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2 **Effect of beam current on defect formation by high-temperature**
3 **implantation of Mg ions into GaN**

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16 We evaluated the beam current dependence of defect formation during Mg ion implantation
17 into GaN at a high temperature of 1100°C with two beam currents. Photoluminescence spectra suggest
18 that low-beam-current implantation reduced the vacancy concentration and activated Mg to a greater
19 extent. Moreover, scanning transmission electron microscopy analysis showed that low-beam-current
20 implantation reduced the density of Mg segregation defects with inactive Mg and increased the density
21 of intrinsic dislocation loops, suggesting decreases in the densities of Ga and N vacancies. The
22 formation of these defects depended on beam current, which is an important parameter for defect
23 suppression.
24

25 Gallium nitride (GaN) has a wide bandgap, high breakdown electric field, and high electron
26 mobility.¹⁻³⁾ To take advantage of its excellent physical properties, GaN is applied to the fabrication
27 of power devices such as metal-oxide-semiconductor field-effect transistors and the edge termination
28 structure.⁴⁻¹⁰⁾ These devices require ion implantation (I/I) to form localized n-type and p-type regions.
29 However, the fabrication of p-type Mg-ion-implanted GaN remains a challenge. Recently, the
30 demonstration of p-type conductivity and the temperature-dependent Hall effect measurement have
31 been reported with annealing above 1400°C at a nitrogen pressure of 1 GPa.¹¹⁻¹⁴⁾ They have indeed
32 been breakthroughs; however, for the practical implementation of novel devices, the ambient-pressure
33 fabrication is desirable because there are strict regulations in the use of high-pressure and high-
34 temperature systems in many countries.

35 Defect control in I/I is essential in fabricating p-type Mg-ion-implanted GaN at ambient
36 pressure. It has been reported that many primary defects, such as vacancies and interstitials, are formed
37 during I/I, and that annealing to remove them results in the formation of various secondary defects.

38 For instance, positron annihilation spectroscopy studies indicated that Mg I/I generates a high
39 concentration of the Ga and N vacancy complex ($V_{\text{Ga}}V_{\text{N}}$).^{15,16} During annealing above 1000°C, these
40 vacancies begin to agglomerate and form vacancy clusters that act as carrier traps, such as $(V_{\text{Ga}})_3(V_{\text{N}})_3$.
41 Moreover, the implanted Mg ions segregated above the concentration of $1 \times 10^{19} \text{ cm}^{-3}$, indicating a
42 high concentration of inactive Mg.¹⁷ Thus, the formation of secondary defects, such as vacancy
43 clusters and Mg segregation defects must be suppressed to obtain high p-type conductivity.

44 To resolve these issues, we propose a technique of I/I at a high temperature (HT). In this
45 HT-I/I technique, primary defects are expected to be annihilated either by their instantaneous return to
46 the original site or diffusion to the surface without forming secondary defects. Previously, the
47 dependence of implanted ion concentration on annealing temperature was reported.^{18–24} In particular,
48 the results of HT-I/I above 1000°C suggested the activation effect of Mg and the suppression of defect
49 formation compared with conventional I/I without intentional heating.¹⁸ If HT-I/I can suppress defect
50 formation, beam current should be an important parameter because it is expressed as the number of
51 defects introduced per unit time. Therefore, low-beam-current implantation should lead to more
52 defects being removed by annealing during heating. In this study, we performed HT-I/I with different
53 beam currents of two orders of magnitude and evaluated the dependence of defect formation on beam
54 current.

55 Mg ions were implanted into unintentionally doped (UID) Ga-polar GaN grown on a free-
56 standing GaN substrate. UID-GaN of 3 μm thickness was grown by metalorganic vapor phase epitaxy
57 (MOVPE). Impurities such as carbon, oxygen, and silicon are at background levels or lower than 10^{16}
58 cm^{-3} . As a protective film during HT-I/I, a 50-nm-thick AlN layer was deposited by MOCVD at 500°C
59 on both sides. Mg ions were implanted at energies of 100 and 40 keV with the corresponding doses of
60 1.3×10^{14} and $2.8 \times 10^{13} \text{ cm}^{-2}$, respectively. A sharp box-shaped profile with a depth of 100 nm was
61 formed at a Mg concentration of $1.0 \times 10^{19} \text{ cm}^{-3}$. The angles of the samples were set to 7 and 15° from
62 the c- and m-axes, respectively, to prevent channeling. The temperature during implantation was
63 1100°C, which was monitored using a thermocouple set on the surface of the sample. A sample was
64 placed on a carbon susceptor and heated by infrared irradiation from the back surface. The beam
65 current was set to 100 and 1 μA . In our implantation device, a pulse of ions was applied for 0.06 ms
66 over a period of 3.25 ms. Beam current densities of 100 and 1 μA were obtained as 31.83 and 0.318
67 μAcm^{-2} , respectively. The implantation duration was set to 68 and 6,800 s for 100 and 1 μA ,
68 respectively, to keep the total dose constant. To regulate the thermal annealing of 100 and 1 μA as-
69 implanted samples, 100 μA as-implanted samples were annealed in vacuum at 1100°C for 6,732 s.
70 These samples are referred to as 100 μA +heated, and the other two samples as 1 and 100 μA as-
71 implanted. After removing the AlN protective film, a 100-nm-thick AlN protective film for post
72 annealing was again deposited. All the samples were annealed at 1250°C for 45 s in N_2 gas at
73 atmospheric pressure. In this paper, we denote the samples obtained after HT-I/I and annealing at

74 1250°C for 45 s as “as-implanted” and “post-annealed”, respectively. To evaluate the effect of beam
75 current on defect formation, we determined the luminescence characteristics and analyzed the defects
76 by using scanning transmission electron microscopy (STEM).

77 Figures 1(a) and 1(b) show the results of secondary ion mass spectrometry of as-implanted
78 and post-annealed samples, respectively. Compared with the 100 μA and 100 μA +heated as-implanted
79 samples, the Mg concentration of the 1 μA as-implanted sample decreased at a depth of around 100
80 nm and increased at depths larger than 200 nm owing to diffusion. Furthermore, post annealing
81 reduced the Mg concentration peak around 50 nm; the Mg ions in the reduced peak may have diffused
82 to the surface. These Mg diffusion behaviors are clearly dependent on the beam current, which may
83 be due to the different density and type of defects formed by 1 μA and 100 μA implantation.

84 We determined the luminescence characteristics of as-implanted and post-annealed samples
85 by photoluminescence (PL) measurement at 77 K. A 325 nm He–Cd laser was utilized at an excitation
86 density of 260 Wcm^{-2} . The luminescence characteristics observed at 77 K are as follows. The high
87 intensity of near-band-edge (NBE) emission at around 3.47 eV suggests a decrease in the concentration
88 of nonradiative recombination centers (NRCs).²⁵⁾ The donor–acceptor-pair (DAP) emission at around
89 3.28 eV suggests the activation of Mg. The green luminescence (GL) at around 2.35 eV is
90 characteristic of Mg-ion-implanted GaN, and the GL is identified as GL2.²⁵⁻²⁷⁾ The origin of GL2 is
91 attributed to the formation of N vacancies (V_{NS}),^{26,27)} which we also assume in our discussion. The PL
92 spectra of as-implanted and post-annealed samples are shown in Figs. 2(a) and 2(d), respectively. The
93 DAP, NBE, their LO replicas, and GL2 were clearly observed in all the samples.

94 In as-implanted samples [Fig. 2(a)], the GL2 intensities of the 1 and 100 μA samples were
95 almost the same but it was about twice as high as that of the 100 μA +heated sample. The 100
96 μA +heated sample was annealed for another 6,732 s after the implantation, whereas heating of the 100
97 and 1 μA samples was stopped as soon as the implantation was completed. Thus, some areas of the 1
98 and 100 μA samples are amorphous, and their GL2 intensities were lower than that of the 100
99 μA +heated sample. The DAP intensity of the 1 μA sample was higher than that of the 100 μA sample
100 but comparable to that of the 100 μA +heated sample. As in the case of GL2, the Mg ions introduced
101 into the last batch should not have been activated. The NBE intensity of the 1 μA sample was the
102 highest among all the samples, and the NBE intensities of the other two samples were comparable.
103 The evaluation of the as-implanted samples did not show that implantation at 1 μA was particularly
104 superior to that at 100 μA in terms of both GL2 and DAP intensities. However, a clearly higher NBE
105 intensity indicated that Mg could be implanted with fewer defects formed at the same dose.

106 In the case of post-annealed samples [Fig. 2(d)], the GL2 intensity of the 1 μA sample was
107 lower than those of the 100 μA and 100 μA +heated samples. These suggest that V_{NS} formation was
108 suppressed during post annealing since low-beam-current implantation decreased the V_{N} concentration
109 in the as-implanted samples. The DAP and NBE intensities of the 1 μA sample was sufficiently higher

110 than those of the 100 μA and 100 μA +heated samples. It is clear that implantation at 1 μA contributed
111 to the increases in DAP and NBE intensities, indicating the further activation of Mg and the decrease
112 in the concentration of V_{Ga} complexed with V_{NS} , which act as NRCs. These results showed that the
113 intensities of GL2, DAP, and NBE clearly depended on the beam current, not on the total annealing
114 time, and low-beam-current implantation is more effective than high-beam-current implantation in
115 further activating Mg and reducing the vacancy concentration.

116 To investigate the beam current dependence of the structural properties, we analyzed the
117 defects by annular dark-field scanning transmission electron microscopy (ADF-STEM). The ADF-
118 STEM cross-sectional images of as-implanted and post-annealed samples are shown in Figs. 3(a)–(c)
119 and 3(d)–(f), respectively. In this experiment, extended defects appearing as bright dots and loops were
120 observed. Bright dots were observed in all the samples and bright loops in post-annealed samples. We
121 investigated the characteristics of bright dots and loops by STEM and high-angle ADF (HAADF-
122 STEM) analyses.

123 Figure 4(a) shows an HAADF-STEM image of the extended defects indicated by a solid
124 triangle in Fig. 3(c). A triangular contrast, which has an upper (0001) facet and (11-23) side facets, can
125 be observed. The energy-dispersive X-ray spectroscopy (EDS) spectrum of the (0001) facet in the red
126 rectangle is shown in Fig. 4(b). This spectrum shows that the Mg concentration increased around the
127 boundary of (0001), indicating that the triangular defects have a high concentration of Mg. A similar
128 defect is confirmed in p-type GaN grown epitaxially by doping with a high concentration of Mg.^{17,28)}
129 Structurally, it is a pyramidal inversion domain (PID) with a high concentration of inactive Mg at the
130 (0001) facet and inhibits p-type conductivity. The triangular defects have features in common with the
131 PID, indicating that the PID-like Mg segregation defects also have a high concentration of inactive
132 Mg.

133 We utilized an inside–outside technique to identify whether the bright circles are extrinsic
134 or intrinsic dislocation loops by following the procedure described in ref. 29. We confirmed that the
135 bright circles were intrinsic dislocation loops. Such defects have been observed in Mg-implanted GaN
136 after UHPA at 1300°C for more than 30 min or at 1400°C for more than 5 min and after Mg
137 segregation defects had disappeared.^{29,30)}

138 We converted the number of these defects into areal density in the c-axis direction and
139 examined the effects of low-beam-current implantation, by comparing with the PL spectrum. Table I
140 shows the defect densities of as-implanted and post-annealed samples. The low-beam-current
141 implantation such as that at 1 μA significantly decreased the density of extended defects, which implies
142 a decrease in the density of vacancies and interstitial atoms. Some of the extended defects are PID-
143 like Mg segregation defects with inactive Mg; the fewer the PID-like Mg segregation defects, the
144 greater the extent of Mg being activated. In addition, the number of intrinsic dislocation loops
145 increased, and since the loops are formed by vacancy aggregation,²⁹⁾ vacancies probably disappeared.

146 These support the PL results suggesting a decrease in the density of vacancies and a higher activation
147 of Mg by low-beam-current implantation. Thus, the evaluations of PL and STEM indicate that low-
148 beam-current implantation reduces the defect density and further activates Mg. There are several
149 considerations from the perspective of secondary defect formation and beam current. The density of
150 primary defects formed by low-beam-current implantation is the same as that formed by high-beam-
151 current implantation, but the density of defects formed per unit time is lower. Furthermore, at a high
152 temperature, some of the formed primary defects should disappear owing to their diffusion to the
153 surface or annihilation. Therefore, we consider that the low density of primary defects per unit time
154 leads to the reduction of the probability of aggregation with other primary defects, and more primary
155 defects diffuse to the surface or are instantly annihilated.

156 In summary, we performed the HT-I/I of Mg into GaN with two beam currents. PL spectra
157 suggest that the reduction in beam current suppressed the vacancy concentration and activated Mg to
158 a greater extent. From the results of STEM analyses, the areal density of extended defects including
159 PID-like Mg segregation defects decreased and that of dislocation loops increased with the reduction
160 in the beam current. These results support the idea that low-beam-current implantation suppresses the
161 formation of vacancies and that the post annealing further activates Mg. We consider that the control
162 of the beam current can lower the annealing temperature and shorten the annealing time for p-type
163 Mg-implanted GaN.

164

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References

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183 1) T. Maeda, T. Narita, S. Yamada, T. Kachi, T. Kimoto, M. Horita, and J. Suda, *J. Appl. Phys.* **129**,
184 185702 (2021).
- 185 2) H. Fujikura, T. Konno, T. Kimura, Y. Narita, and F. Horikiri, *Appl. Phys. Lett.* **117**, 012103 (2020).
- 186 3) T. Narita, K. Tomita, Y. Tokuda, T. Kogiso, M. Horita, and T. Kachi, *J. Appl. Phys.* **124**, 215701
187 (2018).
- 188 4) R. Tanaka, S. Takashima, K. Ueno, H. Matsuyama, M. Edo, and K. Nakagawa, *Appl. Phys.*
189 *Express* **12**, 054001 (2019).
- 190 5) R. Tanaka, S. Takashima, K. Ueno, H. Matsuyama, and M. Edo, *Jpn. J. Appl. Phys.* **59**, SGGD02
191 (2020).
- 192 6) M. Yoshino, Y. Ando, M. Deki, T. Toyabe, K. Kuriyama, Y. Honda, T. Nishimura, H. Amano, T.
193 Kachi, T. Nakamura, *Materials* **12**, 689 (2019).
- 194 7) M. Matys, T. Ishida, K. Pil Nam, H. Sakurai, K. Kataoka, T. Narita, T. Uesugi, M. Bockowski, T.
195 Nishimura, Jun Suda, *Appl. Phys. Express* **14**, 074002 (2021).
- 196 8) M. Matys, T. Ishida, K. Pil Nam, H. Sakurai, T. Narita, T. Uesugi, M. Bockowski, J. Suda, and T.
197 Kachi, *Appl. Phys. Lett.* **118**, 093502 (2021).
- 198 9) J. R. Dickerson, A. A. Allerman, B. N. Bryant, A. J. Fischer, M. P. King, M. W. Moseley, A. M.
199 Armstrong, R. J. Kaplar, I. C. Kizilyalli, O. Aktas, and J. J. Wierer, *IEEE Trans. Electron Devices*
200 **63**, 419 (2016).
- 201 10) T. J. Anderson, J. D. Greenlee, B. N. Feigelson, J. K. Hite, F. J. Kub and K. D. Hobart, *ECS J.*
202 *Solid State Sci. Technol.* **5**, Q176 (2016).
- 203 11) H. Sakurai, M. Omori, S. Yamada, Y. Furukawa, H. Suzuki, T. Narita, K. Kataoka, M. Horita, M.
204 Bockowski, J. Suda, and T. Kachi, *Appl. Phys. Lett.* **115**, 142104 (2019).
- 205 12) H. Sakurai, T. Narita, M. Omori, S. Yamada, A. Koura, M. Iwinska, K. Kataoka, M. Horita, N.
206 Ikarashi, M. Bockowski, J. Suda, and T. Kachi, *Appl. Phys. Express* **13**, 086501 (2020).
- 207 13) K. Sierakowski, R. Jakiela, B. Lucznik, P. Kwiatkowski, M. Iwinska, M. Turek, H. Sakurai, T.
208 Kachi, and M. Bockowski, *Electronics* **9**, 1380 (2020).
- 209 14) M. H. Breckenridge, J. Tweedie, P. Reddy, Y. Guan, P. Bagheri, D. Szymanski, S. Mita, K.
210 Sierakowski, M. Bockowski, R. Collazo, and Z. Sitar, *Appl. Phys. Lett.* **118**, 022101 (2021).
- 211 15) A. Uedono, S. Takashima, M. Edo, K. Ueno, H. Matsuyama, H. Kudo, H. Naramoto, and S.
212 Ishibashi, *Phys. Status Solidi B* **252**, 2794 (2015).
- 213 16) A. Uedono, S. Takashima, M. Edo, K. Ueno, H. Matsuyama, W. Egger, T. Koschine, C.
214 Hugenschmidt, M. Dickmann, K. Kojima, S. F. Chichibu, S. Ishibashi, *Phys. Status Solidi B* **255**,
215 1700521 (2018).
- 216 17) A. Kumar, W. Yi, J. Uzuhashi, T. Ohkubo, J. Chen, T. Sekiguchi, R. Tanaka, S. Takashima, M.
217 Edo, and K. Hono, *J. Appl. Phys.* **128**, 065701 (2020).

- 218 18) M. Takahashi, A. Tanaka, Y. Ando, H. Watanabe, M. Deki, M. Kushimoto, S. Nitta, Y. Honda,
219 K. Shima, K. Kojima, S. F. Chichibu, and H. Amano, *Jpn. J. Appl. Phys.* **59**, 056502 (2020).
220 19) A. Wenzel, C. Liu, and B. Rauschenbach, *Mater. Sci. Eng. B* **59**, 191 (1999).
221 20) S. Whelan, M. J. Kelly, R. Gwilliam, C. Jeynes, and C. Bongiorno, *J. Appl. Phys.* **98**, 013515
222 (2005).
223 21) A. Lardeau-Falcy, M. Coig, M. Charles, C. Licitra, J. Kanyandekwe, F. Milési, J. Eymery, and F.
224 Mazen, *ECS Trans.* **80**, 131 (2017).
225 22) T. Niwa, T. Fujii, and T. Oka, *Appl. Phys. Express* **10**, 091002 (2017).
226 23) A. Kumar, K. Mitsuishi, T. Hara, K. Kimoto, Y. Irokawa, T. Nabatame, S. Takashima, K. Ueno,
227 M. Edo, and Y. Koide, *Nanoscale Res. Lett.* **13**, 403 (2018).
228 24) M. Sumiya, K. Fukuda, S. Takashima, S. Ueda, T. Onuma, T. Yamaguchi, T. Honda, and A.
229 Uedono, *J. Cryst. Growth* **511**, 15 (2019).
230 25) S. F. Chichibu, K. Shima, K. Kojima, S. Takashima, K. Ueno, M. Edo, H. Iguchi, T. Narita, K.
231 Kataoka, S. Ishibashi and A. Uedono, *Jpn. J. Appl. Phys.* **58**, SC0802 (2019).
232 26) K. Kojima, S. Takashima, A. Edo, K. Ueno, M. Shimizu, T. Takahashi, S. Ishibashi, A. Uedono,
233 and S. F. Chichibu, *Appl. Phys. Express* **10**, 061002 (2017).
234 27) M. A. Reshchikov, D. O. Demchenko, J. D. McNamara, S. Fernández-Garrido, and R. Calarco,
235 *Phys. Rev. B* **90**, 035207 (2014).
236 28) T. Narita, N. Ikarashi, K. Tomita, K. Kataoka, and T. Kachi, *J. Appl. Phys.* **124**, 165706 (2018).
237 29) K. Iwata, H. Sakurai, S. Arai, T. Nakashima, T. Narita, K. Kataoka, M. Bockowski, M. Nagao, J.
238 Suda, T. Kachi, and N. Ikarashi, *J. Appl. Phys.* **127**, 105106 (2020).
239 30) T. Nakashima, E. Kano, K. Kataoka, S. Arai, H. Sakurai, T. Narita, K. Sierakowski, M. Bockowski,
240 M. Nagao, J. Suda, T. Kachi and N. Ikarashi, *Appl. Phys. Express* **14**, 011005(2021).

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254 **Figure captions**

255

256 **Fig. 1.** (Color online) Depth profiles of Mg in (a) as-implanted samples and (b) post-annealed samples.
257 (annealed at 1250°C for 45 s)

258

259 **Fig. 2.** (Color online) PL spectra at 77 K for (a) as-implanted samples. (b)–(c) PL intensity axis on a
260 linear scale of (a). (d) PL spectra for post-annealed samples. (e)–(f) PL intensity axis on a linear scale
261 of (d)

262

263 **Fig. 3.** Cross-sectional ADF-STEM images of (a) 100 μA as-implanted sample, (b) 100 μA +heated as-
264 implanted sample, (c) 1 μA as-implanted sample, (d) 100 μA post-annealed sample, (e) 100 μA +heated
265 post-annealed sample, and (f) 1 μA post-annealed sample.

266

267 **Fig. 4.** (a) HAADF-STEM image of Mg segregation defect. (b) EDS spectra measured at the (0001)
268 facet of the triangle (red rectangle) and the GaN matrix.

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Table I. Areal density in the c-axis direction of as-implanted and post-annealed sample. All densities are calculated in the arrowed area of Fig. 3(a)

| Sample name | Extended defect ($\times 10^{10} \text{ cm}^{-2}$) | Dislocation loop ($\times 10^{10} \text{ cm}^{-2}$) |
|---------------------------|---|--|
| 100 μA | 13.70 / 13.50 | 0.00 / 0.23 |
| 100 μA +heated | 11.05 / 12.05 | 0.00 / 0.32 |
| 1 μA | 3.02 / 2.36 | 0.00 / 0.54 |

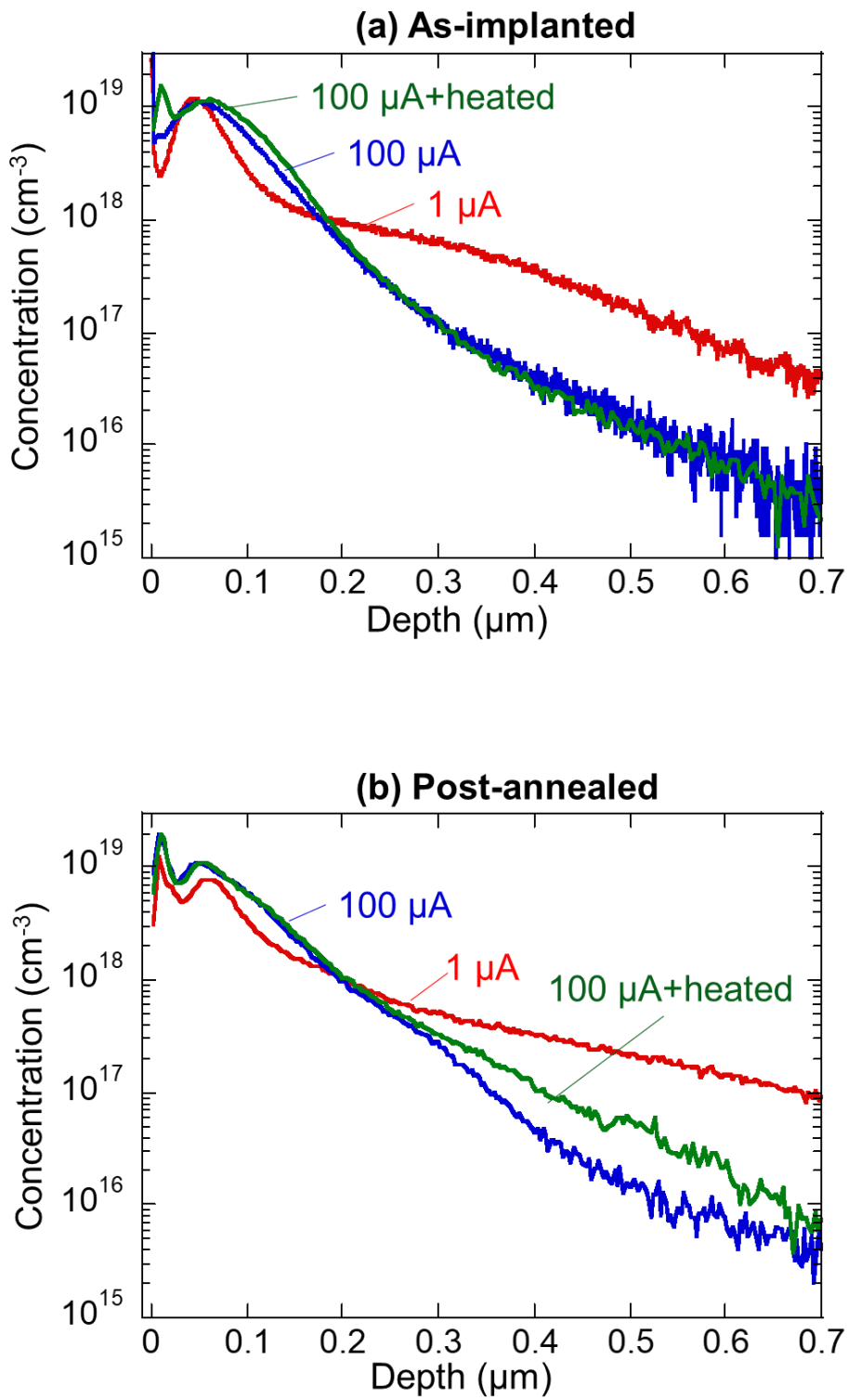


Fig. 1.

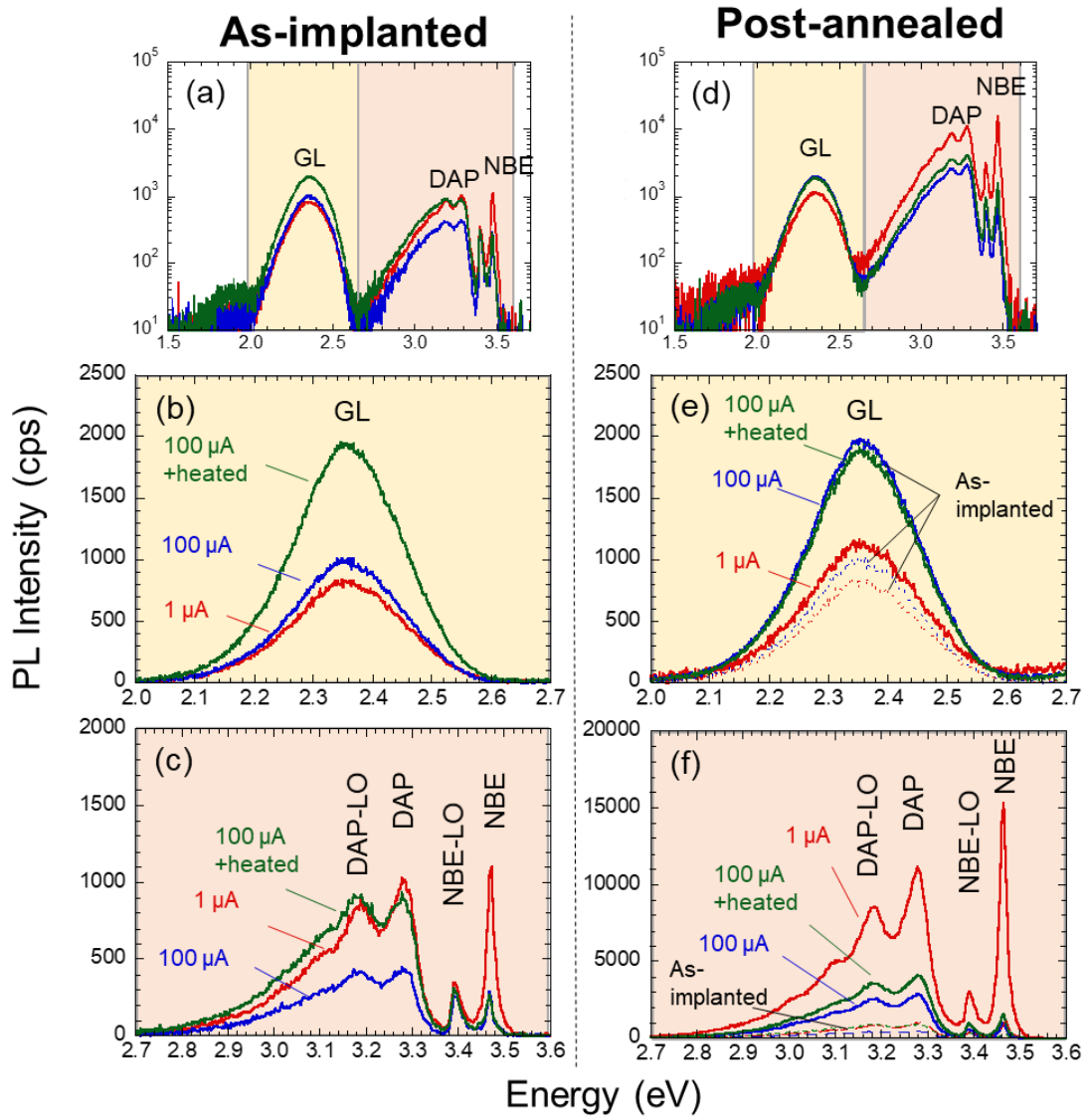


Fig. 2.

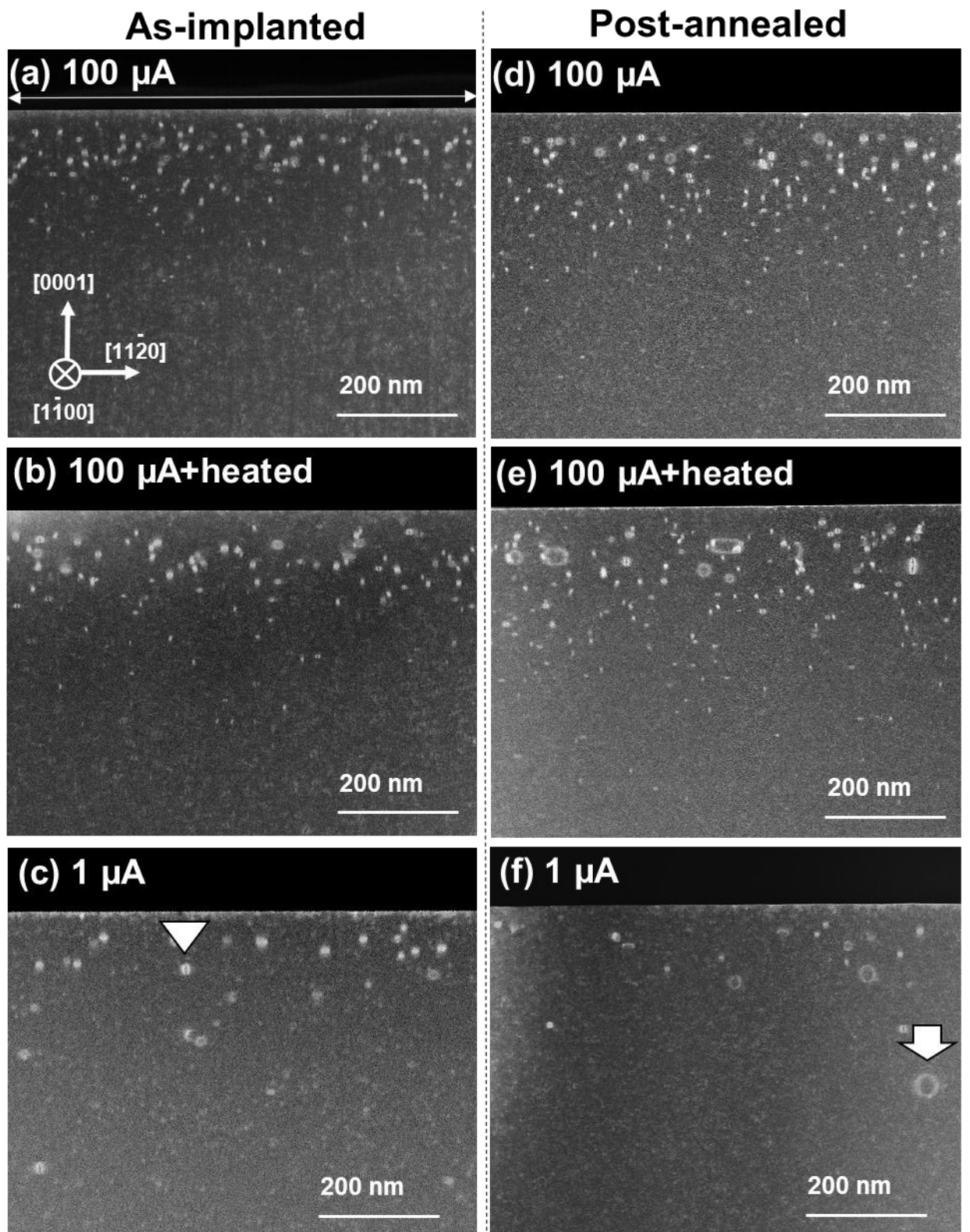


Fig. 3.

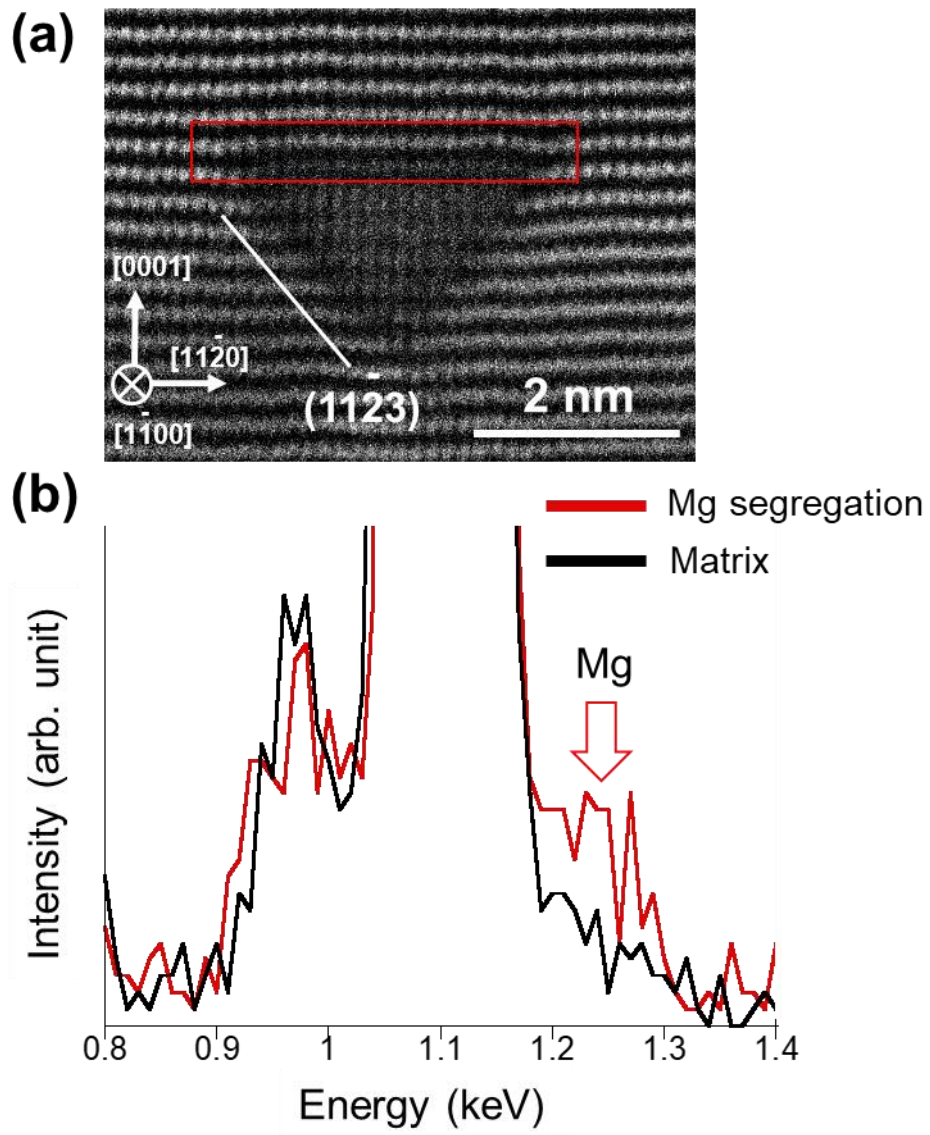


Fig. 4.