

1 **Crystalline structure of TiC ultra thin layers formed on highly**
2 **oriented pyrolytic graphite by chemical reaction from Ti/graphite**
3 **system**

4 Osamu Nakatsuka^{1,*}, Kenji Hisada^{1,†}, Satoshi Oida^{1,‡}, Akira Sakai², and Shigeaki Zaima^{1,3}

5
6 ¹ *Department of Crystalline Materials Science, Graduate School of Engineering, Nagoya*
7 *University, Nagoya 464-8603, Japan*

8 ² *Graduate School of Engineering Science, Osaka University, Toyonaka, Osaka 560-8531,*
9 *Japan*

10 ³ *Institute of Materials and Systems for Sustainability, Nagoya University, Nagoya*
11 *464-8603, Japan*

12 *E-mail: nakatuka@alice.xtal.nagoya-u.ac.jp

13

14

15 We have investigated the atomic-scale reaction between a Ti thin layer and highly oriented
16 pyrolytic graphite (HOPG) mainly by scanning tunneling microscopy. A deposited Ti layer
17 shows an epitaxial orientation structure on a HOPG substrate even in room-temperature
18 deposition, while the bonding between Ti and HOPG is very weak. The chemical reaction
19 between Ti and HOPG takes place, and epitaxial TiC domains on HOPG are formed for
20 annealing at above 600 °C. The TiC domains shows a smooth surface corresponding to the
21 TiC(111) plane after annealing at 700 °C. The formation of TiC(001) facets and significant
22 surface roughening of not only TiC but also HOPG substrates are observed for annealing at
23 as high as 900 °C.

† Present address: DENSO Corporation, Nukata-Gun, Aichi 444-0193, Japan

‡ Present address: IBM T. J. Watson Research Center, New York 10598, USA

24 **1. Introduction**

25 Carbon (C)-related low-dimensional structures such as graphene sheets and nanotubes have
26 attracted attention for nanoelectronic applications because of the high mobility of both
27 their electrons and holes and their characteristic electronic properties¹⁻⁷⁾. For the
28 application of low-dimensional C materials in nanoelectronic devices, it is necessary to
29 control the crystalline structure of not only the C material itself but also insulators, metals,
30 and interfaces of device structures on the atomic scale. Also, in the development of device
31 applications of C nanostructures, there are many challenges for metal/semiconductor
32 contact technologies to achieve low contact resistance, high thermal robustness, high
33 adhesive property, and high reliability⁸⁻¹³⁾.

34

35 Understanding the crystalline structure and chemical reaction at the interface between
36 a metal electrode layer and a C nanostructure is essential for attaining high performance
37 and reliability of the contact structure between metal and C-related materials. However, it
38 is difficult to analyze these properties in atomic-scale C nanostructures in detail because of
39 the limitation of the volume of low-dimensional C nanomaterials. On the other hand,
40 highly oriented pyrolytic graphite (HOPG) consists of high-quality stacked graphite
41 monolayers, and thus it is a good specimen for the fundamental study of metal/C
42 nanostructure contact. There are many previous studies of the behavior and crystalline
43 structure of some metals, such as Cu^{14,15)}, Ni^{15,16)}, Pt^{15,17)}, Ag^{18,19)}, and Co²⁰⁾, deposited on
44 graphite substrates. Also, scanning tunnel microscopy (STM) studies of the nanostructure
45 of metals on graphite substrates have been reported²¹⁻²⁵⁾. However, these previous studies
46 of the metal/graphite substrate system basically focused on the nanostructures of metals
47 deposited on substrates, and studies focusing on the metal/graphite contact structure and
48 interfacial reaction are limited. The relationship between the crystalline structure and the
49 chemical reaction between a metal and a graphite substrate has not yet been understood in

50 detail.

51

52 In this study, we examined the deposition of multi-monolayers of a titanium (Ti) thin
53 film directly on a HOPG substrate and analyzed the atomic structure of the sample surface
54 by STM and photoelectron spectroscopic techniques to investigate the crystalline structures
55 and chemical reaction of a Ti layer with graphite after annealing.

56

57 **2. Experimental methods**

58 HOPG was used as the substrate. Its surface was cleaned by peeling the surface graphite
59 layers off with an adhesive tape followed by annealing at 1000 °C in an ultrahigh-vacuum
60 (UHV) chamber whose base pressure was below 3×10^{-8} Pa. The surface of HOPG was
61 confirmed to be clean by in situ reflection high-energy electron diffraction (RHEED) and
62 STM. After cleaning the surface, Ti layers with thicknesses of 0.05–0.25 nm were
63 successively deposited using an electron beam evaporation system at a pressure below
64 1×10^{-7} Pa at room temperature (RT). Then, Ti/HOPG samples were annealed at
65 temperatures ranging from 600 to 1000 °C at a pressure below 1×10^{-6} Pa in the same UHV
66 chamber.

67

68 The crystalline structure and morphology of the sample surface were observed by in
69 situ RHEED and STM. The chemical bonding structure of surface layers in Ti/HOPG
70 samples was analyzed by ex situ X-ray photoelectron spectroscopy (XPS) using a Mg K α
71 X-ray source (photon energy; $h\nu = 1254$ eV).

72

73 **3. Results and discussion**

74 Figures 1(a)-1(d) show STM images of the Ti (thickness: 0.05 nm)/HOPG and Ti

75 (thickness: 0.25 nm)/HOPG samples. Figures 1(b) and 1(d) are enlarged images of Figs.
 76 1(a) and 1(c), respectively. In the Ti (thickness: 0.05 nm)/HOPG sample, many small
 77 particle-like bright spots are observed on the terrace of HOPG in Fig. 1(a). Also,
 78 protrusions along the step edge lines can be observed. These small particles and
 79 protrusions are considered to be agglomerated structure of deposited Ti. Note that we
 80 firstly observed the STM image shown in Fig. 1(b) and then observed the image in Fig.
 81 1(a). After observing the magnified STM image shown in Fig. 1(b), we found that there are
 82 no small particles in the region first observed in Fig. 1(a). This result suggests that the Ti
 83 particles were swept out and agglomerated at edge lines, and the as-deposited Ti atoms
 84 very weakly interacted with the HOPG surface, which is energetically stable.

85

86 For the deposited Ti thickness as thick as 0.25 nm, the agglomerated structure shows
 87 the characteristic threefold symmetry, as shown in Fig. 1(c). The in situ RHEED pattern
 88 obtained from this sample immediately after the Ti deposition is shown in Fig. 2. A streak
 89 with a ring pattern is clearly observed. The characteristic distances d_1 and d_2 were
 90 estimated to be 0.297 ± 0.004 and 0.513 ± 0.004 nm, respectively, from the diffraction pattern
 91 shown in Fig. 2. These values correspond to the lattice planes of hexagonal Ti(1000) ($d =$
 92 0.296 nm) and Ti($11\bar{2}0$) ($d = 0.513$ nm)²⁶). These values can be associated with the
 93 hexagonal structure of Ti domains and indicate the epitaxial growth of Ti thin layers on
 94 HOPG. The epitaxial relationship between Ti and the HOPG substrate can be determined
 95 to be Ti(0001)//HOPG(0001) and Ti[$11\bar{2}0$]/HOPG[$11\bar{2}0$] from the RHEED pattern.

96

97 We deduced that this characteristic structure is related to the Van der Waals
 98 epitaxy^{27,28}) of a Ti thin layer on the HOPG substrate. Van der Waals epitaxy was also
 99 previously reported in the growth of thin films such as GaSe on MoS₂²⁹) and CrPt₃ on
 100 WSe₂³⁰). It is known that this characteristic growth is observed in the case of thin-film

101 growth on an energetically inactive surface of a substrate. In addition, the condition of the
102 long-range periodic structure with a large lattice mismatch between the grown film and the
103 substrate is required for this epitaxy. In fact, it is well known that an HOPG surface is
104 relatively energetically stable and the lattice mismatch between Ti and HOPG is as large as
105 20%, considering the epitaxial relationship mentioned above. This means that the bonding
106 of the Ti layer with the HOPG substrate just by Van der Waals force is very weak, which is
107 consistent with the sweeping out of the agglomerated Ti particles at step edges just by tip
108 scanning in the very thin Ti layer deposited on HOPG, which is observed in the STM
109 image shown in Fig. 1(a).

110

111 Next, we examined the in situ annealing of Ti/HOPG samples under UHV condition.
112 Figures 3(a)-3(c) show STM images of Ti (thickness: 0.25 nm)/HOPG samples after in situ
113 annealing at 700, 900, and 1000 °C, respectively. The surface profiles at the position
114 indicated in these images by arrows are also shown below the STM images. After
115 annealing at 700 °C, many hexagonal domains with atomically flat and uniform surfaces
116 are formed on the HOPG terrace. On the other hand, after annealing above 900 °C, the
117 symmetry of the domains significantly degrades as they agglomerate with significant
118 roughening of the domain surface. In addition, after annealing at 1000 °C, significant
119 erosion of the HOPG substrate around the domains is observed as a dip at the edge region
120 of the domains as indicated by red arrows in the surface profile shown in Fig. 3(c).

121

122 Figures 4(a)-4(c) show the RHEED patterns of the Ti (thickness: 0.25 nm)/HOPG
123 samples after in situ annealing at 700, 900, and 1000 °C, respectively. After 700 °C
124 annealing in UHV, the characteristic distances d_3 and d_4 shown in Fig. 3(a) were estimated
125 to be 0.306 ± 0.004 and 0.531 ± 0.004 nm, respectively. These values correspond to the
126 lattice planes of Ti carbide, TiC(100) ($d = 0.306$ nm), and TiC(111) ($d = 0.530$ nm)²⁶. The

127 STM and RHEED results indicate the epitaxial growth of a TiC layer on HOPG due to the
128 chemical reaction between Ti and HOPG. The epitaxial relationship between TiC and
129 HOPG is determined to be $\text{TiC}(111)//\text{HOPG}(0001)$ and $\text{TiC}[1\bar{1}0]//\text{HOPG}[11\bar{2}0]$. Moreover,
130 after 900 °C annealing, the streak in the diffraction pattern becomes spotty, indicating the
131 degradation of the TiC layer with the formation of a 3D island-like morphology, which is
132 consistent with the observation from the STM image shown in Fig. 3(b).

133

134 Here, we deduce the morphological transition of the TiC layer on HOPG depending
135 on the annealing condition. After annealing below 700 °C, the TiC layer shows a smooth
136 surface corresponding to the $\text{TiC}(0001)$ plane (c-plane). Ito et al. previously reported the
137 STM observation of the epitaxial growth of a graphite monolayer on a $\text{TiC}(111)$ surface³¹.
138 They found that the graphite monolayer grows showing that the azimuthal relationship
139 with TiC is $\text{C}[11\bar{2}0]//\text{TiC}[1\bar{1}0]$, which is consistent with the observed crystalline structure
140 in this study. They note that, with this relationship, one-seventh of the reciprocal space unit
141 vector of $\text{TiC}[1\bar{1}0]$ corresponds to one-fifteenth of that of the graphite along $[11\bar{2}0]$, and
142 one-fourteenth of the reciprocal lattice of $\text{TiC}[11\bar{2}]$ corresponds to one-tenth of that of
143 graphite along $[10\bar{1}0]$. We deduce that this long-range matching with the superlattice
144 would lead to the formation of a TiC epitaxial layer on HOPG at a temperature of 700 °C.

145

146 On the other hand, after annealing above 900 °C, the agglomeration of TiC is
147 observed and the domains show a 3D-like pyramidal shape surrounded by $\text{TiC}\{1000\}$
148 a-plane facets. This result is taken into account from the viewpoint of the surface energy of
149 TiC. Dudiy et al. reported that the surface energies of $\text{TiC}(100)$, $\text{TiC}(110)$, and $\text{TiC}(111)$
150 with the relaxed structure are theoretically calculated 1.73, 3.78, and 5.63 J/m², and that the
151 TiC a-plane is energetically most stable among these planes³². Therefore, it is considered
152 that flat TiC domains with the c-plane are formed in the initial stage of the reaction

153 between Ti and HOPG by low-temperature annealing, although TiC domains agglomerate
154 with a change to a pyramidal shape to preferably form a-plane facets by high-temperature
155 annealing.

156

157 We also characterized the chemical reaction products after annealing the Ti/HOPG
158 samples by ex situ XPS measurement. Figure 5 shows the photoelectron spectra of the Ti
159 2p core level obtained from the Ti (thickness: 0.25 nm)/HOPG samples as-deposited and
160 after annealing at temperatures ranging from 600 to 1000 °C. In the case of the
161 as-deposited sample, a Ti oxide component at binding energies of 464.5 and 458.8 eV for
162 Ti 2p_{1/2} and Ti 2p_{3/2}, respectively, is dominant³³). In the spectrum for the as-deposited
163 sample, the component related to the metal Ti, which was previously reported to have a
164 binding energy of 453.8 eV for the Ti 2p_{3/2} core level³³), is hardly observed. This is
165 considered to be due to the oxidation of a Ti metal layer during the sample transport in
166 atmosphere from the deposition chamber to the ex situ XPS measurement system. This is
167 because a Ti thin layer is a chemically active material and is easily oxidized even at RT.

168

169 On the other hand, the intensities of peaks at 460.7 and 454.7 eV for Ti 2p_{1/2} and Ti
170 2p_{3/2}, respectively, significantly increase with annealing temperature, which can be
171 considered to be the formation of TiC³⁴) through a chemical reaction between the Ti layer
172 and HOPG takes place in high-temperature annealing. This result is consistent with the
173 STM and RHEED observations mentioned above. Interestingly, note that we can identify
174 the component related to TiC from the as-deposited Ti layer at RT on HOPG; although the
175 intensity is lower than that of the annealed sample. This result suggests that TiC could be
176 formed only by depositing Ti on graphite at RT. However, the component of TiC is limited
177 in the as-grown sample, and we cannot determine in detail where the reaction preferentially
178 occurs, for example, at the step in the HOPG surface. Considering the STM observation

179 mentioned above, we guess that the bonding structure at the interface between Ti and
180 HOPG is dominated by Van der Waals force.

181

182 We summarize the crystalline structure and reaction between Ti and HOPG in Fig. 6.
183 In the as-deposited Ti sample, the deposited Ti forms domains and epitaxially grows on
184 HOPG, despite the fact than the bonding between Ti and HOPG by Van der Waals epitaxy
185 is very weak [Fig. 6(a)]. After annealing at a low temperature in the range of 600–700 °C,
186 the formation of epitaxial TiC domains through the reaction between Ti and HOPG takes
187 place, and the surface of TiC becomes uniform and flat [Fig. 6(b)]. With the annealing
188 temperature increased to 900 °C, TiC domains agglomerate with the formation of {111}
189 facets [Fig. 6(c)], and the surface roughening of HOPG becomes significant especially
190 around TiC domains after annealing at temperatures as high as 1000 °C [Fig. 6(d)].

191

192 **4. Conclusions**

193 We investigated the atomically crystalline structure of a Ti thin layer deposited on HOPG
194 by in situ RHEED, STM, and ex situ XPS. We observed the epitaxial growth of Ti
195 immediately after its deposition on HOPG at RT. After annealing above 600 °C, the
196 formation of TiC proceeds with increasing annealing temperature, and STM and RHEED
197 observations revealed the epitaxial growth of TiC with characteristic domains. The
198 roughening of the surface of not only TiC but also HOPG substrates for annealing at
199 temperatures as high as 900 °C. Considering these results, we have to control the
200 atomic-scale interfacial reaction between the metal electrode and the low-dimensional
201 graphite structure to form low-resistance contacts with high reliability for C
202 nanoelectronics.

203 **References**

- 204 ¹⁾ S. Frank, P. Poncharal, Z. L. Wang, and W. A. de Heer, *Science* **280**, 1744 (1998).
205 ²⁾ M. Shim, A. Javey, N. W. S. Kam, and H. Dai, *J. Am. Chem. Soc.* **123**, 11512 (2001).
206 ³⁾ T. Dürkop, S. A. Getty, E. Cobas, and M. S. Fuhrer, *Nano Lett.* **4**, 35(2004).
207 ⁴⁾ S. V. Morozov, K. S. Novoselov, F. Schedin, D. Jiang, A. A. Firsov, and A. K. Geim,
208 *Phys. Rev. B* **72**, 201401(2005).
209 ⁵⁾ K. Geim and K. S. Novoselov, *Nat. Mater.* **6**, 183 (2007).
210 ⁶⁾ S. J. Kang, C. Kocabas, T. Ozel, M. Shim, N. Pimparkar, M. A. Alam, S. V. Rotkin,
211 and J. A. Rogers, *Nat. Nanotechnol.* **2**, 230 (2007).
212 ⁷⁾ F. Schwierz, *Nat. Nanotechnol.* **5**, 487 (2010).
213 ⁸⁾ E. J. H. Lee, K. Balasubramanian, R. T. Weitz, M. Burghard, and K. Kern, *Nat.*
214 *Nanotechnol.* **3**, 486, (2008).
215 ⁹⁾ P. Blake, R. Yang, S. V. Morozov, F. Schedin, L. A. Ponomarenko, A. A. Zhukov, R. R.
216 Nair, I. V. Grigorieva, K. S. Novoselov, and A. K. Geim, *Solid State Commun.* **149**,
217 1068 (2009).
218 ¹⁰⁾ S. Russo, M. F. Craciun, M. Yamamoto, A. F. Morpurgo, and S. Tarucha, *Physica E* **42**,
219 677 (2010).
220 ¹¹⁾ K. Nagashio, T. Nishimura, K. Kita, and A. Toriumi, *Appl. Phys. Lett.* **97**, 143514
221 (2010).
222 ¹²⁾ Y. Matsuda, W.-Q. Deng, and W. A. Goddard III, *J. Phys. Chem. C* **114**, 17845 (2010).
223 ¹³⁾ F. Xia, V. Perebeinos, Y.-M. Lin, Y. Wu and P. Avouris, *Nat. Nanotechnol.* **6**, 179
224 (2011).
225 ¹⁴⁾ P. B. Abel, A. L. Korenyi-Both, F. S. Honey, and S. V. Pepper *J. Mater. Res.* **9**, 617
226 (1994).
227 ¹⁵⁾ P. Marcus and C. Hinnen, *Surf. Sci.* **392**, 134 (1997).
228 ¹⁶⁾ M. F. De Riccardis, D. Carbone, V. Martina, M. Re, B. Bozzini, and L. D'Urzo, *Appl.*
229 *Surf. Sci.* **255**, 4309 (2009).
230 ¹⁷⁾ D.-Q. Yang, G.-X. Zhang, and E. Sacher, M. José-Yacamán, and N. Elizondo, *J. Phys.*
231 *Chem. B* **110**, 8348 (2006).
232 ¹⁸⁾ M.-H. Schaffnery, J.-F. Jeanneret, F. Patthey, and W.-D. Schneider, *J. Phys. D* **31**, 3177
233 (1998).
234 ¹⁹⁾ H. Zhang, Q. Fu, Y. Yao, Z. Zhang, T. Ma, D. Tan, and X. Bao, *Langmuir* **24**, 10874
235 (2008).
236 ²⁰⁾ S. W. Poon, J. S. Pan, and E. S. Tok, *Phys. Chem. Chem. Phys.* **8**, 3326 (2006).

- 237 ²¹⁾ E. Ganz, K. Sattler, and J. Clarkes, *Surf. Sci.* **219**, 33 (1989).
238 ²²⁾ V. Maurice and P. Marcus, *Surf. Sci.* **275**, 65 (1992).
239 ²³⁾ T. Endo, T. Sunada, T. Sumomogi, and H. Maeta, *Mater. Charact.* **48**, 159 (2002).
240 ²⁴⁾ I. Lopez-Salido, D. C. Lim, and Y. D. Kim, *Surf. Sci.* **588**, 6 (2005).
241 ²⁵⁾ I. Lopez-Salido, D. C. Lim, R. Dietsche, N. Bertram, and Y. D. Kim, *J. Phys. Chem. B*
242 **110**, 1128 (2006).
243 ²⁶⁾ Inorganic Material Database (AtomWork) [<http://crystdb.nims.go.jp/>].
244 ²⁷⁾ A. Koma, K. Sunouchi, and T. Miyajima, *Microelectron. Eng.* **2**, 129 (1984).
245 ²⁸⁾ A. Koma, *Thin Solid Films* **216**, 72 (1992).
246 ²⁹⁾ K. Ueno, K. Sasaki, N. Takeda, K. Saiki, and A. Koma, *Appl. Phys. Lett.* **70**, 1104
247 (1997).
248 ³⁰⁾ M. Maret, B. Gilles, J.P. Simon, M. Verdier, I. Guhr, B. Riedlinger, M. Albrecht, and G.
249 Schatz, *J. Cryst. Growth* **275**, 2289 (2005).
250 ³¹⁾ H. Ito, H. Ichinose, C. Oshima, and T. Ichinokawa, *Surf. Sci. Lett.* **254**, L437 (1991).
251 ³²⁾ S. V. Dudiy and B. I. Lundqvist, *Phys. Rev. B* **69**, 125421 (2004).
252 ³³⁾ L. Ramqvist, K. Hamrin, G. Johansson, A. Fahlman, and C. Nordling, *J. Phys. Chem.*
253 *Solids* **30**, 1835 (1969).
254 ³⁴⁾ H. Ihara, Y. Kumashiro, A. Itoh, and K. Maeda, *Jpn. J. Appl. Phys.* **12**, 1462 (1973).
255

256 **Figure Captions**

257

258 Fig. 1. (Color online) The STM images of the (a) Ti (thickness: 0.05 nm)/HOPG sample
259 and (b) magnified one, and (c) Ti (thickness: 0.25 nm)/HOPG sample and (d) magnified
260 one.

261

262 Fig. 2. In situ RHEED pattern of the Ti (thickness: 0.05 nm)/HOPG sample immediately
263 after Ti deposition.

264

265 Fig. 3. (Color online) STM images of Ti (thickness: 0.25 nm)/HOPG samples after in situ
266 annealing at (a) 700, (b) 900, and (c) 1000 °C.

267

268 Fig. 4. RHEED patterns of Ti (thickness: 0.25 nm)/HOPG samples after in situ annealing at
269 (a) 700, (b) 900, and (c) 1000 °C.

270

271 Fig. 5. (Color online) Photoelectron spectra of the Ti 2p core-level obtained from Ti
272 (thickness: 0.25 nm)/HOPG samples as-deposited and after annealing at 600, 700, 800, 900,
273 and 1000 °C.

274

275 Fig. 6. (Color online) Schematic diagrams of crystalline structure and reaction between Ti
276 and HOPG observed in this study. (a) As-deposited, and after annealing at (b) 600–700, (c)
277 900, and (d) 1000 °C.

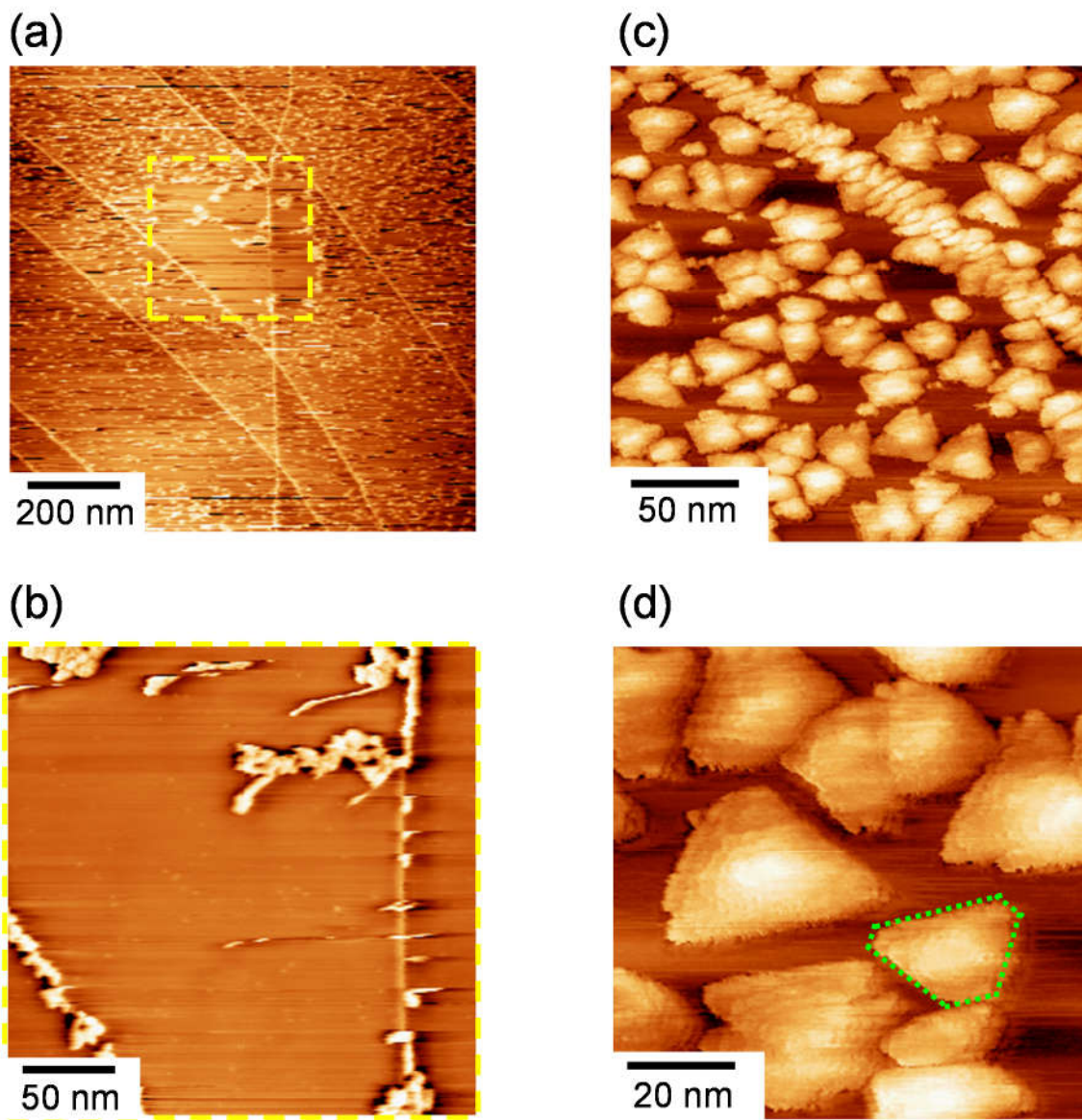


Fig. 1

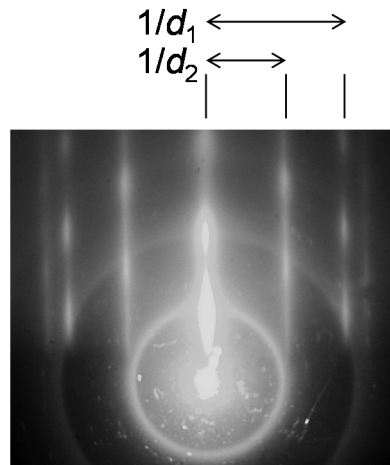


Fig. 2

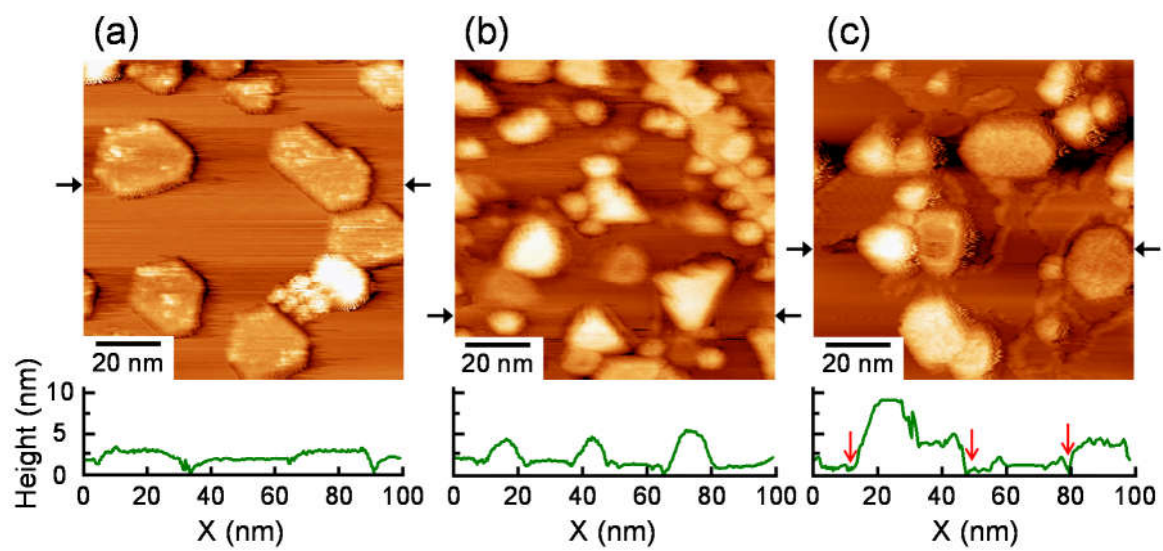


Fig. 3

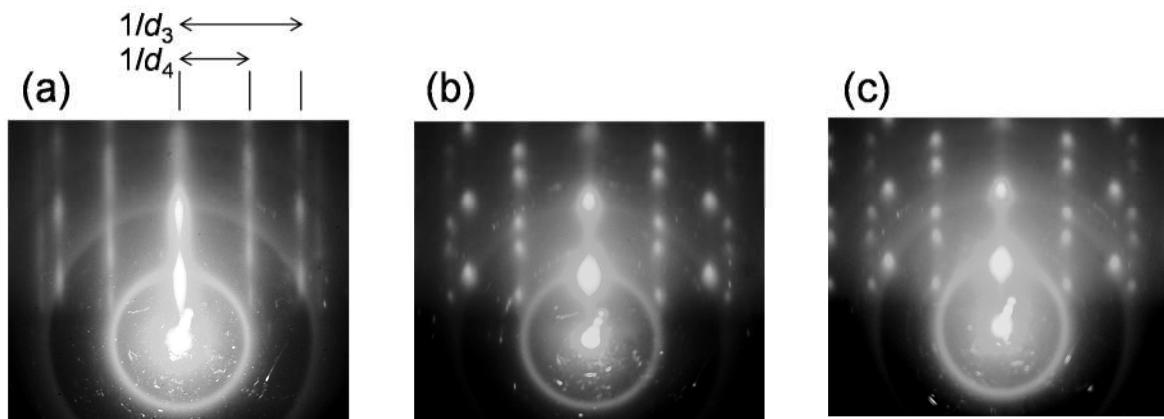


Fig. 4

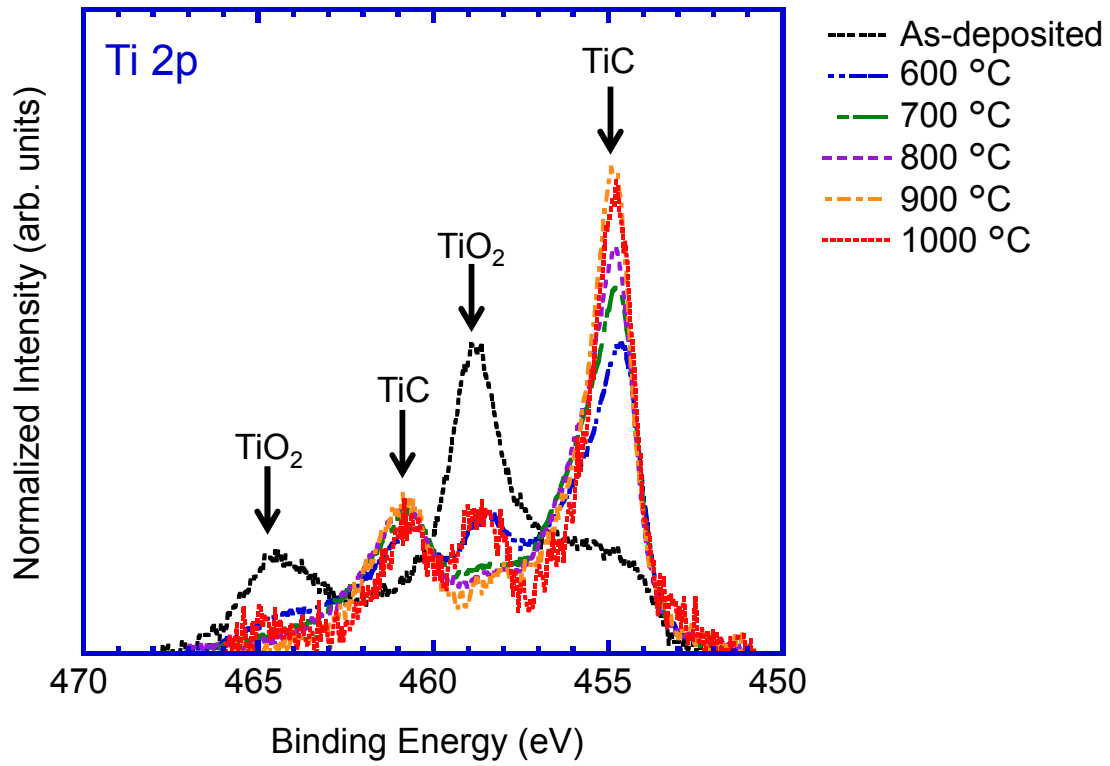


Fig. 5

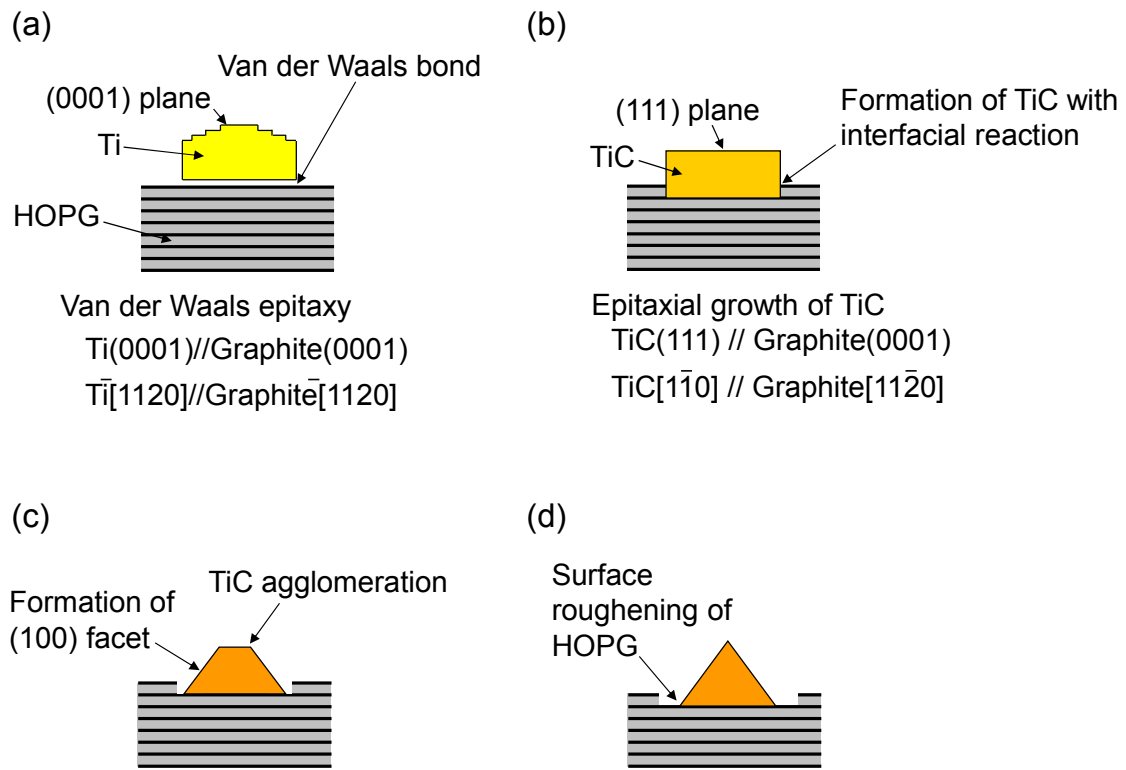


Fig. 6