# Optimum Breakdown Charge for Maximum Dielectric Strength in Spark Conditioning in Vacuum under a Non-uniform Electric Field

H. Kojima, T. Takahashi, N. Hayakawa

Nagoya University Furo-cho, Chikusa-ku Nagoya 464-8603, Japan

#### K. Hasegawa and M. Sakaki

MEIDENSHA CORPORATION 515 Kaminakamizo, Higashimakado Numazu 410-8588, Japan

#### ABSTRACT

Spark conditioning is an effective method of increasing the dielectric strength in vacuum. In this paper, we determined the breakdown (BD) charge dependence for spark conditioning in vacuum under various non-uniform electric fields and with different anode materials. We discuss the BD mechanism in terms of BD arising from anode heating or cathode heating. Experimental results revealed that for anode heating BD, the longer the gap distance, the larger the BD charge needed to obtain maximum dielectric strength. In contrast, for cathode heating BD, the BD charge for maximum dielectric strength was smaller than that for anode heating BD. From these results, we estimated the optimal BD charge to maximize the 50% BD electric field strength, i.e. the dielectric strength, for spark conditioning in vacuum.

Index Terms — Vacuum discharge, conditioning, non-uniform electric field, vacuum circuit breakers, breakdown.

## **1 INTRODUCTION**

VACUUM interrupters (VIs) and vacuum circuit breakers (VCBs) are expected to be used for high-voltage applications because of their high insulation performance, good current interruption performance, and environmental friendliness [1-4]. Spark conditioning is used to improve the dielectric strength in vacuum [5]. The conditioning effect has a maximum value with respect to discharge current, and the breakdown (BD) charge at which the maximum conditioning effect is obtained depends on the cathode material [6].

We have previously investigated the effects of the BD charge on spark conditioning under a non-uniform electric field, based on the mechanism of BD initiation [6]. It is expected that the larger BD charge leads to getting the higher conditioning effect, but the excess BD charge causes the damaging effect of the electrode. Therefore, the maximum dielectric strength by spark conditioning is achieved by the balance between the conditioning and the damaging effects on the electrode surface, and these effects depend on the BD mechanism. Although various theories of BD in vacuum have been proposed [7, 8], assuming that clumps on the electrodes are sufficiently removed in the final stage of spark conditioning, cathode heating BD (CH-BD) or anode heating BD (AH-BD) will mainly occur. By applying a high voltage to the electrode system during spark conditioning, electron emission from the cathode heats the electrode and vaporizes the electrode material. In CH-BD, cathode heating leads to BD [7]; in contrast, in AH-BD, anode heating leads to BD [8]. The possible BD mechanism is dominated by the conditions and configuration of the electrode system, and the difference in the BD mechanism mediates the effect of the conditioning on the magnitude of BD charge. Therefore, to increase the dielectric strength in vacuum, we must clarify the optimum BD charge for spark conditioning with various gap lengths and electrode materials, based on the BD mechanism.

In this paper, we focus on the BD charge dependence for spark conditioning in vacuum for various non-uniform electric fields, gap lengths, and anode materials. Then, we discuss the optimal BD charge for the spark conditioning in vacuum.

Manuscript received on X Month 2005, in final form XX Month 2005.

## 2 EXPERIMENTAL SETUP AND PROCEDURE

Figure 1 shows the experimental setup. We repeatedly applied a negative lightning impulse voltage generated by an impulse generator (I.G.) to the rod-plane electrode system in a vacuum chamber. We used the up-down method for the repetitive impulse voltage application. In the up-down method, the initial charged voltage of the I.G. was 30 kV, and the charged voltage  $\Delta V$  was increased in 10 kV increments. The applied voltage and the discharge current were observed with a voltage divider and a current transformer (CT), respectively, and were recorded with a digital oscilloscope. The vacuum pressure in the chamber was kept around  $10^{-6}$  Pa.

Figure 2 shows the configuration of the rod-plane electrode system. The gap distance *d* was from 1 to 30 mm and the tip radius *r* of the rod electrode was 1, 2, or 5 mm. The rod electrode (cathode) was made with oxygen-free high thermal conductivity copper (Cu), and the grounded plane electrode (anode) was made with Cu or stainless steel (SS). We changed the discharge current at BD by varying the limiting resistance  $R_L$  to 30  $\Omega$ , 60  $\Omega$ , 2 k $\Omega$ , 11 k $\Omega$ , and 20 k $\Omega$ . Figure 3 is typical waveforms of applied voltage and discharge current when BD occurred during the conditioning process. The discharge current was suppressed by higher  $R_L$ . Most of BDs during conditioning process occurred at the peak of applied voltage.

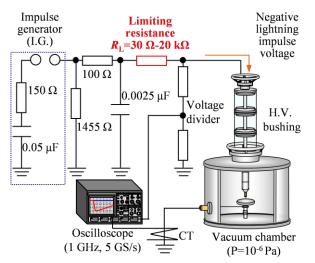


Figure 1. Schematic of the experimental setup.

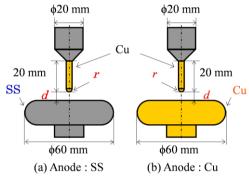
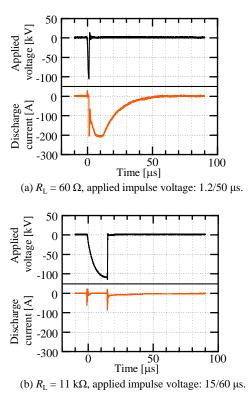


Figure 2. Schematic of the electrode configuration.



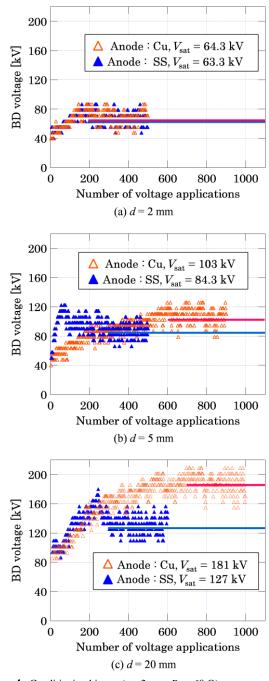
**Figure 3.** Typical waveforms of applied voltage and discharge current when BD occurred during conditioning process. Anode: SS, d = 5 mm. The charged voltage of I.G. was 130 kV.

## 3 RESULTS AND DISCUSSION

## 3.1 CONDITIONING HISTORY

Figure 4 shows the conditioning history for gap distances of d = 2, 5, and 20 mm. In all cases, r = 2 mm and  $R_{\rm L} = 60 \Omega$ . The 50% BD voltage ( $V_{\rm sat}$ ) after conditioning saturation is shown by horizontal lines. When the maximum BD voltage did not change over 300 voltage applications, we considered the conditioning saturated. This saturation was also confirmed by stabilization of Weibull plot of BD voltage [9]. For d = 2 mm (Figure 4a),  $V_{\rm sat}$  was similar regardless of the anode material. However, for d = 5 and 20 mm (Figures 4b and 4c), BD voltage decreased during the conditioning process for the SS anode. This implies that the electrodes were damaged during the conditioning.  $V_{\rm sat}$  was lower for the SS anode than for the Cu anode.

Figure 5 shows the development of the BD spot on the rod cathode surface for the Cu cathode and SS anode, observed by the digital camera at d = 5 mm (conditions shown in Figure 4b). On the Cu anode (Figure 5a), the BD spot started at the tip of the electrode and moved to the wider region of the electrode with a lower field as the number of voltage applications increased, which is consistent with [10], i.e. the development of conditioning was dominated by the conditioning degree at the tip of electrode. For the SS anode, the BD spot tended to be concentrated at the tip of the cathode, which implies that the cathode tip was heavily damaged.

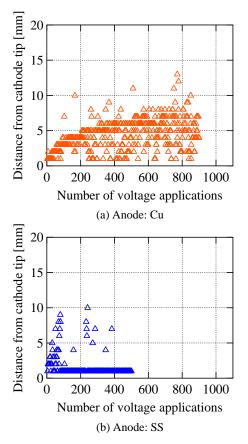


**Figure 4.** Conditioning history (r = 2 mm,  $R_L = 60 \Omega$ ).

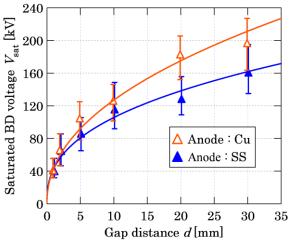
#### 3.2 GAP DISTANCE DEPENDENCE OF CONDITIONING CHARACTERISTICS

Figure 6 shows the relation between  $V_{\text{sat}}$  and d. In this figure, the error bars indicate the maximum and minimum values of BD voltage after the conditioning saturation. For short gap distances ( $d \le 2$  mm),  $V_{\text{sat}}$  was similar for both anode materials, whereas for long gap distances (d = 5-30 mm),  $V_{\text{sat}}$  for the Cu anode was higher than that for the SS anode.

Figure 7 shows microscope images of the surface at the tip of the Cu cathode before and after conditioning. Before conditioning, machining traces were observable as several straight trails. After conditioning, for short gap distances,



**Figure 5.** Development of the BD spot on the cathode surface (r = 2 mm,  $R_L = 60 \Omega$ , d = 5 mm). The resolution of the BD spot is 1 mm. On the vertical axis, 0 mm is the cathode tip, and 20 mm is the cathode base (see Figure 2).



**Figure 6.** Relationship between  $V_{\text{sat}}$  and  $d (r = 2 \text{ mm}, R_{\text{L}} = 60 \Omega)$ .

severe melting was visible and the machining traces had disappeared, regardless of the anode material. In contrast, for long gap distances, the cathode for the Cu anode melted less than the cathode for the SS anode.

For short gap distances, CH-BD tends to occur because a high electric field appears at a low voltage. CH-BD does not depend on the anode, so it is considered that similar conditioning characteristics were obtained for different anode

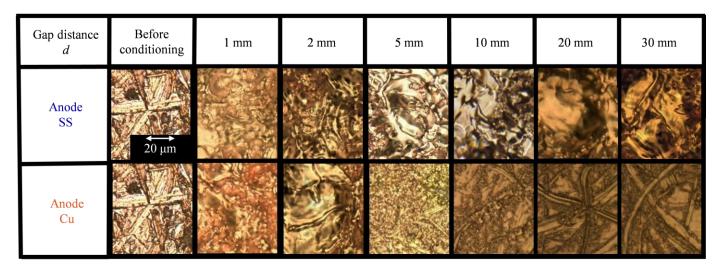


Figure 7. Microscope images of the surface at the tip of the Cu cathode before and after conditioning saturation (r = 2 mm,  $R_L = 60 \Omega$ ).

materials for short gap distances. In addition, when CH-BD mainly occurred, the BD charge was large [6] and the cathode was heavily damaged. In contrast, for long gap distances with the Cu anode, AH-BD tends to occur because a high voltage is needed to form a high electric field; therefore, the electron accelerating voltage between the electrodes becomes large. During AH-BD, it is considered that the cathode temperature immediately prior to BD is lower than that during CH-BD, and the damage to the cathode surface is suppressed. However, for the SS anode, it is thought that CH-BD tended to occur even for long gap distances. This agrees with the heavy damage of the cathode for the SS anode (Figure 7). As mentioned in Section 3.1, this agrees with the decrease in BD voltage during the conditioning (Figure 4) and the concentration of the BD spot at the cathode tip (Figure 5). As a result,  $V_{sat}$  for the Cu anode was higher than that for the SS anode where CH-BD mainly occurred (Figure 6).

#### 3.3 EFFECT OF GAP DISTANCE AND ANODE MATERIAL ON BREAKDOWN CHARGE DEPENDENCE

Figure 8 shows the relationship between BD field strength  $E_{\text{sat}}$  at the cathode tip and charge  $Q_{\text{BD}}$  per single BD after conditioning saturation.  $E_{\text{sat}}$  was calculated from BD voltage by using the charge simulation method, and  $Q_{\text{BD}}$  is the time integration value of the discharge current waveform in a BD after conditioning saturation as shown in Figure 3. The symbols represent the average values for each combination of r and d. The error bars represent the minimum and maximum values of  $E_{\text{sat}}$  and  $Q_{\text{BD}}$  after the conditioning saturation.

For the SS anode,  $E_{\text{sat}}$  depended on  $Q_{\text{BD}}$  in spite of *d* and *r* (Figure 8a) and CH-BD mainly occurred, as discussed in Section 3.2. CH-BD occurred after the cathode temperature reached the melting point, which depended on the electric field at the cathode surface regardless of the applied voltage. The cathode temperature just before BD will have an almost constant value, namely, the melting point [6]. Thus, the amount of molten cathode material or damaging effect of cathode depends only on the discharge magnitude  $Q_{\text{BD}}$ .

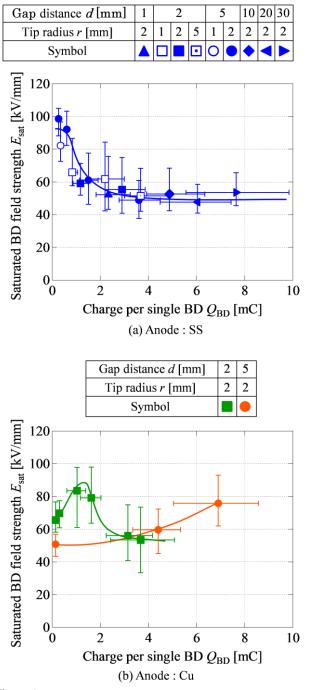
For the Cu anode, the relation between  $E_{\text{sat}}$  and  $Q_{\text{BD}}$  depended on *d* (Figure 8b). For d = 5 mm, AH-BD tended to occur, as discussed in Section 3.2. AH-BD can occur before the cathode temperature gets up to the melting point. In other words, for AH-BD, less of the cathode melted than for CH-BD. The damage to the cathode surface was thereby suppressed, improving the conditioning effect, and  $V_{\text{sat}}$  increased with  $Q_{\text{BD}}$ .

For d = 2 mm, CH-BD mainly occurred as discussed in Section 3.2, similar to the SS anode. In Figure 8b, the damage to the cathode surface was dominant at  $Q_{BD} \ge 3$  mC. At lower  $Q_{BD}$  values ( $Q_{BD} \le 3$  mC), the conditioning progressed more than at higher  $Q_{BD}$  values, and  $E_{sat}$  increased. This higher  $E_{sat}$ could increase the collision energy of the electrons with the anode at BD, and the BD mechanism would change from CH-BD to AH-BD, resulting in the appearance of the maximum  $E_{sat}$  in Figure 8b.

Figure 9 shows the dependence of the BD characteristics on the gap distance for AH-BD. For longer gap distances, the voltage accelerating electrons between electrodes became large. Thus, a smaller field emission electron current is sufficient for BD. This means that BD occurs in a lower electric field at longer gap distances. Because the cathode heating before BD is caused by the field emission electron current, the amount of the molten cathode depends on the electric field on the cathode surface. Thus, for longer gap distances, the cathode temperature immediately prior to BD is lower than that for shorter gap distances. Therefore, for AH-BD, the longer the gap distance, the larger the  $Q_{\rm BD}$  needed to obtain the maximum  $E_{\rm sat}$ .

#### 3.4 OPTIMUM BREAKDOWN CHARGE FOR SPARK CONDITIONING

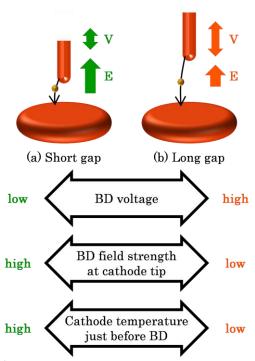
For AH-BD, at longer gap distances,  $Q_{BD}$  to obtain the maximum  $E_{sat}$  increased, whereas at shorter gap distances, CH-BD mainly occurred. For CH-BD, the cathode temperature just before BD was constant at the melting point. However, for AH-BD, the cathode temperature did not reach the melting point before BD. Thus, the cathode temperature just before BD for CH-BD was higher than that for AH-BD. This means



**Figure 8.** Relationship between  $E_{\text{sat}}$  and  $Q_{\text{BD}}$ . The error bars represent the minimum and maximum values after the conditioning saturation.

that for CH-BD the same amount of molten cathode, (and thus the same conditioning effect) could be obtained with a smaller  $Q_{\rm BD}$  than for AH-BD (Figure 8). Therefore,  $E_{\rm sat}$  reached its maximum value at a smaller  $Q_{\rm BD}$  value.

For CH-BD, the cathode temperature just before BD did not depend on the cathode shape (Figure 8a). Therefore, because the amount of molten cathode, which indicates the conditioning effect and the damage to the cathode, depended on  $Q_{\rm BD}$  only, and the  $Q_{\rm BD}$  value required to obtain the maximum  $E_{\rm sat}$  value did not depend on the cathode shape. We



**Figure 9.** Gap distance dependence of BD characteristics in the case of BD based on anode heating (AH-BD).

expect that these conditioning characteristics should be able to identify the optimal  $Q_{BD}$  for the maximum  $E_{sat}$  in any electrode configurations.

A previous study showed that the damaging effect during the spark conditioning decreases the minimum BD voltage rather than the 50% BD voltage [9]. In the conditioning of VIs, the minimum BD voltage is as important to reliability as 50% BD voltage is [11]. In this paper, we focused on maximizing  $E_{sat}$  to optimize conditioning. However, based on the BD mechanisms, the damage is also dominated by  $Q_{BD}$ . Therefore, we expect that our method could also be used to maximize and optimize the minimum BD voltage.

#### **4 CONCLUSION**

We focused on the dependence of conditioning characteristics in vacuum on the gap distance and BD charge for various anode materials. Our results for the optimal BD charge for spark conditioning in vacuum to maximize the 50% BD electric field  $E_{sat}$  can be summarized as follows:

(1) For short gap distances, CH-BD occurs regardless of the anode materials. Therefore, damage to the cathode surface becomes dominant.

(2) For long gap distances and anode materials with a high melting point, CH-BD mainly occurs. In contrast, for anode materials with a low melting point, AH-BD mainly occurs, and damage to the cathode surface is suppressed.

(3) When AH-BD mainly occurs, the longer the gap distance is, the larger the optimal  $Q_{BD}$  to obtain the maximum BD field strength.

(4) When CH-BD occurs, the optimal BD charge to obtain the maximum BD field strength does not depend on the electrode shape.

## REFERENCES

- H. Fink and R. Renz, "Future trends in vacuum technology applications", 20th Int. Symp. Discharge and Elect. Ins. in Vacuum, pp. 25-29, 2002.
- [2] H. Saitoh, H. Ichikawa, A. Nishijima, Y. Matsui, M. Sakaki, A. Honma, and H. Okubo, "Research and development on 145kV/40kA one break vacuum circuit breaker", IEEE/PES Transmission and Distribution Conf. and Exhibition: Asia Pac., Vol. 2, pp. 1465-1468, 2002.
- [3] M. Homma, M. Sakaki, E. Kaneko, and S. Yanabu, "History of vacuum circuit breakers and recent developments in Japan", IEEE Trans. Dielectr. Electr. Insul., Vol. 13, No. 1, pp. 85-92, 2006.
- [4] H. Okubo, H. Kojima, K. Kato, N. Hayakawa, and M. Hanai, "Advanced electrical insulation techniques for higher voltage vacuum interrupters", Cigre B3/D1 Colloquium, No. 218, 2013.
- [5] P. G. Slade, *The Vacuum Interrupter: Theory, Design, and Application*, CRC Press, London, Ch. 1, 2008.
- [6] H. Kojima, T. Takahashi, N. Hayakawa, K. Hasegawa, H. Saito, and M. Sakaki, "Dependence of spark conditioning on breakdown charge and electrode material under a non-uniform electric field in vacuum" IEEE Trans. Dielectr. Electr. Insul., Vol. 23, No. 5, pp. 3224-3230, 2016.
- [7] P. A. Chatterton, "A theoretical study of field emission initiated vacuum breakdown", Proc. Phys. Soc., Vol. 88, pp. 231-245, 1966.
- [8] W. S. Boyle, P. Kisliuk, and L. H. Germer, "Electrical breakdown in high vacuum", J. Appl. Phys., Vol. 26, pp. 720-725, 1955.
- [9] H. Kojima, Y. Otake, T. Takahashi, N. Hayakawa, K. Hasegawa, H. Saito, and M. Sakaki, "Analysis of conditioning and damaging process in vacuum based on breakdown probability distribution", IEEE Trans. Dielectr. Electr. Insul., Vol. 23, No. 1, pp.43-48, 2016.
- [10] F. Miyazaki, Y. Inagawa, K. Kato, M. Sakaki, H. Ichikawa, and H. Okubo, "Electrode conditioning characteristics in vacuum under impulse voltage application in nonuniform electric field", IEEE Trans. Dielectr. Electr. Insul., Vol. 12, No. 1, pp. 17-23, 2005.
- [11] T. Shioiri, N. Asari, J. Sato, K. Sasage, K. Yokokura, M. Homma, and K. Suzuki, "Investigation of dielectric breakdown characteristics of double-break vacuum interrupter and dielectric breakdown probability distribution in vacuum interrupter", Elect. Eng. Jpn., Vol. 175, No. 2, pp. 13-20, 2011 [IEEJ Trans. Power and Energy, Vol. 128, No. 12, pp. 1452-1458, 2008].



Naoki Hayakawa (M'90) was born on 9 September 1962. He received a Ph.D. in 1991 in electrical engineering from Nagoya University, Japan. Since 1990, he has been at Nagoya University and presently he is a Professor at Nagoya University in the Department of Electrical Engineering and Computer Science. From 2001 to 2002, he was a guest scientist at the Forschungszentrum Karlsruhe/Germany. Prof.

Hayakawa is a senior member of IEEJ and a member of CIGRE.



Kosuke Hasegawa was born on 21 May 1985. He received an M.E. in electrical engineering from Shizuoka University in 2010. He worked at Japan AE Power Systems Corporation from 2010 to 2012. He is an engineer at MEIDENSHA CORPORATION, Japan. He is a member of IEEJ.



**Masayuki Sakaki** was born on 18 September 1956. He received a B.S. in electrical engineering from Science University of Tokyo in 1979. He worked for 23 years in the Engineering Department Switchgear Factory at MEIDENSHA CORPORATION. From 2002 to 2012, he worked at Japan AE Power Systems Corporation. He is currently an Administration Officer of the R&D group at MEIDENSHA CORPORATION. He is a member of IEEJ.



**Hiroki Kojima** (M'11) was born on 7 December 1975. He received his Ph.D. in 2004 in energy engineering and science from Nagoya University, Japan. He was a Research Fellow with the Japan Society for the Promotion of Science from 2000 to 2003. Since 2004, he has been at Nagoya University and presently he is an Associate Professor at Nagoya University in the Department of Electrical Engineering and Computer Science. Dr. Kojima is a

member of IEEJ.



**Tatsuya Takahashi** was born on 26 April 1991. He received his B.S. in electrical and electronic engineering in 2014 and his M.E. degree in electrical engineering and computer science in 2016 from Nagoya University, Japan. Currently, he works at Mitsubishi Electric Corporation, Japan.