

THE $E_6 \otimes SO(10)$ PREON MODEL BASED ON GLOBAL $SU(18)$ COLOR-FLAVOR

Yuko OKAMOTO and R.E. MARSHAK

Physics Department, Virginia Polytechnic Institute and State University, Blacksburg, VA 24061, USA

Received 12 August 1985

We have constructed a version of the chiral three-preon model $E_6 \otimes SO(10)$ based on the global color-flavor symmetry group $SU(18)$. By applying the 't Hooft anomaly matching condition to the subgroup $SU(16) \times SU(2)$ of $SU(18)$ together with a few physical constraints, we obtain a unique solution that gives rise to three generations of the spinorial representation 16 of $SO(10)$ without exotics. Except for $N=18$, no solution at all exists for the global color-flavor group $SU(N)$ ($16 < N < 22$) when $SU(N)$ breaks to $SU(16) \times SU(N-16)$.

The $E_6 \otimes SO(10)$ preon model was first derived [1]^{*1} as the unique chiral three-fermion preon model on the basis of some fairly general conditions. One starts with $G_{MC} \otimes G_{CF}$, where G_{MC} is a simple metacolor group and $G_{CF} = U(N)$ is the largest global color-flavor group, part of which must be gauged. It can then be shown that $G_{MC} = E_6$ and that the gauged color-flavor subgroup is $SO(10)$ (with $N=16$) if one insists on single irreducible representations of both the metacolor and color-flavor groups as well as asymptotic freedom in the metacolor sector. In the GTM $E_6 \otimes SO(10)$ model, three families are predicted if $SO(10)$ descends through the phenomenologically interesting $SU(4)_C \times SU(2)_L \times SU(2)_R$ group [2]. The chief drawback of this $E_6 \otimes SO(10)$ model is the predicted superabundance of exotic fermions and the consequent loss of asymptotic freedom in the composite color-flavor sector.

Recently, Silveira and Zee have proposed [3] a new version of the $E_6 \otimes SO(10)$ preon model in which they start with $SU(27)$ as the global color-flavor symmetry group. These authors end up with the same gauged $SO(10)$ color-flavor subgroup and, in the process, make constructive use of the 't Hooft anomaly matching condition [4] to eliminate undesirable exotics. However, it seems to us that their willingness to surrender asymptotic freedom in the metacolor sector – through

the use of $SU(27)$ – is too steep a price to pay. Furthermore, we differ from Silveira and Zee and agree with Bars [5] that in the application of the 't Hooft anomaly matching condition no index with absolute value greater than unity should be allowed and that the meta-Pauli principle should be imposed.

Keeping these last three points in mind, we have searched for a version of the $E_6 \otimes SO(10)$ preon model (in the spirit of Silveira and Zee) that retains metacolor asymptotic freedom, that gives a unique solution to the 't Hooft condition with "Bars" indices and that predicts at least three families of ordinary quarks and leptons. We find that there is only one such model and this requires $G_{CF} = SU(18)$ so that the preons belong to the representation (27; 18) of $(E_6; SU(18))$. Since the 't Hooft condition is not satisfied for $SU(18)$ when the meta-Pauli principle is imposed, this group must spontaneously break down to a subgroup H_{CF} which, in turn, contains gauged $SO(10)$ as its subgroup. We take H_{CF} to be $SU(16) \times SU(2)$; the $U(1)$ symmetry that would appear in the breakdown $G_{CF} \rightarrow H_{CF}$ is assumed to be dynamically broken by metacolor forces as in ref. [3]. We note that the Appelquist–Carazzone decoupling conditions [4] are not imposed here, since we cannot give E_6 invariant mass term to the preons.

The particle content of the model is summarized in table 1. The 't Hooft anomaly matching condition comes from the three $SU(16)$ currents, (the $SU(2)$ currents give no contribution) and we have:

^{*1} We shall refer to this paper as GTM.

Table 1

		E_6	SU(16)	SU(2)	Indices
preons	P_1	27	\square	1	
	P_2	27	1	\square	
composites	$P_1 P_1 P_1$	1	$\square\square\square, \begin{smallmatrix} \square \\ \square \end{smallmatrix}, \begin{smallmatrix} \square & \square \\ \square & \square \end{smallmatrix}$	1	l_1^+, l_1^-, l_1^0
	$P_1 P_1 P_2$	1	$\square\square, \begin{smallmatrix} \square \\ \square \end{smallmatrix}$	\square	l_2^+, l_2^-
	$P_1 P_2 P_2$	1	\square	$\square\square, \begin{smallmatrix} \square \\ \square \end{smallmatrix}$	l_3^+, l_3^-
	$P_2 P_2 P_2$	1	1	$\square\square\square, \begin{smallmatrix} \square & \square \\ \square & \square \end{smallmatrix}$	l_4^+, l_4^0

$$209l_1^+ + 65l_1^- + 247l_1^0 + 40l_2^+ + 24l_2^- + 3l_3^+ + l_3^- = 27. \tag{1}$$

By the meta-Pauli principle^{#2}, eq. (1) reduces to

$$247l_1^0 + 40l_2^+ + 24l_2^- + 3l_3^+ + l_3^- = 27. \tag{2}$$

Finally, with the Bars condition on the indices, namely $|l_i| \leq 1$, we obtain the *unique* solution

$$l_3^+ = l_2^- = 1, \tag{3}$$

with zero value for all other indices.

The above solution predicts two massless composite fermions, $(\square, \square\square)$ and $(\begin{smallmatrix} \square \\ \square \end{smallmatrix}, \square)$, under $(SU(16), SU(2))$, which correspond to the $SO(10)$ representations $3(16)$ and $2(120)$. The representation 120, however, is real, and fermions corresponding to 120 acquire large masses at the grand unification scale [6]. Hence, we are left with exactly three generations of the spinorial representation 16 of $SO(10)$ without any exotics.

One objection to any version of the $E_6 \otimes SO(10)$ preon model would be that $SO(10)$ is not asymptotically free in the preon color-flavor sector [7]. However, since we have been implicitly assuming that the scale of metacolor confinement Λ_{MC} is larger than Λ_{GUT} , the $SO(10)$ coupling constant would only become large near or even above the Planck mass scale (using the usual renormalization group analysis). Because of the absence of exotic fermions in our version of the

$E_6 \otimes SO(10)$ model, we have no problem with asymptotic freedom in the composite color-flavor sector.

In this note we have shown that in an $E_6 \otimes SO(10)$ preon model the 't Hooft anomaly matching scheme with a few physical constraints suggests the spontaneous breakdown of the global symmetry $G_{CF} = SU(18)$ into $SU(16) \times SU(2)$. The solution to the anomaly consistency equation uniquely predicts three generations of $SO(10)$ GUT without exotics. In fact, if we start with $G_{CF} = SU(N)$ ($16 < N < 22$), the subgroup, $SU(16) \times SU(N-16)$ does not yield any solution at all to the 't Hooft anomaly matching equation (except for $N = 18$). Details will be given elsewhere [8].

We are grateful to Gary Staebler for useful discussions. This work was partially supported by DOE Contract No. DE-AS05-80ER10713.

References

- [1] J.M. Gipson, Y. Tosa and R.E. Marshak, Phys. Rev. D32 (1985) 284.
- [2] See, R.E. Marshak and R.N. Mohapatra, Whither grand unification, in: Festschrift in honor of Y. Ne'eman (Cambridge U.P., London, 1985), to be published.
- [3] V. Silveira and A. Zee, Phys. Lett. 157B (1985) 191.
- [4] G. 't Hooft, in: Recent development in gauge theories, NATO Advanced Study Institute (Cargèse, 1979), eds. G. 't Hooft et al. (Plenum, New York, 1980).
- [5] I. Bars, Nucl. Phys. B208 (1982) 77.
- [6] H. Georgi, Nucl. Phys. B156 (1979) 126.
- [7] See X.Y. Li and R.E. Marshak, Towards a realistic composite model of quarks and leptons, DESY preprint (1985); see also Y. Tosa and R.E. Marshak, Phys. Rev. D27 (1983) 616.
- [8] Y. Okamoto and R.E. Marshak, in preparation.

^{#2} It should be pointed out that we use the meta-Pauli principle in a crucial way to obtain the unique solution to the anomaly matching equation, while the meta-Pauli principle does not play a vital role in the Silveira-Zee model [3].