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#### **Key Points:**

- Envelope widths of natural and active seismic sources were systematically analyzed at five volcanoes
- Comparison with envelope simulations showed the existence of strong heterogeneities in a thin surface layer to 1 km depth
- Residuals between observed envelope widths and those of our estimated 1D models correlated with tomographic images

Supporting Information:

Supporting Information S1

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## Scattering and Attenuation Characteristics at Volcanoes Inferred From Envelope Widths of Natural and Active Seismic Sources

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**Abstract** Seismic scattering and attenuation at volcanoes, thought to be strongest on the Earth, can be used to map volcanic feeding systems. We systematically analyzed high-frequency (5-10 Hz) seismograms of volcano-tectonic earthquakes at Galeras volcano (Colombia) and active sources at Kirishima, Unzen, Bandai, and Iwate volcanoes (Japan) to investigate their scattering and attenuation characteristics. The envelope widths estimated from these seismograms were compared with those calculated by Monte Carlo envelope simulations for 1D layered models parameterized by the scattering mean free path and the quality factor of medium attenuation for S waves. Our results indicated a surficial, highly heterogeneous, and attenuative layer up to around 1 km thickness at all studied volcanoes. The strongest heterogeneities at volcanoes thus exist in a thin surface layer, likely comprising unconsolidated and/or highly fractured materials. Using the space-weighting function for diffusive wavefields, we mapped the residuals between observed envelope widths and those calculated with our estimated 1D models at Kirishima, Unzen, Bandai, and Iwate. These maps showed spatial distributions of the envelope-width residuals were unique to each volcano and correlated with P wave velocity tomographic images. Areas of positive residuals correspond to low-velocity anomalies and thus to heterogeneous, strongly scattering rocks, whereas areas of negative residuals correspond to high-velocity anomalies and thus less heterogeneous volcanic or basement rocks. Our results demonstrate that envelope widths can improve the characterization of scattering and attenuation structures beneath volcanoes.

## 1. Introduction

Estimation of volcano structures is fundamental to our understanding of volcanic processes related to magma ascent, gas migration, and magmatic and phreatic eruptions. Seismological studies have played an essential role in imaging volcano structures. A wide variety of *P* and *S* wave velocity ( $V_p$  and  $V_s$ , respectively) structures have been estimated at various volcanoes by tomographic methods to invert arrival times of body waves from natural and active seismic sources (e.g., Chouet & Matoza, 2013; Koulakov & Shapiro, 2015; Lees, 2007). Furthermore, temporal changes in the  $V_p/V_s$  ratios have been estimated by 4D seismic tomography (e.g., Koulakov et al., 2013; Londoño & Kumagai, 2018; Patanè et al., 2006). These studies have shown that volcanoes display seismologically detectable structural changes especially in association with magmatic intrusions.

Seismic heterogeneous characteristics beneath volcanoes are represented by the scattering mean free path  $(l_0)$  and the quality factor of medium attenuation  $(Q_i)$  for *S* waves. By fitting the seismograms excited by active sources (dynamite blasts) to those calculated with the analytical solutions of the diffusion equation or the radiative transfer theory (Paasschens, 1997) in a uniform medium or half-space,  $l_0$  of 0.1–1 km and  $Q_i$  of 10–100 in frequency bands of 2–20 Hz have been estimated at various volcanoes (e.g., Del Pezzo et al., 2006; Friedrich & Wegler, 2005; Parsiegla & Wegler, 2008; Wegler, 2003; Wegler & Lühr, 2001; Yamamoto & Sato, 2010). These results indicate that volcanoes expose the strongest structural heterogeneities on the Earth and their seismic wavefields are dominated by multiple scattering (Sato et al., 2012). The space-weighting function (SWF) or sensitivity kernels of  $l_0$  and  $Q_i$  for multiple scattering wavefields have been used to determine the spatial distributions of scattering and attenuation at volcanoes (e.g., Akande et al., 2013; Prudencio, Ibanez, et al., 2013; Prudencio et al., 2015; Prudencio, Aoki et al., 2017; Prudencio, Taira et al., 2017; Sketsiou et al., 2020).





**Figure 1.** Topography of Galeras (200 m contour intervals), source locations of VT earthquakes (red circles), and seismic stations (black triangles). The location of Galeras is indicated by the red triangle in the regional map (bottom right).

Seismogram envelope broadening and peak delays, which have been reported for both regional tectonic earthquakes (e.g., Obara & Sato, 1995; Saito et al., 2005; Takahashi et al., 2007; Tripathi et al., 2010) and volcano-tectonic (VT) earthquakes (De Siena et al., 2016), have been explained by regional-scale heterogeneous structures. Kumagai et al. (2018) estimated envelope widths (p, the ratio of the cumulative to the peak amplitude for S waves in high-frequency [5–10 Hz] seismograms) of VT earthquakes at Nevado del Ruiz (Colombia) and Taal (Philippines) and compared them to those calculated by Monte Carlo envelope simulations following the method of Yoshimoto (2000) for layered  $l_0$  and  $Q_i$  models. They showed that a highly heterogeneous and attenuative layer at these volcanoes exists only to 1 km depth, below which  $l_0$  and  $Q_i$  are similar to those in the normal crust. Such a strongly heterogenous surface layer was also reported by Wegler (2004) at Vesuvius volcano (Italy) based on comparison of seismograms produced by active sources with the analytical solution of the diffusion equation in a surface layer over a uniform half-space. These studies suggest that  $l_0$  and  $Q_i$  structures are strongly depth dependent at volcanoes, far from the uniform  $l_0$  and  $Q_i$  distributions assumed in most previous studies using the analytical solutions. Therefore, further studies are required to estimate 1D  $l_0$  and  $Q_i$  structures at other volcanoes and to develop a method for investigating their

#### Table 1

One-Dimensional P Wave Velocity Structure Used in Our Hypocenter Determinations at Galeras

Depth (km)	P wave velocity (m/s)
-5.5	3,500
-2.2	3,700
-0.2	4,000
3.8	6,000
21.8	6,800

3D distributions, which will improve our understanding of magmatic and hydrothermal systems.

Recently, the amplitude source location (ASL) method (Battaglia & Aki, 2003; Kumagai et al., 2010; Yamasato, 1997) and the seismic amplitude ratio method (Taisne et al., 2011) have been widely used to determine the source locations of long period (LP) and very long period (VLP) events and tremor with emergent onsets (e.g., Caudron et al., 2018; Ichihara & Matsumoto, 2017; Ichimura et al., 2018; Kumagai et al., 2013; Kurokawa et al., 2016; Maeda et al., 2019; Mori & Kumagai, 2019; Ogiso & Yomogida, 2015;





Figure 2. Topography (100 m contour intervals) and temporary seismic stations (black triangles) at (a) Kirishima, (b) Unzen, (c) Bandai, and (d) Iwate. Red circles indicate the epicenters of active sources produced by dynamite blasts, and the red labels (S1, S2, etc.) indicate the shot number. The locations of Kirishima, Unzen, Bandai, and Iwate are shown by red triangles in the inset regional maps.

Pérez-Guillén et al., 2019; Tan et al., 2019; Walter et al., 2017). As the assumption of isotropic radiation of *S* waves used in these methods holds for strong structural heterogeneities at volcanoes (Morioka et al., 2017), accommodating depth-dependent scattering and attenuation structures contributes to improving the accuracy of ASL estimates and hence seismic monitoring of active volcanism (Kumagai et al., 2019).

In this paper, we applied the method of Kumagai et al. (2018) based on envelope widths to VT earthquakes at Galeras volcano (Colombia) and active sources at Kirishima, Unzen, Bandai, and Iwate volcanoes (Japan) to estimate their 1D  $l_0$  and  $Q_i$  structures. Our results confirm that a highly heterogenous and attenuative surface layer commonly exists to a depth of around 1 km beneath these volcanoes. We used the 2D SWF of Del Pezzo et al. (2016) to map the distributions of the residuals between observed envelope widths and those calculated with 1D models ( $\Delta p$ ) at Kirishima, Unzen, Bandai, and Iwate and found that the  $\Delta p$  distributions correlate with tomographic images. We discuss the similarity and diversity in our estimated 1D structures and  $\Delta p$  distributions and their relations to volcanic structures.

## 2. Data and Methods

#### 2.1. VT Earthquakes at Galeras

Galeras in Colombia (Figure 1) is an active andesitic volcano producing repetitive dome-forming and Vulcanian eruptions (e.g., Bain et al., 2019; Calvache & Williams, 1997) accompanied by VT earthquakes,



Table 2         Parameter Values Used in Our Monte Carlo Simulations					
Parameter	Value				
$r_0$ $\Delta r$ $\Delta z$ $\Delta t$ f	0.1, 0.5, 1.0,, 13.5, 14.0 km 50 m 100 m 0.02 s 7.5 Hz				

LP events, and tremor (e.g., Gil Cruz & Chouet, 1997; Stix et al., 1997; Torres et al., 1996). For our analysis, we used VT earthquakes observed by broadband and short-period seismometers at Galeras between April 2015 and February 2019 (Figure 1). No eruptions occurred during this period, but a swarm of VT earthquakes occurred during May–July 2018. Following Kumagai et al. (2018), we determined hypocenters and origin times by the method of Benz et al. (1996) based on the finite-difference method of Podvin and Lecomte (1991) to calculate traveltimes of first-arrival waves. Table 1 lists the velocity model in our hypocenter determinations.

#### 2.2. Active Sources at Kirishima, Unzen, Bandai, and Iwate

Kirishima (Kyushu, southwestern Japan) is a volcanic complex, where sub-Plinian eruptions occurred at Shimoe-dake during 26–27 January 2011 (Ichihara, 2016; Kozono et al., 2014; Nakada et al., 2013; Shimbori et al., 2013). Unzen (Kyushu) produced dome-forming eruptions accompanied by deadly pyroclastic flows at Fugen-dake in June 1991 (Nakada et al., 1999; Sato et al., 1992; Yamasato et al., 1993). Bandai is a stratovolcano in northeastern Japan, at which a phreatic eruption triggered a large sector collapse in 1888 (Nakamura, 1978; Sekiya & Kikuchi, 1889). Iwate (northeastern Japan) became active in 1998 with VT earth-quakes, LP and VLP events, tremor, and crustal deformation (Miura et al., 2000; Nishimura et al., 2000; Tanaka, Hamaguchi, Yamawaki et al., 2002), although no eruptions followed.

For our analyses at these volcanoes, we used active sources produced by dynamite blasts at Kirishima in 1994 (Kagiyama et al., 1995; Matsushima et al., 1995; Tomatsu et al., 1997, 2001; Tsutsui et al., 1996), Unzen in 1995 (Matsushima et al., 1997; Nishi, 2002), Bandai in 1997 (Yamawaki et al., 2004), and Iwate in 2000 (Tanaka, Hamaguchi, Yamawaki et al., 2002; Tanaka, Hamaguchi, Nishimura, et al., 2002) (Figure 2). These seismic explorations were conducted by the Japanese national project to better understand volcano structures (Kagiyama et al., 1995). More than 100 short-period seismometers with local data loggers were temporarily deployed to observe earthquakes excited by chemical explosions at each volcano (Figure 2), and their detailed  $V_p$  structures were estimated (Nishi, 2002; Tanaka, Hamaguchi, Nishimura, et al., 2002; Tomatsu et al., 1997, 2001; Tsutsui et al., 1996; Yamawaki et al., 2004). In our analyses of the envelope widths of active sources at these volcanoes, we excluded some explosions located far from the volcanic centers, and some stations very close to sea level because they may not provide accurate structural information.

#### 2.3. Envelope Widths

We used the method of Kumagai et al. (2018) to estimate envelope widths of VT earthquakes at Galeras and active sources at Kirishima, Unzen, Bandai, and Iwate. We measured *P* wave arrival times of VT earthquakes at Galeras and used those reported by Kagiyama et al. (1995) at Kirshima, Matsushima et al. (1995) at Unzen, Yamawaki et al. (2004) at Bandai, and Tanaka et al. (2002) at Iwate. We obtained envelope waveforms from vertical band-passed (5–10 Hz) seismograms. We measured the peak amplitude (*A*) in the envelope waveform after the *S* wave arrival time. Noise portions before the *P* wave arrival time in the time-integrated envelope waveform were used to estimate the linear trend. We estimated the offset of the integrated waveform from the trend line at the end of coda wave arrivals ( $I_e$ ) and that at the *S* wave arrival time ( $I_p$ ) to obtain the cumulative amplitude  $I = I_e - I_p$  (see Kumagai et al., 2018, for more details). The envelope width was estimated as p = I/A.

## 2.4. Monte Carlo Envelope Simulations

Table 3								
One-Dimensional	S	Wave	Velocity	Structure	Used	in	Our	Monte
Carlo Simulations								

Depth (km)	S wave velocity (m/s)
0.0	1,750
1.0	2,250
4.0	3,000
15.0	3,500

We simulated envelope waveforms in depth-dependent  $l_0$  and  $Q_i$  structures using the method of Yoshimoto (2000). Following Kumagai et al. (2018), we assumed scalar *S* waves, a linear increase of  $V_s$  with depth, and constant  $l_0$  and  $Q_i$  values in each layer. Each particle isotropically radiated from a source moves over distance  $V_s\Delta t$  without scattering for a time step  $\Delta t$ , and its energy attenuates by  $\exp(-2\pi f\Delta t/Q_i)$ , where *f* is frequency. The particle changes the propagation direction with the probability given by  $V_s\Delta t/l_0$ . We used 3D domains with horizontal dimensions of  $32 \times 32$  km and vertical dimensions of 15 km for Galeras and 5 km for



Table 4

Values of Scattering Mean Free Path ( $l_0$ ), the Quality Factor of Medium Attenuation for S Waves ( $Q_i$ ), and Layer Depths Used for Trial Models in Our Monte Carlo Simulations for Galeras

Layer number	Layer boundary depth (km)	<i>l</i> <sub>0</sub> (km)	$Q_i$
1	1.0	0.5, 1, 10	50, 100, 200
2	5.0	1, 10, 100, 1,000	100, 200, 300
3	15.0	10, 100, 1,000	100, 200, 300

the other volcanoes, in which absorbing boundary conditions were used except for a free flat surface. We assumed three- and two-layered structures at Galeras and the other volcanoes, respectively. Assuming radial symmetry, we set receivers at equal distances from the epicenter, and we accumulated energy particles that reached the receiver volume defined by  $r_0 \pm \Delta r$  and  $\Delta z$  in the radial and vertical directions, respectively, where  $r_0$  is the radius from the epicenter. We used source depths of 1.5, 4.5, and 7.5 km for Galeras and 50 m for the other volcanoes. In our Monte Carlo simulations, we used the parameter values listed in Table 2, various  $l_0$  and  $Q_i$  values and layer boundary depths (see section 2.5), and linearly interpolated  $V_s$  structures using the  $V_s$  values listed in Table 3. The Monte Carlo method of Yoshimoto (2000) does not include surface waves. The dominance of *S* waves in seismograms in a high-frequency band of 5–10 Hz has been indicated by the finite-difference full-waveform simulation study of Morioka et al. (2017) as well as ASL studies assuming isotropic radiation of *S* waves (e.g., Battaglia & Aki, 2003; Kumagai et al., 2010, 2019; Ogiso & Yomogida, 2015). Therefore, our use of the Monte Carlo simulations without including surface waves is justified.

#### 2.5. Model Estimates

We exactly followed the grid-search approach of Kumagai et al. (2018) to obtain the best fitting layered model at Galeras. Observed envelope widths were divided into source depth ranges of 0–3, 3–6, and 6–9 km denoted as n = 1, 2, and 3, respectively. The simulated envelope widths estimated from envelope simulations at source depths of 1.5, 4.5, and 7.5 km were fit to the observed envelope widths in the respective depth ranges using the polynomial function:

$$p_n^k(r) = \sum_{j=1}^m {}_k a_n^j r^j,$$
 (1)

where *r* is source-station distance,  $p_n^k$  represent the simulated envelope widths,  $_k a_n^j$  are coefficients for the polynomial function of the *n*th depth range in the *k*th model, and *m* is the polynomial order. We used the following normalized residuals:

$$R_{n}^{k} = \frac{\sum_{l=1}^{N} \left\{ p_{n}^{o}(r_{l}) - p_{n}^{k}(r_{l}) \right\}^{2}}{\sum_{l=1}^{N} \left\{ p_{n}^{o}(r_{l}) \right\}^{2}},$$
(2)

where  $r_l$  is the source-station distance of the *l*th data points,  $p_n^o$  is the observed envelope widths in the *n*th depth range, and N is the number of data points. We evaluated the total normalized residual as  $R^k$ 

Table 5
Values of Scattering Mean Free Path $(l_0)$ , the Quality Factor of Medium Attenuation for S Waves $(Q_i)$ , and Layer Boundary
Depths Used for Trial Models in Our Monte Carlo Simulations for Kirishima, Unzen, Bandai, and Iwate

Layer number	Layer boundary depth (km)	<i>l</i> <sub>0</sub> (km)	Qi
1	0.5, 0.75, 1.0, 1.25	0.25, 0.5, 1, 10	10, 25, 50, 100, 200
2	5.0	1, 10, 100, 1,000, 5,000	25, 50, 100, 200





**Figure 3.** Envelope widths of VT earthquakes plotted against sourcestation distances at (a) Galeras, (b) Nevado del Ruiz, and (c) Taal (data at Nevado del Ruiz and Taal are from Kumagai et al., 2018). Source depths, measured from 3.4 and 4.4 km above sea level at Galeras and Nevado del Ruzi, respectively, and sea level at Taal, are shown by the color scale.

 $= (R_1^k + R_2^k + R_3^k)/3$ . Three layers with boundaries at depths of 1, 5, and 15 km with  $l_0$  and  $Q_i$  values increasing in deeper layers were assumed (Table 4). We adopted combinations of  $l_0$  and  $Q_i$  values in the three layers that provided  $R^k < 0.5$ . Using these combinations, we adjusted the boundary depth between the first and second layers to either 0.5 or 2 km and that between the second and third layers to either 4 or 6 km. We estimated the averaged  $l_0$  and  $Q_i$  model as

$$\bar{l}_{0} = \frac{\sum_{k=1}^{M} w_{k} l_{0}^{k}}{\sum_{k=1}^{M} w_{k}}, \quad \overline{Q}_{i} = \frac{\sum_{k=1}^{M} w_{k} Q_{i}^{k}}{\sum_{k=1}^{M} w_{k}}.$$
(3)

Here  $l_0^k$  and  $Q_i^k$  are  $l_0$  and  $Q_i$  in each layer for the *k*th model,  $w_k = e^{-R^k}$ , and *M* is the number of models. The layer boundary depths were estimated as the weighted averages of the adjusted depths.

At Kirishima, Unzen, Bandi, and Iwate, we followed the same procedure but used only one depth range (0–3 km) for the observed envelope widths and two layers (trial  $l_0$  and  $Q_i$  values and boundary depths are listed in Table 5) and a source depth of 50 m in the Monte Carlo simulations. We conducted a grid search for the trial models to evaluate  $R^k = R_1^k$  at each volcano. We estimated the averaged models by using Equation 3 with the layer boundary depths taken as the weighted averages of those determined for the individual volcanoes.

#### 3. Results

#### 3.1. Observed Envelope Widths

Observed envelope widths at Galeras were smaller than those estimated at Nevado del Ruiz and Taal (Figure 3). Although the envelope widths at Galeras did not obviously increase up to r of around 5 km as observed at Nevado del Ruiz and Taal, nearly constant values were observed at larger distances at all three volcances. As shown by Kumagai et al. (2018), the envelope width increases monotonically with increasing r when uniform distributions of  $l_0$  and  $Q_i$  are assumed. Therefore, a depth-dependent structure is required to explain the observed envelope widths at Galeras.

Observed envelope widths at Kirishima, Unzen, Bandai, and Iwate (Figure 4) generally increased up to *r* of around 5 km, and then they were nearly constant or slightly increased. Data scatter along the trends was largest at Iwate and least at Bandai.

We note that envelope widths depend on the source-station distance as well as the source depth as indicated by Kumagai et al. (2018). The observed envelope widths at Kirishima, Unzen, Bandai, and Iwate were relatively larger than those at Galeras at the same source-station distance (Figures 3a and 4). The sources of VT earthquakes that we used at Galeras (Figure 1) were deeper than the active sources at Kirishima, Unzen, Bandai, and Iwate (Figure 2). These shallow active sources were strongly affected by shallow heterogenous layers, which resulted in larger envelope widths.





Figure 4. Envelope widths of active sources plotted against source-station distance at (a) Kirishima, (b) Unzen, (c) Bandai, and (d) Iwate.

#### 3.2. One-Dimensional Models

Our results of the grid search using polynomial order m = 3 at Galeras are shown in Figure S1 in the supporting information. This polynomial order was used at Taal and Nevado del Ruiz (Kumagai et al., 2018) and provided the minimum  $R^k$  values compared with those of other orders at Galeras. Using the envelope widths of models with residuals within 1.1 times the minimum  $R^k$  values (green lines in Figure 5), we obtained the averaged  $l_0$  and  $Q_i$  models shown in Figure 6 and listed in Table 6. Our estimated structures at Galeras indicated that the first layer was 0.5 km thick with  $l_0 = 710$  m and  $Q_i = 120$ , whereas  $l_0$  was larger than  $10^2$  km and  $Q_i$  was around 100–200 in the second and third layers.

Our results of the grid search with m = 5 at Kirishima, Unzen, Bandai, and Iwate are shown in Figures S2–S5 and Table S1. We used m = 5 because the residuals were minimized at this order or did not change greatly when higher orders were used (Figure S6). Using the envelope widths of models with residuals within 1.1 times the minimum  $R^k$  values (green lines in Figure 7), we obtained the averaged models shown in Figure 8 and listed in Table 6. Our estimated models at these volcanoes were very similar: The first layers were 600–900 m thick with  $l_0 = 250$  m and  $Q_i$  around 25, and the second layers had  $l_0$  around 10<sup>3</sup> km and  $Q_i$  around 50–70. We note that  $l_0 = 250$  m was the minimum value explored in our trial models because of the long computational times required in the Monte Carlo simulations of such strongly scattering media. It is therefore possible that  $l_0$  in the first layer could be smaller than our estimates, although our estimations reasonably fit the observed trends at these volcanoes (Figure 7).

#### 3.3. Mapping of Envelope Width Residuals

The 2D SWF (Del Pezzo et al., 2016), derived from energy distributions in diffusive wavefields, is given as

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**Figure 5.** Observed and simulated envelope widths at Galeras. Blue circles indicate observed envelope widths for VT earthquakes in three source-depth ranges. Red solid and dashed lines are the envelope widths simulated by the model with the minimum residual and their fits by Equation 1, respectively. Green solid and dashed lines are the envelope widths of models with residuals within 1.1 times the minimum residual and their fits by Equation 1, respectively.

$$\kappa_{l}^{k}(x, y) = \frac{1}{2\pi\sigma_{x}\sigma_{y}} \exp\left[-\left\{\frac{(x-x_{l})^{2}}{2\sigma_{x}^{2}} + \frac{(y-y_{l})^{2}}{2\sigma_{y}^{2}}\right\}\right] + \frac{1}{2\pi\sigma_{x}\sigma_{y}} \exp\left[-\left\{\frac{(x-x_{k})^{2}}{2\sigma_{x}^{2}} + \frac{(y-y_{k})^{2}}{2\sigma_{y}^{2}}\right\}\right] + \frac{1}{4\pi\sigma_{x}\sigma_{y}} \exp\left[-\left\{a(x-x_{0})^{2} + 2b(x-x_{0})(y-y_{0}) + c(y-y_{0})^{2}\right\}\right]$$
(4)

Here

$$a = \frac{\cos^2\theta}{12\sigma_x^2} + \frac{\sin^2\theta}{2\sigma_y^2}, \quad b = -\frac{\sin^2\theta}{24\sigma_x^2} + \frac{\sin^2\theta}{4\sigma_y^2}, \quad c = \frac{\sin^2\theta}{24\sigma_x^2} + \frac{\cos^2\theta}{2\sigma_y^2}, \quad (5)$$

$$\sigma_x = \delta_x d, \ \sigma_y = \delta_y d, \ x_0 = \frac{x_k + x_l}{2}, \ y_0 = \frac{y_k + y_l}{2}, \tag{6}$$

where x and y are the Cartesian coordinates;  $(x_k, y_k)$  and  $(x_l, y_l)$  are source and receiver locations, respectively; d is the source-station distance;  $\delta_x = \delta_y = 0.2$ ; and  $\theta = \tan^{-1}\{(y_k - y_l)/(x_l - x_k)\}$ . Although the SWF was only given by Del Pezzo et al. (2016) for sources and stations positioned on the x-axis, they actually used the general form for the arbitrary source and station positions given in Equation 4. Because the sources and stations were located in the first heterogenous layer at Kirishima, Unzen, Bandai, and Iwate, and because the scatter in observed envelope widths is sensitive to the first layer (Kumagai et al., 2018), the SWF can be used to map the individual residual values between the observed envelope widths and those calculated with the best fit layered model (Table S1) at each volcano. However, the SWF may not be used for the envelope-width residuals at Galeras, Taal, or Nevado del Ruiz, because the sources are located in the second and third less heterogeneous layers and the diffusive wavefields between the sources and stations are not entirely established.

Following Sketsiou et al. (2020), we used both a regionalization approach and a tomography approach to map the envelope-width residuals using the SWF. In the regionalization approach, we estimated the distributions of spatially weighted envelope-width residuals defined as

$$W(n_x, n_y) = \sum_{l=1}^{N_l} \sum_{k=1}^{N_k} P_l^k(n_x, n_y) / (\Delta x \Delta y),$$
(7)

where

$$P_l^k(n_x, n_y) = K_l^k(n_x, n_y) \Delta p_l^k$$
(8)

$$K_{l}^{k}(n_{x}, n_{y}) = \frac{\kappa_{l}^{k}(n_{x}, n_{y})}{\sum_{i=1}^{N_{x}} \sum_{j=1}^{N_{y}} \kappa_{l}^{k}(i, j)}.$$
(9)

Here  $\Delta p_l^k$  is the envelope-width residual for the *l*th event at the *k*th station;  $N_k$  and  $N_l$  are the numbers of stations and events, respectively;  $n_x$  and  $n_y$  represent the grid points;  $N_x$  and  $N_y$  are the numbers of grid points in the *x* and *y* directions, respectively; and  $\Delta x$  and  $\Delta y$  are the grid intervals in the *x* and *y* directions, respectively. In the tomography





Figure 6. Estimated models with residuals within 1.1 times the minimum residual (black lines) and the averaged model (red line), respectively, at Galeras.

approach, we solved the following observation equation to estimate  $P_l^k$   $(n_x, n_y)$  by the damped least squares method (e.g., Menke, 1989):

$$\mathbf{d} = \mathbf{A}\mathbf{m} \tag{10}$$

where

$$\mathbf{d} = \left[\Delta p_1^1 \cdots \Delta p_{N_l}^{N_k} 0 \cdots 0\right]^T,\tag{11}$$

$$\mathbf{A} = \begin{bmatrix} K_{1}^{1}(1, 1) & \cdots & \cdots & K_{1}^{1}(N_{x}, N_{y}) \\ \vdots & \ddots & \vdots \\ \vdots & & \ddots & \vdots \\ K_{N_{l}}^{N_{k}}(1, 1) & \cdots & \cdots & K_{N_{l}}^{N_{k}}(N_{x}, N_{y}) \\ \lambda & 0 & \cdots & 0 \\ \vdots & \ddots & & \vdots \\ \vdots & & \ddots & 0 \\ 0 & \cdots & 0 & \lambda \end{bmatrix},$$
(12)

$$\mathbf{m} = \left[P_1^1(1, 1) \cdots P_{N_l}^{N_k}(N_x, N_y)\right]^T.$$
(13)

Here *T* denotes transpose and  $\lambda$  is the damping factor.

At each volcano, we plotted  $\Delta p_l^k$  along each ray path (Figure 9) and the number of rays in each grid with  $\Delta x = \Delta y = 500$  m (Figure S7). We also constructed envelope-width residual maps by the regionalization and tomography approaches (Figures 10 and 11, respectively). In the tomography approach, we used  $\lambda = 0.2$  at all four volcanoes (Figure S8). The maps obtained by the two approaches at the individual volcanoes display similar positive and negative locations, but the short-wavelength patterns are enhanced in the maps obtained by the tomography approach. As discussed by Sketsiou et al. (2020), the regionalization approach tends to be affected by greatly deviated data or outliers. The differences in our maps obtained by the two approaches may be caused by such effects. In the following discussion, we use the maps obtained by the tomography approach (Figure 11).

At Kirishima, areas of positive residuals were located near the two shot positions (S1 and S5), whereas areas of negative residuals were located near the other shot positions (S2, S3, and S6) (see shot positions in Figure 2). Areas of positive residuals also existed near the summits of Shimoe-dake and Ohachi. At Unzen, an area of strong negative residuals existed beneath the western flank. At Bandai, areas of negative residuals were near some shot positions (S3, S4, and S5) and beneath the volcanic center and an area of positive residuals existed beneath the southeastern flank. At Iwate, broad areas of positive and negative residuals

existed along the eastern and western flanks, respectively.

#### Table 6

Averaged Models of the Scattering Mean Free Path $(l_0)$ and the Quality
Factor of Medium Attenuation for S Waves (Q <sub>i</sub> ) at Galeras, Kirishima,
Unzen, Bandai, and Iwate Volcanoes

Volcano	Layer boundary depth (km)	<i>l</i> <sub>0</sub> (km)	Qi
Galeras	0.5	0.71	120
	5.0	200	150
	15.0	530	230
Kirishima	0.69	0.25	25
	5.0	2,000	50
Unzen	0.83	0.25	25
	5.0	2,000	72
Bandai	0.58	0.25	25
	5.0	2,000	50
Iwate	0.91	0.25	27
	5.0	1,900	66

### 4. Discussion

Our grid search results of the observed envelope widths indicated that a highly heterogeneous and attenuative surface layer exists to around 1 km depth at all studied volcanoes (Figures 6 and 8). Although the envelope-width residuals mapped using the 2D SWF display various spatial distributions at Kirishima, Unzen, Bandai, and Iwate (Figure 11), our estimated 1D models are similar to those at Vesuvius (Wegler, 2004), Nevado del Ruiz, and Taal (Kumagai et al., 2018). These results clearly indicate that strong heterogeneities at volcanoes are surface characteristics.

The values of  $l_0$  and  $Q_i$  estimated in the first layer for the VT earthquakes and active sources were systematically different:  $l_0 = 700-1,000$  m and  $Q_i = 100-200$  at Galeras, Nevado del Ruiz, and Taal (VT earthquakes)





Figure 7. Observed and simulated envelope widths at (a) Kirishima, (b) Unzen, (c) Bandai, and (d) Iwate. Symbols as in Figure 5.

and  $l_0 = 250$  m and  $Q_i$  of around 25 at Kirishima, Unzen, Bandai, and Iwate (active sources); the estimated  $l_0$ and Q<sub>i</sub> values in the second layer were also systematically different (Figure 12). These systematic differences may result from the different resolutions attained for VT earthquakes and active sources. For natural earthquakes, the observed envelope widths were divided into source-depth ranges of 0-3, 3-6, and 6-9 km for comparison with those calculated for various structures at depths of 1.5, 4.5, and 7.5 km. We used this approach because all possible combinations of source-station pairs were not practical to compute in our grid search by the Monte Carlo method, because it would require extensive computational time. Therefore, the 1D models estimated from natural earthquakes may be poorly resolved, especially at shallow depths, where fewer earthquakes occurred. On the other hand, the active sources were located at shallow depths of around 50 m. There may be depth-dependent structures within the first layer; particularly, the upper most part may be highly heterogeneous and attenuative. Therefore, the results from active sources may be strongly affected by this uppermost part of the first layer, which also affected estimates in the second layer. Despite such systematic differences, our results clearly show that a thin, surficial, heterogeneous layer underlain by less heterogenous layers is a common feature. As discussed by Kumagai et al. (2018), this surface layer may correspond to low-velocity surface layers commonly observed in tomographic images (e.g., Aoki et al., 2009; Molina et al., 2005; Nishi, 2002; Tanaka, Hamaguchi, Nishimura, et al., 2002; Tomatsu et al., 2001; Yamawaki et al., 2004; Zollo et al., 2003) and low-resistivity surface layers inferred from electromagnetic surveys at various volcanoes (e.g., Aizawa et al., 2008; Alanis et al., 2013; Kanda et al., 2008; Nurhasan et al., 2006; Yamaya et al., 2013). This layer may consist of unconsolidated and/or highly fractured materials, hydrothermal fluids, and altered clay minerals, all of which are common to volcanic edifices and produce strongly scattering wavefields.

At Galeras, 3D distributions of scatterers estimated by Carcolé et al. (2006) using the method of Nishigami (1991) and assuming single scattering indicate that strong scattering regions are broadly





Figure 8. Estimated models at (a) Kirishima, (b) Unzen, (c) Bandai, and (d) Iwate. Symbols as in Figure 6.

distributed beneath the summit at depths between 4 and 8 km. However, we do not observe this feature in our estimated 1D model at Galeras. This difference is likely due to their assumption of single scattering, which is inappropriate at Galeras because of the dominance of multiple scattering, especially at shallow depths in volcanic edifices.

Although the 1D models estimated at Kirishima, Unzen, Bandai, and Iwate were quite similar, the scatter in the observed envelope widths showed different features among these volcanoes (Figures 4 and 7). The observed envelope widths at Bandai were less scattered than those at the other volcanoes, of which Iwate showed the broadest scatter. Their envelope-width residual maps clearly show different spatial distributions (Figure 11). Despite the less scattered envelope widths at Bandai, distinct areas of negative and positive residuals were observed there (Figure 11c). At Iwate, the scattered envelope-width residuals appeared as broad areas of positive and negative residuals in the edifice (Figure 11d). Negative and positive residuals were distributed in the volcano flanks at Kirishima (Figure 11a), and an area of strong negative residuals existed at Unzen (Figure 11b).

Positive (negative) envelope-width residuals represent smaller  $l_0$  (larger) and/or larger (smaller)  $Q_i$  values from the 1D model (Kumagai et al., 2018). In other words, positive residuals indicate more heterogeneous and/or less attenuative bodies and negative residuals less heterogeneous and/or more attenuative bodies. Indeed, we noted correlations between our residual maps and 3D  $V_p$  tomographic images from previous studies at these volcanoes (Nishi, 2002; Tanaka, Hamaguchi, Nishimura, et al., 2002; Tomatsu et al., 2001; Yamawaki et al., 2004). Comparison of our map at Kirishima (Figure 11a) and the *P* wave velocity image of Tomatsu et al. (2001) at depths between -1.5 and -0.5 km below sea level (Figures 10a and 10b of Tomatsu et al., 2001) indicates that the areas of positive residuals observed near shot locations S1 and S5 (Figure 2a) and the summit of Shimoe-dake correspond to low velocities of around  $V_p = 2$  km/s. We also found at Kirishima that the areas of negative residuals near shot locations S2 and S6 correspond to high





**Figure 9.** Envelope-width residuals ( $\Delta p$ , the residuals between the observed envelope widths and those calculated with the best-fit layered model) plotted along ray paths using the color scale shown at the bottom of each map: (a) Kirishima, (b) Unzen, (c) Bandai, and (d) Iwate.

velocities and that near S3 corresponds to low velocities. At Unzen, the clear area of negative residuals in the western flank (Figure 11b) corresponds to high velocities around  $V_p = 5$  km/s at depths between 0 and 1 km below sea level (Figures of 7a and 7b of Nishi, 2002), which were interpreted as pre-Unzen volcanic rock based on the geological map of Hoshizumi et al. (1999). We also note an area of positive residuals around Fugen-dake (Figure 11b), where the 1991 Unzen eruptions occurred. At Bandai, areas of negative residuals at shot locations S3, S4, and S5 and in the volcanic center and the area of positive residuals in the southeastern flank (Figures 2c and 11c) correspond to high- and low-velocity anomalies, respectively, at depths between -1 and 1 km below sea level in tomographic images (see Figures 9a–9c of Yamawaki et al., 2004). Although the tomographic images at Iwate of Tanaka, Hamaguchi, Nishimura, et al. (2002) were shown only in limited plan view and cross section, they indicate low- and high-velocity anomalies at



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**Figure 10.** Envelope-width residual maps based on the regionalization approach using Equations 7 and 8 ( $\Delta x = \Delta y = 500$  m): (a) Kirishima, (b) Unzen, (c) Bandai, and (d) Iwate. *W* values (Equation 7) in individual grids with more than five rays are plotted using the color scale shown at the bottom of each map. The gray triangles in Figure 10a are the locations of Shimoe-dake (Sm) and Ohachi (Oh) and that in Figure 10b is the location of Fugen-dake, where the 1991 Unzen eruptions occurred.

shallow depths in the eastern and western flanks, respectively, which may correspond to the respective areas of positive and negative residuals (Figure 11d).

These comparisons consistently indicate that areas of positive and negative envelope-width residuals correspond to low- and high-velocity anomalies except for one area near shot location S3 at Kirishima. These correlations suggest that areas of positive residuals correspond to heterogeneous rocks comprising unconsolidated materials and/or fractures, whereas areas of negative residuals correspond to less heterogenous volcanic or basement rocks. To quantify  $l_0$  and  $Q_i$  values in these anomalies, a tomographic approach based on inversion of  $\Delta p_l^k$  may prove useful. However, one of the difficulties in performing such an inversion is to construct the data kernel, in which the envelope widths for perturbations of  $l_0$  and  $Q_i$  in individual anomalies must be calculated for all combinations of source-station pairs using the Monte Carlo





**Figure 11.** Envelope-width residual maps based on the tomography approach using Equations 7 and 10 ( $\Delta x = \Delta y = 500$  m): (a) Kirishima, (b) Unzen, (c) Bandai, and (d) Iwate. Symbols as in Figure 10.

method, which is computationally expensive. Takeuchi (2016) developed an efficient method for calculating envelope seismograms and their partial derivatives with the Monte Carlo method, which was applied to estimate 3D scattering and attenuation structures in the lithosphere in southwest Japan (Ogiso, 2019). However, this method uses the Born approximation assuming single scattering; therefore, it is not applicable to the simulation of seismograms at volcanoes, where multiple scattering dominates at shallow depths. Akande et al. (2019) used 3D sensitivity kernels based on the analytical equation of Paasschens (1997) for attenuation imaging at Campi Flegrei caldera (Italy). This approach was also used by Shito et al. (2020) to estimate scattering and attenuation structures beneath Kyushu, Japan. The 3D kernels, however, are based on uniform  $l_0$  and  $Q_i$  distributions, but the assumption of uniform  $l_0$  and  $Q_i$  is not appropriate at volcanoes. Future studies should therefore develop a tomographic method with efficient





**Figure 12.** Comparison of estimated 1D models at various volcanoes. Red and blue lines indicate models based on VT earthquakes and active sources, respectively. Models at Nevado del Ruiz (Ruiz) and Taal are from Kumagai et al. (2018).

calculations of partial derivatives of  $\Delta p_l^k$  to estimate 3D  $l_0$  and  $Q_i$  distributions beneath volcanoes, which would contribute to our better understanding of magmatic and hydrothermal systems.

## 5. Conclusions

The envelope widths of VT earthquakes at Galeras and active sources produced by dynamite blasts at Kirishima, Unzen, Bandai, and Iwate were estimated by using high-frequency (5-10 Hz) seismograms. We compared the observed envelope widths with those simulated by the Monte Carlo method and estimated the 1D  $l_0$  and  $Q_i$  models at each volcano using the grid-search approach of Kumagai et al. (2018). Our models indicated the existence of a heterogeneous  $(l_0 = 250 -$ 710 m) and attenuative ( $Q_i = 25-120$ ) surface layer with a thickness of 500-900 m, which is underlain by less heterogeneous layers  $(l_0 = 10^2 - 10^3 \text{ km} \text{ and } Q_i = 50 - 200)$ . This feature is similar to those found at Vesuvius (Wegler, 2004), Taal, and Nevado del Ruiz (Kumagai et al., 2018), indicating that such strong structural heterogeneities at volcanoes are surface characteristics. These surface layers correspond to low-velocity and low-resistive surface layers that may consist of unconsolidated and/or highly fractured materials, hydrothermal fluids, and altered clay minerals.

Using the 2D SWF of Del Pezzo et al. (2006), we mapped the residuals between the observed envelope widths and those calculated with the best fit layered models at Kirishima, Unzen, Bandai, and Iwate. These maps show areas of negative and positive residuals distributed beneath the volcano flanks at Kirishima, an area of strong negative residuals at Unzen, distinct areas of negative and positive residuals at Bandai, and broad areas of positive and negative residuals in the eastern and western flanks of Iwate, respectively. The spatial distributions of these residuals correlate with *P* wave velocity tomographic images at each volcano: Areas of positive residuals correspond to low-velocity anomalies and thus heterogeneous, strongly scattering rocks, whereas areas of negative residuals correspond to high-velocity anomalies, a tomographic approach based on inversion of the envelope-width residuals may prove useful, but it would require an efficient method to calculate envelope seismograms and their partial derivatives with respect to  $l_0$  and  $Q_i$  perturbations. Our study demonstrates that the envelope width is useful for improving the characterization of scattering and attenuation structures beneath volcanoes.

## Data Availability Statement

Observed and simulated envelope widths and source and station information for our studied volcanoes are found online (http://www.eps.nagoya-u.ac.jp/~kumagai/2020JB020249.zip). Waveform data at Galeras and those at Kirishima, Unzen, Bandai, and Iwate are available upon request from the Colombian Geological Survey (https://www2.sgc.gov.co/Paginas/servicio-geologico-colombiano.aspx) and the Earthquake Research Institute, The University of Tokyo (http://www.eri.u-tokyo.ac.jp/kyodoriyou/facilities/), respectively.

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