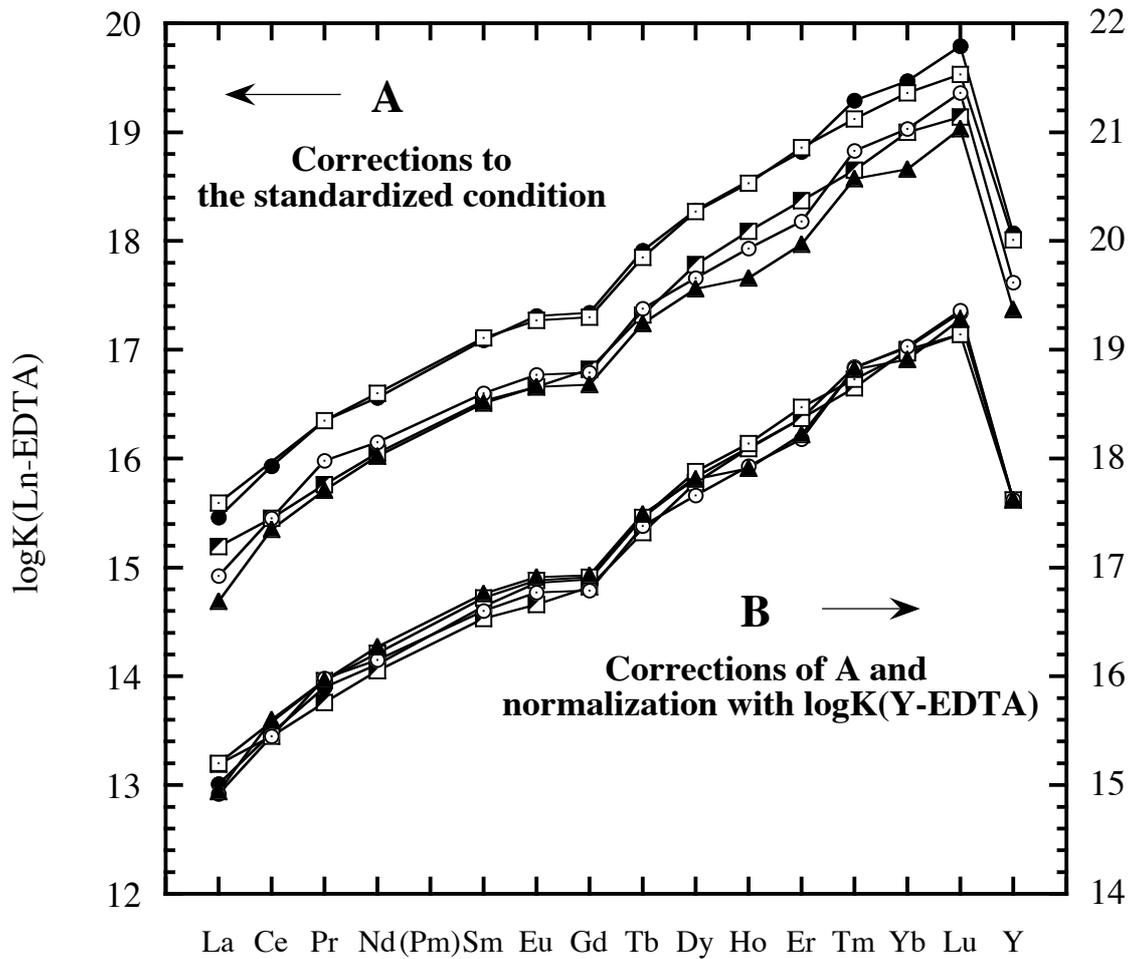


Fig. 2. A: Corrected data of $\log K(\text{Ln-EDTA})$ and $\log K(\text{Y-EDTA})$ by polarography (Schwarzenbach et al., 1954) and potentiometry (Suzuki et al., 1980; Wheelwright et al., 1953) to the standardized condition listed in Table 2. B: The corrected data of A are further normalized to $\log K(\text{Y-EDTA})=17.62$ (Table 3). Symbols are the same as in Fig. 1. The half-filled squares (Betts and Dahlinger, 1959) are the reported ones free from the corrections and normalization here.

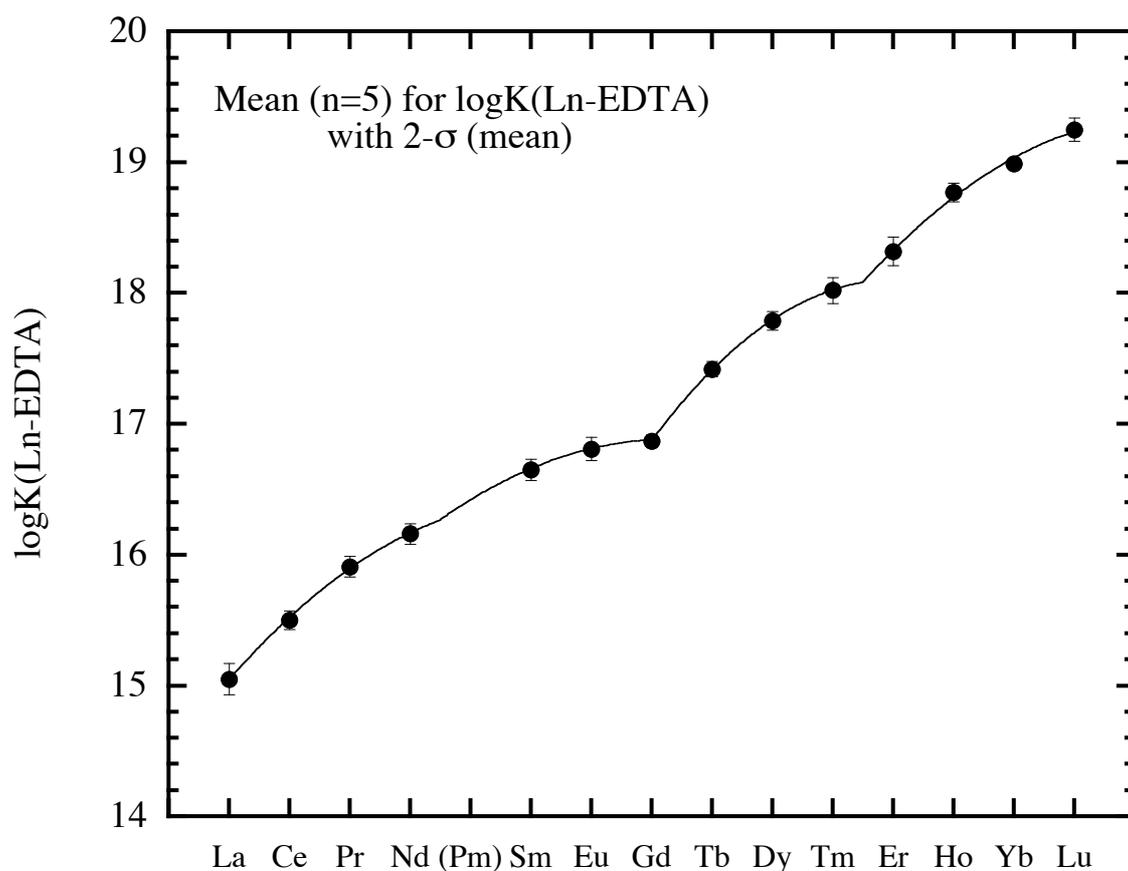


Fig. 3. The series change of the accepted means with $2\text{-}\sigma(\text{mean})$ errors for $\log K(\text{Ln-EDTA})$ from the five data sets (the last column of Table 3). Smooth curves for the respective tetrads are the least-squares fittings with second-order polynomials of the atomic number of Ln with the following constraint: the cusps between the first and second tetrad curves and between the third and fourth ones are located at the mid-points between Nd and Pm and between Ho and Er, respectively (see text).