

## PAPER

## Coding Theoretic Construction of Quantum Ramp Secret Sharing

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**SUMMARY** We show a construction of a quantum ramp secret sharing scheme from a nested pair of linear codes. Necessary and sufficient conditions for qualified sets and forbidden sets are given in terms of combinatorial properties of nested linear codes. An algebraic geometric construction for quantum secret sharing is also given.

**key words:** algebraic geometry code, non-perfect secret sharing, quantum secret sharing, ramp secret sharing

## 1. Introduction

Secret sharing (SS) [1] is a cryptographic scheme to encode a secret to multiple shares being distributed to participants, so that only qualified (or authorized) sets of participants can reconstruct the original secret from their shares. Traditionally both secret and shares were classical information (bits). Several authors [2]–[4] extended the traditional SS to quantum one so that a quantum secret can be encoded to quantum shares.

When we require unqualified sets of participants to have zero information of the secret, the size of each share must be larger than or equal to that of secret. By tolerating partial information leakage to unqualified sets, the size of shares can be smaller than that of secret. Such an SS is called a ramp (or non-perfect) SS [5]–[7]. The quantum ramp SS was proposed by Ogawa et al. [8]. In their construction [8] as well as its improvement [9], the size of shares can be  $L$  times smaller relative to quantum secret than its previous construction [2]–[4], where  $L$  is the number of qudits in quantum secret. We call a quantum state in a  $q$ -dimensional system as a qudit.

In their construction [8], each share is a quantum state on a  $q$ -dimensional complex linear space, and  $q$  has to be larger than or equal to the number  $n$  of participants. When  $n$  is large,  $q$  also has to be large. But it is not clear whether or not such a large dimensional quantum systems are always readily available. To deal with such a situation, we need a quantum ramp SS allowing  $n > q$ . We stress that we study the ramp (non-perfect) SS while [2]–[4] and their subsequent developments [10], [11] studied the perfect SS, and that none of the results in this paper are contained in [2]–[4], [11], [12].

On the other hand, the present paper can be regarded as a generalization of [3], [12]. Because [3], [12] studied connection between perfect quantum SS and the Calderbank-Shor-Steane (CSS) quantum error-correcting codes [13], [14], while our proposed encoding (6) of quantum secret into quantum shares is the same as that of the  $q$ -ary CSS codes. The connection between quantum ramp SS and quantum error correction seems first studied in [10]. Our new contributions that are not given in [10] are (a) necessary and sufficient conditions for qualified sets and forbidden sets that can be easily checked by a digital computer, (b) a quantum procedure for partially reconstructing the quantum secret by an intermediate set of shares, and (c) a construction of quantum ramp SS that allows arbitrarily large  $n$  for a fixed  $q$ . Item (a) completely characterizes the qualified and the forbidden sets. Item (b) above clarifies how much quantum information in the secret can be reconstructed by an intermediate set, which is a share set neither qualified nor forbidden (unauthorized). We note that item (c) above does not contradict with  $q > \sqrt{(n+2)/2}$  [10, Eq. (5)], because [10, Eq. (5)] considered perfect quantum SS.

It is well-known that all linear classical ramp SS can be constructed from a pair of linear codes  $C_2 \subseteq C_1 \subseteq \mathbf{F}_q^n$  [15], [16], where  $\mathbf{F}_q$  is the finite field with  $q$  elements. Smith [4] studied connection between *perfect* linear classical SS and *perfect* quantum SS by using the monotone span program that can express any *perfect* linear classical SS, but he did not consider ramp SS. In this paper we shall show the following.

**Theorem 1:** Let  $J \subseteq \{1, \dots, n\}$  and  $\bar{J} = \{1, \dots, n\} \setminus J$ . For  $\vec{x} = (x_1, \dots, x_n) \in \mathbf{F}_q^n$  define  $P_J(\vec{x}) = (x_i)_{i \in J}$ . We define  $\bar{P}_J$  to be an  $\mathbf{F}_q$ -linear map from  $C_1/C_2$  to  $P_J(C_1)/P_J(C_2)$  sending  $\vec{x} + C_2 \in C_1/C_2$  to  $P_J(\vec{x}) + P_J(C_2) \in P_J(C_1)/P_J(C_2)$ . A quantum ramp SS can be constructed from **any**  $C_2 \subseteq C_1 \subseteq \mathbf{F}_q^n$ , regardless of  $n$  and  $q$ .

1. The constructed quantum SS encodes a quantum secret of  $(\dim C_1 - \dim C_2)$  qudits to  $n$  shares. Each share is a qudit.
2. A set  $J$  of participants can reconstruct

$$\dim \bar{P}_J(\ker(\bar{P}_{\bar{J}})) \quad (1)$$

qudits out of  $(\dim C_1 - \dim C_2)$  qudits of the encoded quantum secret. If

$$\dim \bar{P}_J(\ker(\bar{P}_{\bar{J}})) = \dim C_1 - \dim C_2 \quad (2)$$

then the set  $J$  of participants can reconstruct the secret

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perfectly. This means that  $J$  is a qualified set. In this case  $\bar{J}$  has no information of the secret, which means that  $\bar{J}$  is a forbidden (also called unauthorized) set.

3. The condition (2) is equivalent to both

$$\dim P_J(C_1) - \dim P_J(C_2) = \dim C_1 - \dim C_2 \quad (3)$$

and

$$\dim P_{\bar{J}}(C_1) - \dim P_{\bar{J}}(C_2) = 0. \quad (4)$$

Condition (4) is equivalent to

$$\dim C_2^\perp \cap \ker(P_J) - \dim C_1^\perp \cap \ker(P_J) = 0, \quad (5)$$

where  $C_1^\perp$  denotes the dual code of linear code  $C_1$  with respect to the standard inner product in  $\mathbf{F}_q^n$ .

4. Both (3) and (4) are also a necessary condition for  $J$  to be a qualified set.

We shall explain how the above theorem is useful with a concrete application as below. It was shown that all linear classical secret sharing can be expressed by  $C_2 \subset C_1 \subset \mathbf{F}_q^n$  [17]. The above theorem reveals a clear connection between the access structures (families of qualified and forbidden sets) of classical and quantum secret sharing arising from the same pair  $C_2 \subset C_1$ . By using that connection, we can translate a result for classical secret sharing into quantum one. For example, Iwamoto et al. [18] proposed a construction method of classical linear secret sharing that minimizes share sizes with an arbitrarily given access structure. By using the above theorem, the construction method in [18] was easily translated for quantum secret sharing that minimizes share sizes [19]. We note that before [19] even construction of quantum secret sharing for arbitrary access structures is not well-studied. We also note that for purely classical SS, a similar result to this paper was already reported in [22].

This paper is organized as follows: Section 2 proposes the encoding of secrets and shows Item 1 in Theorem 1. Section 3 proposes the decoding of secrets and it shows Items 2 and 3 in Theorem 1. Section 4 proves Item 4 in Theorem 1 by computing the Holevo information of the set  $J$ . It also computes the coherent information as a byproduct. Section 5 shows that Theorem 1 completely characterizes the qualified and forbidden sets of the quantum ramp SS by Ogawa et al. [8]. Section 6 gives an algebraic geometric (AG) construction. A major benefit of the AG construction is that  $n$  can become arbitrarily large for a fixed  $q$  [20]. Section 7 gives concluding discussions.

## 2. Encoding Secrets

We shall propose a construction of a quantum ramp SS from a nested pair of linear codes  $C_2 \subseteq C_1 \subseteq \mathbf{F}_q^n$ . Our proposal is a quantum version of classical ramp SS proposed by Chen et al. [15, Section 4.2]. Let  $\mathcal{G}_i$  and  $\mathcal{H}_j$  be  $q$ -dimensional complex linear spaces. We also assume that orthonormal bases of  $\mathcal{G}_i$  and  $\mathcal{H}_j$  are indexed by  $\mathbf{F}_q$  as  $\{|s\rangle\}_{s \in \mathbf{F}_q}$ . The quantum secret is  $\dim C_1 - \dim C_2$  qudits on  $\bigotimes_{i=1}^{\dim C_1 - \dim C_2} \mathcal{G}_i$ . Fix an  $\mathbf{F}_q$ -linear isomorphism  $f : \mathbf{F}_q^{\dim C_1 - \dim C_2} \rightarrow C_1/C_2$ .

Also,  $\{|s\rangle \mid s \in \mathbf{F}_q^{\dim C_1 - \dim C_2}\}$  is an orthonormal basis of  $\bigotimes_{i=1}^{\dim C_1 - \dim C_2} \mathcal{G}_i$ . We shall encode a quantum secret to  $n$  qudits in  $\bigotimes_{j=1}^n \mathcal{H}_j$  by a complex linear isometric embedding. To specify such an embedding, it is enough to specify the image of each basis state  $|s\rangle \in \bigotimes_{i=1}^{\dim C_1 - \dim C_2} \mathcal{G}_i$ . We encode  $|s\rangle$  to

$$\frac{1}{\sqrt{|C_2|}} \sum_{\vec{x} \in f(\vec{s})} |\vec{x}\rangle \in \bigotimes_{j=1}^n \mathcal{H}_j. \quad (6)$$

We note that the proposed encoding (6) is equivalent to that of CSS codes [13], [14]. When  $\dim C_1 - \dim C_2 = 1$ , Eq. (6) is a special case of encoding considered in [10, Section 1]. Marin and Markham [10] mostly studied the case in which the sizes of quantum shares are the same as that of quantum secret, while in our study the size of quantum secret is generally larger than those of quantum shares. Recall that by definition of  $f$ ,  $f(\vec{s})$  is a subset of  $C_1$ ,  $f(\vec{s}) \cap f(\vec{s}_1) = \emptyset$  if  $\vec{s} \neq \vec{s}_1$ , and  $f(\vec{s})$  contains  $|C_2|$  vectors. From these properties we see that (6) defines a complex linear isometric embedding. The quantum system  $\mathcal{H}_j$  is distributed to the  $j$ -th participant.

**Example 2:** We show a slightly modified variant of Ogawa et al. [8] as an example. Let  $q = 7$ ,  $n = 5$ ,  $L = 3$ ,  $\alpha_1 = 3$ ,  $\alpha_2 = 5$ ,  $\alpha_3 = 6$ ,  $\alpha_4 = 1$ ,  $\alpha_5 = 4$ . For  $s_1, s_2, s_3 \in \mathbf{F}_7$ ,  $|s_1 s_2 s_3\rangle$  is encoded to

$$\frac{1}{\sqrt{7}} \sum_{r \in \mathbf{F}_7} \bigotimes_{j=1}^5 |r + s_1 \alpha_j + s_2 \alpha_j^2 + s_3 \alpha_j^3\rangle. \quad (7)$$

This encoding can be described by

$$\begin{aligned} C_1 &= \{(r + s_1 \alpha_j + s_2 \alpha_j^2 + s_3 \alpha_j^3)_{j=1, \dots, 5} \mid \\ &\quad r, s_1, s_2, s_3 \in \mathbf{F}_7\}, \\ C_2 &= \{(r, r, r, r, r) \mid r \in \mathbf{F}_7\}, \\ f(s_1, s_2, s_3) &= \{(r + s_1 \alpha_j + s_2 \alpha_j^2 + s_3 \alpha_j^3)_{j=1, \dots, 5} \mid r \in \mathbf{F}_7\}. \end{aligned}$$

## 3. Decoding Secrets

### 3.1 Preliminary Algebra

In this subsection we show Item 3 in Theorem 1 in order to introduce the proposed decoding procedure. The equivalence between (4) and (5) follows from Forney's second duality lemma [21, Lemma 7] and  $\ker(P_J) = \{(x_1, \dots, x_n) \in \mathbf{F}_q^n \mid x_i = 0 \text{ if } i \in J\}$ .

Equation (3) is equivalent to  $\tilde{P}_J$  being an isomorphism, and (4) is equivalent to  $\tilde{P}_{\bar{J}}$  being the zero map. From these observations we see that (3) and (4) imply (2) and vice versa. This finishes the proof of Item 3 in Theorem 1.

**Remark 3:** Equation (5) corresponds to [22, Eq. (3)] for classical ramp SS.

### 3.2 Proposed Decoding Procedure

Suppose that the quantum secret is

$$\sum_{\vec{s} \in \mathbf{F}_q^{\dim C_1 - \dim C_2}} \alpha(\vec{s}) |\vec{s}\rangle \in \bigotimes_{i=1}^{\dim C_1 - \dim C_2} \mathcal{G}_i. \quad (8)$$

It is encoded to  $n$  qudits as

$$\sum_{\vec{s} \in \mathbf{F}_q^{\dim C_1 - \dim C_2}} \alpha(\vec{s}) \frac{1}{\sqrt{|C_2|}} \sum_{\vec{x} \in f(\vec{s})} |\vec{x}\rangle \in \bigotimes_{j=1}^n \mathcal{H}_j. \quad (9)$$

Decompose  $\ker(\tilde{P}_J)$  to a direct sum  $V \oplus (\ker(\tilde{P}_J) \cap \ker(\tilde{P}_J))$ , and decompose  $C_1/C_2$  to  $W \oplus V \oplus \ker(\tilde{P}_J)$ . Let  $\mathcal{G}(J)$  to be the complex linear space spanned by  $\{|\vec{s}\rangle \mid f(\vec{s}) \in V\}$ . We have  $\dim \mathcal{G}(J) = |\tilde{P}_J(\ker(\tilde{P}_J))|$  because

$$\begin{aligned} & \dim \tilde{P}_J(\ker(\tilde{P}_J)) \\ &= \dim \ker(\tilde{P}_J) - \dim \ker(\tilde{P}_J) \cap \ker(\tilde{P}_J) \\ &= \dim V. \end{aligned} \quad (10)$$

The space  $\bigotimes_{i=1}^{\dim C_1 - \dim C_2} \mathcal{G}_i$  can be decomposed as  $\mathcal{G}(J) \otimes \mathcal{G}_{\text{rest}}$ , where  $\mathcal{G}_{\text{rest}}$  is the complex linear space spanned by  $\{|\vec{s}_{KW}\rangle \mid f(\vec{s}_{KW}) \in W \oplus \ker(\tilde{P}_J)\}$ , and  $|\vec{s}_J\rangle \otimes |\vec{s}_W + \vec{s}_K\rangle \in \mathcal{G}(J) \otimes \mathcal{G}_{\text{rest}}$  is identified with  $|\vec{s}\rangle \in \bigotimes_{i=1}^{\dim C_1 - \dim C_2} \mathcal{G}_i$  for  $\vec{s} = \vec{s}_J + \vec{s}_W + \vec{s}_K$  with  $\vec{s}_J \in f^{-1}(V)$ ,  $\vec{s}_W \in f^{-1}(W)$  and  $\vec{s}_K \in f^{-1}(\ker(\tilde{P}_J))$ . This identification is a unitary map between  $\mathcal{G}(J) \otimes \mathcal{G}_{\text{rest}}$  and  $\bigotimes_{i=1}^{\dim C_1 - \dim C_2} \mathcal{G}_i$ , because it is linear and preserves the inner product.

**Example 4:** We retain the notations from Example 2. Let  $J = \{1, 2, 3\}$  and  $\bar{J} = \{4, 5\}$ . Firstly we examine  $\ker(\tilde{P}_{\bar{J}}) \subset C_1/C_2$ . When  $(s_1, s_2, s_3) = (2, 1, 0)$  or  $(s_1, s_2, s_3) = (0, 0, 1)$ ,  $P_{\bar{J}}(f(s_1, s_2, s_3)) = P_{\bar{J}}(C_2)$ , from which we see that  $\ker(\tilde{P}_{\bar{J}})$  is two-dimensional linear space spanned by  $f(2, 1, 0)$  and  $f(0, 0, 1)$ . On the other hand,  $P_J(f(2, 1, 0)) \neq P_J(C_2)$  and  $P_J(f(0, 0, 1)) = P_J(C_2)$ , which mean that  $\ker(\tilde{P}_{\bar{J}}) \cap \ker(\tilde{P}_J)$  is one-dimensional linear space spanned by  $f(0, 0, 1)$ . We also observe that  $V$  is the one-dimensional space spanned by  $f(2, 1, 0)$ , that  $\ker(\tilde{P}_J)$  is the one-dimensional space spanned by  $f(0, 0, 1)$ . There is some freedom in choosing  $W$ , for example, we can choose  $W$  as the one-dimensional space spanned by  $f(1, 0, 0)$ .

$\mathcal{G}(J)$  is the 7-dimensional complex linear space spanned by  $\{|2a\rangle \otimes |a\rangle \otimes |0\rangle \mid a \in \mathbf{F}_7\}$ , while  $\mathcal{G}_{\text{rest}}$  is the 49-dimensional complex linear space spanned by  $\{|s_1\rangle \otimes |0\rangle \otimes |s_3\rangle \mid s_1, s_3 \in \mathbf{F}_7\}$ .

In this section we shall prove that a set  $J$  of participants can reconstruct the part of the quantum secret (8) from (9). The reconstructed part is a state in  $\mathcal{G}(J)$ . By reordering indices we may assume  $J = \{1, \dots, |J|\}$ . We also assume

$$\dim \tilde{P}_J(\ker(\tilde{P}_{\bar{J}})) > 0, \quad (11)$$

otherwise the set  $J$  can reconstruct no part of the secret by the proposed decoding procedure.

The restriction of  $\tilde{P}_J \circ f$  to  $V$  is injective by the definition of  $V$ . This and the definitions of  $V$  and  $W$  imply that there exists an  $\mathbf{F}_q$ -linear isomorphism  $g_1$  from  $P_J(C_1)/P_J(C_2)$  to  $\mathbf{F}_q^{\dim P_J(C_1) - \dim P_J(C_2)}$  with the following

condition. When we write  $\vec{s} = \vec{s}_J + \vec{s}_W + \vec{s}_K$  in the same way as the previous paragraph for  $\vec{s} \in \mathbf{F}_q^{\dim C_1 - \dim C_2}$  then  $g_1(\tilde{P}_J(f(\vec{s}))) = (\vec{s}_J, \vec{s}_W) \in \mathbf{F}_q^{\dim P_J(C_1) - \dim P_J(C_2)}$ . If (2) holds then we have  $V = C_1/C_2$  and we regard  $\vec{s}_W$  and  $\vec{s}_K$  as  $\vec{0}$  and  $\vec{s}_J$  as  $\vec{s}$ . Observe that  $g_1$  is inverting the restriction of  $\tilde{P}_J \circ f$  to  $V$ .

On the other hand, there also exists an  $\mathbf{F}_q$ -linear epimorphism  $g_2$  from  $P_J(C_1)$  to  $\mathbf{F}_q^{\dim P_J(C_2) \cap \ker(P_{\bar{J}})}$  that is one-to-one on every coset belonging to the factor linear space  $P_J(C_1)/P_J(C_2 \cap \ker(P_{\bar{J}}))$ . The above map can be constructed as follows: Find a direct sum decomposition of  $P_J(C_1) = P_J(C_2 \cap \ker(P_{\bar{J}})) \oplus U$ . For  $\vec{x} \in P_J(C_1)$ , find a decomposition  $\vec{x} = \vec{x}_1 + \vec{x}_2$  such that  $\vec{x}_1 \in P_J(C_2 \cap \ker(P_{\bar{J}}))$  and  $\vec{x}_2 \in U$ . Then map  $\vec{x}_1$  by a some fixed linear isomorphism from  $P_J(C_2 \cap \ker(P_{\bar{J}}))$  to  $\mathbf{F}_q^{\dim P_J(C_2) \cap \ker(P_{\bar{J}})}$ , while ignoring  $\vec{x}_2$ . Observe that  $g_2$  is extracting the  $P_J(C_2 \cap \ker(P_{\bar{J}}))$ -component.

By a construction similar to  $g_2$ , there also exists an  $\mathbf{F}_q$ -linear epimorphism  $g_3$  from  $P_J(C_1)/P_J(C_2 \cap \ker(P_{\bar{J}}))$  to  $\mathbf{F}_q^{\dim P_J(C_2) - \dim P_J(C_2 \cap \ker(P_{\bar{J}}))}$  that is one-to-one on every coset belonging to the factor linear space  $P_J(C_1)/P_J(C_2)$  such that the value of  $g_3$  is determined by  $\vec{s}_W$ ,  $\vec{s}_K$ , and  $P_{\bar{J}}(\vec{x})$  independently of  $\vec{s}_J$ . Observe also that  $g_3$  is extracting the  $P_J(C_2)$ -component from the factor linear space  $P_J(C_1)/P_J(C_2 \cap \ker(P_{\bar{J}}))$ .

Consider the  $\mathbf{F}_q$ -linear map  $g_4$  from  $P_J(C_1)$  to  $\mathbf{F}_q^{\dim P_J(C_1)}$  sending  $\vec{v} \in P_J(C_1)$  to  $(g_1(\vec{v} + P_J(C_2)), g_2(\vec{v}), g_3(\vec{v} + P_J(C_2 \cap \ker(P_{\bar{J}}))))$ . We see that  $g_4$  is an  $\mathbf{F}_q$ -linear isomorphism because it is surjective and the domain and the image of  $g_4$  have the same dimension.

For  $\vec{v} \in P_J(C_1)$ , we can construct a unitary operation sending  $|\vec{v}\rangle \in \bigotimes_{j=1}^{|J|} \mathcal{H}_j$  to  $|g_4(\vec{v}), \vec{0}\rangle \in \bigotimes_{j=1}^{|J|} \mathcal{H}_j$ , where  $\vec{0}$  is the zero vector of length  $|J| - \dim P_J(C_1)$ . Since this unitary operation does not change  $\mathcal{H}_{|J|+1}, \dots, \mathcal{H}_n$ , it can be executed only by the first to the  $|J|$ -th participants. Applying the unitary operation to (9) gives

$$\sum_{\vec{s} \in \mathbf{F}_q^{\dim C_1 - \dim C_2}} \alpha(\vec{s}) \frac{1}{\sqrt{|C_2|}} \sum_{\vec{x} \in f(\vec{s})} |\vec{s}_J, \vec{s}_W, g_2(P_J(\vec{x})), g_3(P_J(\vec{x}) + P_J(C_2 \cap \ker(P_{\bar{J}}))), \vec{0}, P_{\bar{J}}(\vec{x})\rangle. \quad (12)$$

$g_2(P_J(\vec{x}))$  can become any vector in  $\mathbf{F}_q^{\dim P_J(C_2) \cap \ker(P_{\bar{J}})}$  independently of  $\vec{s}_J$ ,  $\vec{s}_W$ ,  $\vec{s}_K$  and  $P_{\bar{J}}(\vec{x})$ . Hereafter we denote  $g_2(P_J(\vec{x}))$  by  $\vec{u}_1$ . For a fixed  $\vec{s} \in \mathbf{F}_q^{\dim C_1 - \dim C_2}$   $P_{\bar{J}}(\vec{x})$  can become any vector in the coset  $\tilde{P}_{\bar{J}}(f(\vec{s})) \in P_{\bar{J}}(C_1)/P_{\bar{J}}(C_2)$ , and  $\vec{s}_W$  determines which coset of  $P_{\bar{J}}(C_1)/P_{\bar{J}}(C_2)$  contains  $P_{\bar{J}}(\vec{x})$  independently of  $\vec{s}_J$ ,  $\vec{s}_K$  and  $\vec{u}_1$ . Hereafter we denote the coset  $\tilde{P}_{\bar{J}}(f(\vec{s})) = P_{\bar{J}}(\vec{x}) + P_{\bar{J}}(C_2)$  by  $g_5(\vec{s}_W)$ . By the definition of  $g_3$ ,  $g_3(P_J(\vec{x}) + P_J(C_2 \cap \ker(P_{\bar{J}})))$  is determined by only  $\vec{s}_W$ ,  $\vec{s}_K$  and  $P_{\bar{J}}(\vec{x})$ , that is, independent of  $\vec{s}_J$ . Hereafter we denote  $g_3(P_J(\vec{x}) + P_J(C_2 \cap \ker(P_{\bar{J}})))$  by  $g_6(\vec{s}_W, \vec{s}_K, P_{\bar{J}}(\vec{x}))$ . By using these notations we can rewrite (12) as

$$\sum_{\vec{s} \in \mathbb{F}_q^{\dim C_1 - \dim C_2}} \alpha(\vec{s}) |\vec{s}_J\rangle \frac{1}{\sqrt{|C_2|}} \sum_{\substack{\vec{u}_1 \in \mathbb{F}_q^{\dim P_J(C_2 \cap \ker(P_{\bar{J}}))} \\ \vec{u}_2 \in g_5(\vec{s}_W)}} \alpha(\vec{s}_W, \vec{s}_K) |\vec{s}_W, \vec{s}_K\rangle, \tag{13}$$

which means that the part  $|\vec{s}_J\rangle$  of the quantum secret (8) is reconstructed but in general entangled with the rest of quantum system.

If the quantum secret is a product state written as

$$\sum_{\vec{s} \in \mathbb{F}_q^{\dim C_1 - \dim C_2}} \alpha(\vec{s}) |\vec{s}\rangle = \left( \sum_{\vec{s}_J \in V} \alpha(\vec{s}_J) |\vec{s}_J\rangle \right) \otimes \left( \sum_{\vec{s}_W, \vec{s}_K} \alpha(\vec{s}_W, \vec{s}_K) |\vec{s}_W, \vec{s}_K\rangle \right)$$

then (13) can be written as

$$\left( \sum_{\vec{s}_J \in V} \alpha(\vec{s}_J) |\vec{s}_J\rangle \right) \otimes \left( \sum_{\vec{s}_W, \vec{s}_K} \alpha(\vec{s}_W, \vec{s}_K) \frac{1}{\sqrt{|C_2|}} \sum_{\substack{\vec{u}_1 \in \mathbb{F}_q^{\dim P_J(C_2 \cap \ker(P_{\bar{J}}))} \\ \vec{u}_2 \in g_5(\vec{s}_W)}} |\vec{s}_W, \vec{u}_1, g_6(\vec{s}_W, \vec{s}_K, \vec{u}_2), \vec{0}, \vec{u}_2\rangle \right),$$

and the reconstructed secret is not entangled with the rest of quantum system.

Observe also that the number of qudits in the reconstructed part is  $\dim V = \dim P_J(\ker(P_{\bar{J}}))$  and if (2) holds then the entire secret is reconstructed. Because the complement of any qualified set is forbidden by [8, Proposition 3], we see that the set  $\bar{J}$  of participants has no information on the quantum secret (8) if (2) holds. This finishes the proof of Item 2 in Theorem 1.  $\square$

**Example 5:** We retain the notations from Example 4. We have  $J = \{1, 2, 3\}$ ,  $\dim P_J(C_1) = 3$ , and  $\dim P_J(C_2) = 1$ .  $\dim P_J(C_1)/P_J(C_2) = 2$ .

When we express

$$\vec{s} = \underbrace{a(2, 1, 0)}_{=\vec{s}_J} + \underbrace{s_3(0, 0, 1)}_{=\vec{s}_K} + \underbrace{s_1(1, 0, 0)}_{=\vec{s}_W},$$

and fix  $r$  in (7), the index vector  $\vec{x}$  in (7) becomes

$$\vec{x} = (r + a + 3s_1 + 6s_3, r + 5s_1 + 6s_3, r + 6a + 6s_1 + 6s_3, r + 3a + s_1 + s_3, r + 3a + 4s_1 + s_3).$$

$g_1((x_1, x_2, x_3) + P_J(C_2)) = (3x_2 - x_1 - 2x_3, 2x_2 - x_1 - x_3) = (a, s_1)$ . We have  $C_2 \cap \ker(P_{\bar{J}}) = \{0\}$  and  $g_2$  is the zero map. We have  $g_3(x_1, x_2) = 2x_1 - x_3 = r + 3a + 6s_3$  and  $g_4(x_1, x_2) = (a, s_1, r + 3a + 6s_3)$ . Therefore, after applying the proposed

decoding procedure, the state (7) of encoded shares becomes

$$\frac{1}{\sqrt{7}} \sum_{r \in \mathbb{F}_7} |a, s_1, r + 3a + 6s_3, r + 3a + s_1 + s_3, r + 3a + 4s_1 + s_3\rangle = \frac{1}{\sqrt{7}} \sum_{r' \in \mathbb{F}_7} |a, s_1, r' + 6s_3, r' + s_1 + s_3, r' + 4s_1 + s_3\rangle$$

where  $r' = r + 3a$ .

We see that  $s_1$  determines, independently of both  $a$  and  $s_3$ , the coset  $\{(r' + s_1 + s_3, r' + 4s_1 + s_3) \mid r' \in \mathbb{F}_7\}$ , which is  $g_5(\vec{s}_W)$ .  $P_{\bar{J}}(\vec{x}) = (r' + s_1 + s_3, r' + 4s_1 + s_3)$ ,  $s_1$  and  $s_3$  uniquely determine  $g_3(x_1, x_2, x_3) = r' + 6s_3$  which is  $g_6$ .

#### 4. Holevo Information and Coherent Information of a Set of Shares

##### 4.1 Holevo Information

In this section we prove that both (3) and (4) are necessary for  $J$  to be a qualified set. We use the Holevo information [23] defined as follows. Let  $\mathcal{S}_{\text{in}}$  and  $\mathcal{S}_{\text{out}}$  be sets of density matrices,  $\Gamma$  a completely positive trace-preserving map from  $\mathcal{S}_{\text{in}}$  to  $\mathcal{S}_{\text{out}}$ ,  $\{\rho_1, \dots, \rho_m\} \subset \mathcal{S}_{\text{in}}$ , and  $P$  a probability distribution on  $\{\rho_1, \dots, \rho_m\}$ . The Holevo information is defined as

$$K(P, \{\rho_1, \dots, \rho_m\}, \Gamma) = H\left(\sum_{i=1}^m P(\rho_i) \Gamma(\rho_i)\right) - \sum_{i=1}^m P(\rho_i) H(\Gamma(\rho_i)), \tag{14}$$

where  $H(\cdot)$  denotes the von Neumann entropy counted in  $\log_q$ . The Holevo information essentially expresses the classical information that can be transferred over  $\Gamma$  [23].

Let  $\Gamma_J$  be the completely positive trace-preserving map from  $\mathcal{S}(\bigotimes_{i=1}^{\dim C_1 - \dim C_2} \mathcal{G}_i)$  to  $\mathcal{S}(\bigotimes_{j \in J} \mathcal{H}_j)$  induced by the encoding procedure proposed in Sect. 2, where  $\mathcal{S}(\cdot)$  denotes the set of density matrices on a complex space  $\cdot$ . By  $K_J$  we denote

$$K(\text{uniform distribution}, \{|\vec{s}\rangle\langle\vec{s}| \mid \vec{s} \in \mathbb{F}_q^{\dim C_1 - \dim C_2}\}, \Gamma_J). \tag{15}$$

By [8, Theorem 1] if

$$K_J < \dim C_1 - \dim C_2 \tag{16}$$

then  $J$  is not a qualified set. The encoding procedure in Sect. 2 is a pure state scheme [8, Section 2], that is, the quantum state of all the shares is pure if the encoded quantum secret is pure. By [8, Proposition 3], if  $\bar{J}$  is not a forbidden set, then  $J$  is not a qualified set. By [8, Theorem 1] if

$$K_{\bar{J}} > 0 \tag{17}$$

then  $\bar{J}$  is not a forbidden set.

We shall prove the next proposition. By (3), (4), (16) and (17), Proposition 6 implies that both (3) and (4) are necessary for  $J$  to be a qualified set.

**Proposition 6:**

$$K_J = \dim P_J(C_1) - \dim P_J(C_2). \quad (18)$$

**Proof.**  $\Gamma_J(|\vec{s}\rangle\langle\vec{s}|)$  is the partial trace of (9) over  $\bigotimes_{j \in \bar{J}} \mathcal{H}_j$ . By the definition of partial trace

$$\begin{aligned} & \Gamma_J(|\vec{s}\rangle\langle\vec{s}|) \\ &= \frac{1}{|C_2|} \sum_{\vec{x}_1, \vec{x}_2 \in f(\vec{s})} |P_J(\vec{x}_1)\rangle\langle P_J(\vec{x}_2)| \underbrace{\langle P_{\bar{J}}(\vec{x}_1) | P_{\bar{J}}(\vec{x}_2) \rangle}_{=1 \Leftrightarrow \vec{x}_2 \in \vec{x}_1 + \ker(P_{\bar{J}})} \\ &= \frac{1}{|C_2|} \sum_{\vec{u} \in P_{\bar{J}}(f(\vec{s}))} \sum_{\vec{x}_1 \in f(\vec{s}) \cap P_{\bar{J}}^{-1}(\vec{u})} \sum_{\vec{x}_2 \in f(\vec{s}) \cap P_{\bar{J}}^{-1}(\vec{u})} |P_J(\vec{x}_1)\rangle\langle P_J(\vec{x}_2)| \\ &= \frac{1}{|C_2|} \sum_{\vec{u} \in P_{\bar{J}}(f(\vec{s}))} \left( \sum_{\vec{x}_1 \in f(\vec{s}) \cap P_{\bar{J}}^{-1}(\vec{u})} |P_J(\vec{x}_1)\rangle \right) \\ & \quad \left( \sum_{\vec{x}_2 \in f(\vec{s}) \cap P_{\bar{J}}^{-1}(\vec{u})} \langle P_J(\vec{x}_2)| \right) \\ &= \frac{1}{|C_2|} \sum_{\vec{u} \in P_{\bar{J}}(f(\vec{s}))} \left( \sum_{\vec{x}_1 \in f(\vec{s}) \cap ((\vec{0}, \vec{u}) + \ker(P_{\bar{J}}))} |P_J(\vec{x}_1)\rangle \right) \\ & \quad \left( \sum_{\vec{x}_2 \in f(\vec{s}) \cap ((\vec{0}, \vec{u}) + \ker(P_{\bar{J}}))} \langle P_J(\vec{x}_2)| \right). \quad (19) \end{aligned}$$

For  $\vec{u}_1, \vec{u}_2 \in P_{\bar{J}}(f(\vec{s}))$ , if  $f(\vec{s}) \cap ((\vec{0}, \vec{u}_1) + \ker(P_{\bar{J}})) = f(\vec{s}) \cap ((\vec{0}, \vec{u}_2) + \ker(P_{\bar{J}}))$  then  $\vec{x}_1$  and  $\vec{x}_2$  in (19) are taken over the same set  $P_J(\vec{x}) + P_J(C_2 \cap \ker(P_{\bar{J}}))$ , where  $\vec{x}$  is any vector in  $f(\vec{s}) \cap ((\vec{0}, \vec{u}_1) + \ker(P_{\bar{J}}))$ . Otherwise  $\vec{x}_1$  and  $\vec{x}_2$  in (19) are taken over two disjoint sets in  $P_J(f(\vec{s}))$ . So (19) is equal to

$$\frac{1}{|C_2|} \sum_{A \in P_J(f(\vec{s}))/\sim} \left( \sum_{\vec{v} \in A} |\vec{v}\rangle \right) \left( \sum_{\vec{v} \in A} \langle \vec{v}| \right), \quad (20)$$

where  $\sim$  is the equivalence relation that defines  $\vec{v}_1, \vec{v}_2 \in P_J(\mathbb{F}_q^n)$  to be equivalent if  $\vec{v}_1 \in \vec{v}_2 + P_J(C_2 \cap \ker(P_{\bar{J}}))$ . (20) is an equal mixture of  $|P_J(C_2)/P_J(C_2 \cap \ker(P_{\bar{J}}))|$  projection matrices to non-overlapping orthogonal spaces, therefore its von Neumann entropy is  $\dim P_J(C_2) - \dim P_J(C_2 \cap \ker(P_{\bar{J}}))$ , which is the second term in the right hand side of (14).

By (20), the density matrix of the first term in RHS of (14) is

$$\begin{aligned} & \frac{1}{q^{\dim C_1 - \dim C_2}} \sum_{\vec{s} \in \mathbb{F}_q^{\dim C_1 - \dim C_2}} \frac{1}{|C_2|} \sum_{A \in P_J(f(\vec{s}))/\sim} \\ & \left( \sum_{\vec{v} \in A} |\vec{v}\rangle \right) \left( \sum_{\vec{v} \in A} \langle \vec{v}| \right) \\ &= \frac{1}{|C_1|}, \sum_{A \in P_J(C_1)/P_J(C_2 \cap \ker(P_{\bar{J}}))} \left( \sum_{\vec{v} \in A} |\vec{v}\rangle \right) \left( \sum_{\vec{v} \in A} \langle \vec{v}| \right). \quad (21) \end{aligned}$$

The von Neumann entropy of (21) is

$$\dim P_J(C_1) - \dim P_J(C_2 \cap \ker(P_{\bar{J}})) \quad (22)$$

by the same argument as the last paragraph. By (14)  $K_J = \dim P_J(C_1) - \dim P_J(C_2)$ .  $\square$

## 4.2 Coherent Information

We use the same notation as (14). Denote by  $\Gamma_E$  the channel to the environment so that any pure state is mapped to a pure state by  $\Gamma \otimes \Gamma_E$ . The channel to the environment for  $\Gamma_J$  is  $\Gamma_{\bar{J}}$ . Then the coherent information of the input state  $\rho$  and the channel  $\Gamma$  is defined by [23]

$$H(\Gamma(\rho)) - H(\Gamma_E(\rho)). \quad (23)$$

Equation (23) can become negative. The quantum capacity is expressed by the maximum of the coherent information over  $\rho$  [24].

The coherent information of  $\Gamma_J$  and the completely mixed secret  $\frac{1}{q^{\dim C_1 - \dim C_2}} \sum_{\vec{s} \in \mathbb{F}_q^{\dim C_1 - \dim C_2}} |\vec{s}\rangle\langle\vec{s}|$  is (22) subtracted by (22) with  $J$  substituted by  $\bar{J}$ . Therefore the coherent information is

$$\begin{aligned} & \dim P_J(C_1) - \dim C_2 \cap \ker(P_{\bar{J}}) \\ & - (\dim P_{\bar{J}}(C_1) - \dim C_2 \cap \ker(P_{\bar{J}})). \quad (24) \end{aligned}$$

We consider to maximize (24) by replacing  $C_1$  by  $D$  such that  $C_2 \subset D \subset C_1$ . This amounts to maximize (23) over the quantum state completely mixed over the subspace spanned by  $\{|\vec{s}\rangle \mid f(\vec{s}) \subset D\}$ .

**Lemma 7:** Let  $D$  be as above. Define

$$D' = C_2 + (D \cap \ker(P_{\bar{J}})).$$

Then we have

$$\begin{aligned} & \dim P_J(D) - \dim C_2 \cap \ker(P_{\bar{J}}) \\ & - (\dim P_{\bar{J}}(D) - \dim C_2 \cap \ker(P_{\bar{J}})) \\ &= \dim P_J(D') - \dim C_2 \cap \ker(P_{\bar{J}}) \\ & - (\dim P_{\bar{J}}(D') - \dim C_2 \cap \ker(P_{\bar{J}})). \quad (25) \end{aligned}$$

**Proof.** Let  $D = D' \oplus D''$ . Then  $\dim D'' = \dim P_{\bar{J}}(D'')$  because  $D'' \cap \ker(P_{\bar{J}}) = \{\vec{0}\}$ . Therefore the  $D''$  component in  $D$  does not help to increase the value of (24). Thus  $D'$  yields the same value for (24) as  $D$  and we have (25).  $\square$

So we see that  $D = C_2 + (C_1 \cap \ker(P_{\bar{J}}))$  maximizes the coherent information to its maximum value

$$\begin{aligned} & \dim P_J(C_2 + (C_1 \cap \ker(P_{\bar{J}}))) - \dim C_2 \cap \ker(P_{\bar{J}}) \\ & - (\dim P_{\bar{J}}(C_2 + (C_1 \cap \ker(P_{\bar{J}}))) - \dim C_2 \cap \ker(P_{\bar{J}})) \\ & \quad = \dim P_{\bar{J}}(C_2) \\ &= \dim P_J(C_2 + (C_1 \cap \ker(P_{\bar{J}}))) - \\ & \quad (\dim C_2 \cap \ker(P_{\bar{J}}) + \dim P_{\bar{J}}(C_2) - \dim C_2 \cap \ker(P_{\bar{J}})) \\ & \quad \quad = \dim P_J(C_2) \\ &= \dim \tilde{P}_J(\ker \tilde{P}_{\bar{J}}). \end{aligned}$$

We remark that the proposed decoding procedure in Sect. 3 reconstructs precisely that number of qudits in the secret.

### 5. Analysis of the Conventional Scheme

In this section we show that the conventional quantum ramp secret SS [8] can be regarded as a special case of the proposed construction, and its qualified and forbidden sets can be identified by Theorem 1. Let  $\alpha_1, \dots, \alpha_n$  be pairwise distinct nonzero<sup>†</sup> elements in  $\mathbf{F}_q$ , which correspond to  $x_1, \dots, x_n$  in [8]. Denote  $(\alpha_1, \dots, \alpha_n)$  by  $\vec{\alpha}$ . Let  $\vec{v} \in (\mathbf{F}_q \setminus \{0\})^n$ . Then the generalized Reed-Solomon code  $\text{GRS}_{n,k}(\vec{\alpha}, \vec{v})$  is [25, Section 10.§8]

$$\{(v_1 h(\alpha_1), \dots, v_n h(\alpha_n)) \mid \deg h(x) \leq k - 1\}, \quad (26)$$

where  $h(x)$  is a univariate polynomial over  $\mathbf{F}_q$ . Let  $\vec{1} = (1, \dots, 1) \in \mathbf{F}_q^n$  and  $\vec{\alpha}^L = (\alpha_1^L, \dots, \alpha_n^L) \in \mathbf{F}_q^n$ . The conventional scheme [8] is a special case of the proposed construction with  $C_1 = \text{GRS}_{n,k}(\vec{\alpha}, \vec{1})$  and  $C_2 = \text{GRS}_{n,k-L}(\vec{\alpha}, \vec{\alpha}^L)$ . Observe that  $C_2 \subseteq C_1$ ,  $\dim C_1 = k$ , and  $\dim C_2 = k - L$ . By the property of the generalized Reed-Solomon codes (see e.g. [25, Section 11.§4]), any subset  $J \subseteq \{1, \dots, n\}$  satisfies both (3) and (4) if  $|J| \geq \dim C_1$  and  $|\bar{J}| \leq \dim C_2$ . Observe that the original restriction  $n = \dim C_1 + \dim C_2$  [8] is removed here.

### 6. Algebraic Geometric Construction

In this section we give a construction of  $C_1 \supset C_2$  based on algebraic geometry (AG) codes. A major benefit of the AG codes is that  $n$  can become arbitrarily large for a fixed  $q$  [20]. For terminology and mathematical notions of AG codes, please refer to [20]. Let  $F/\mathbf{F}_q$  be an algebraic function field of one variable over  $\mathbf{F}_q$ ,  $P_1, \dots, P_n$  pairwise distinct places of degree one in  $F$ , and  $G_1, G_2$  divisors of  $F$  whose supports contain none of  $P_1, \dots, P_n$ . We assume  $G_1 \geq G_2$ . Denote by  $\mathcal{L}(G_1)$  the  $\mathbf{F}_q$ -linear space associated with  $G_1$ . The functional AG code associated with  $G_1, P_1, \dots, P_n$  is defined as

$$C(G_1, P_1, \dots, P_n) = \{(f(P_1), \dots, f(P_n)) \mid f \in \mathcal{L}(G_1)\}.$$

Since  $G_1 \geq G_2$  we have  $C(G_1, P_1, \dots, P_n) \supseteq C(G_2, P_1, \dots, P_n)$ . We further assume  $C(G_1, P_1, \dots, P_n) \neq C(G_2, P_1, \dots, P_n)$ .

**Theorem 8:** The ramp quantum SS constructed from  $C(G_1, P_1, \dots, P_n) \supseteq C(G_2, P_1, \dots, P_n)$  encodes  $\dim C(G_1, P_1, \dots, P_n) - \dim C(G_2, P_1, \dots, P_n)$  qudits to  $n$  shares. We have

$$\begin{aligned} & \dim C(G_1, P_1, \dots, P_n) - \dim C(G_2, P_1, \dots, P_n) \\ & \geq \deg G_1 - \deg G_2 - g(F), \end{aligned} \quad (27)$$

where  $g(F)$  denotes the genus of  $F$ . A set  $J \subseteq \{1, \dots, n\}$  is a

<sup>†</sup>In [8]  $\alpha_i = 0$  was not explicitly prohibited, but an author of [8] informed that  $\alpha_i$  must be nonzero for all  $i = 1, \dots, n$ .

qualified set and its complement  $\bar{J}$  is a forbidden set if

$$|J| \geq \max\{1 + \deg G_1, n - (\deg G_2 - 2g(F) + 1)\}. \quad (28)$$

**Proof.** Equation (27) follows just from

$$\begin{aligned} & \dim C(G_1, P_1, \dots, P_n) \\ & = \dim \mathcal{L}(G_1) - \dim \mathcal{L}(G_1 - P_1 - \dots - P_n), \end{aligned} \quad (29)$$

and the Riemann-Roch theorem [20]

$$\deg G_1 - g(F) + 1 \leq \dim \mathcal{L}(G_1) \leq \max\{0, \deg G_1 + 1\}, \quad (30)$$

where the left inequality of (30) becomes equality if

$$\deg G_1 \geq 2g(F) - 1. \quad (31)$$

Firstly we claim that (3) and (4) hold if

$$|J| \geq 1 + \deg G_1, \quad (32)$$

$$|\bar{J}| \leq \deg G_2 - 2g(F) + 1. \quad (33)$$

By reordering indices we may assume that  $J = \{1, \dots, |J|\}$ . Observe that

$$P_J(C(G_1, P_1, \dots, P_n)) = C(G_1, P_1, \dots, P_{|J|}). \quad (34)$$

If (32) holds then by (30) we have  $\mathcal{L}(G_1 - P_1 - \dots - P_{|J|}) = \{0\}$ , which means that  $\mathcal{L}(G_1)$  is isomorphic to  $C(G_1, P_1, \dots, P_{|J|})$  as an  $\mathbf{F}_q$ -linear space by (29). By the same argument we also see that  $\mathcal{L}(G_1)$  is isomorphic to  $C(G_1, P_1, \dots, P_n)$ . Thus we have seen that (32) implies (3).

If (33) holds then

$$\deg(G_2 - P_{|J|+1} - \dots - P_n) \geq 2g(F) - 1,$$

which implies by (31)

$$\dim \mathcal{L}(G_2 - P_{|J|+1} - \dots - P_n) = \deg G_2 - |\bar{J}| - g(F) + 1. \quad (35)$$

By the same argument

$$\dim \mathcal{L}(G_2) = \deg G_2 - g(F) + 1. \quad (36)$$

Equations (29), (35) and (36) imply  $\dim C(G_2, P_{|J|+1}, \dots, P_n) = |\bar{J}|$ , which in turn implies  $C(G_2, P_{|J|+1}, \dots, P_n) = \mathbf{F}_q^{|\bar{J}|}$ . Therefore we see that (33) implies (4).

Finally noting (28)  $\Rightarrow$  (32) and (33) finishes the proof.  $\square$

**Remark 9:** As the generalized Reed-Solomon codes is a special case of AG codes with  $g(F) = 0$  [20], Sect. 5 can also be deduced from Theorem 8 instead of using [25, Section 11.§4].

**Theorem 10:** We retain notations from Theorem 8 and assume  $\deg G_1 < n$ . The number (1) of qudits in quantum secret that can be decoded by  $J$  is

$$\begin{aligned}
& \dim[\mathcal{L}(G_1 - \sum_{j \in \bar{J}} P_j) + \mathcal{L}(G_2)] \\
& - \dim[(\mathcal{L}(G_1 - \sum_{j \in \bar{J}} P_j) + \mathcal{L}(G_2)) \\
& \cap (\mathcal{L}(G_1 - \sum_{j \in J} P_j) + \mathcal{L}(G_2))]. \tag{37}
\end{aligned}$$

**Proof.** Equation (1) is equal to

$$\dim \ker(\tilde{P}_{\bar{J}}) - \dim \ker(\tilde{P}_J) \cap \ker(\tilde{P}_{\bar{J}}). \tag{38}$$

Since we assume  $\deg G_1 < n$ , the evaluation map  $h \in \mathcal{L}(G_1) \mapsto (h(P_1), \dots, h(P_n)) \in \mathbf{F}_q^n$  is injective and we can deal with the space of functions in  $\mathcal{L}(G_1)$  to count the dimensions of (38).

For  $h_1 + \mathcal{L}(G_2) \in \mathcal{L}(G_1)/\mathcal{L}(G_2)$ , its corresponding coset belongs to  $\ker(\tilde{P}_{\bar{J}})$  if and only if there exists  $h_2 \in \mathcal{L}(G_2)$  such that  $h_1(P_j) - h_2(P_j) = 0$  for all  $j \in \bar{J}$ , which is equivalent to  $h_1 - h_2 \in \mathcal{L}(G_1 - \sum_{j \in \bar{J}} P_j)$ . In other words, the coset  $h_1 + \mathcal{L}(G_2)$  satisfies the above condition if and only if there exists  $h'_1 \in \mathcal{L}(G_1 - \sum_{j \in \bar{J}} P_j)$  such that  $h_1 \equiv h'_1 \pmod{\mathcal{L}(G_2)}$ . The dimension of space of cosets  $h_1 + \mathcal{L}(G_2)$  with the above condition is given by

$$\dim[\mathcal{L}(G_1 - \sum_{j \in \bar{J}} P_j) + \mathcal{L}(G_2)] - \dim \mathcal{L}(G_2). \tag{39}$$

Moreover, while satisfying the condition of the last paragraph, the coset corresponding to  $h_1 + \mathcal{L}(G_2)$  belongs to  $\ker(\tilde{P}_J)$  if and only if there exists another  $h''_1 \in \mathcal{L}(G_1 - \sum_{j \in J} P_j)$  such that  $h_1 \equiv h''_1 \pmod{\mathcal{L}(G_2)}$ . The dimension of space of cosets  $h_1 + \mathcal{L}(G_2)$  with the above two conditions is given by

$$\begin{aligned}
& \dim[(\mathcal{L}(G_1 - \sum_{j \in \bar{J}} P_j) + \mathcal{L}(G_2)) \\
& \cap (\mathcal{L}(G_1 - \sum_{j \in J} P_j) + \mathcal{L}(G_2))] \\
& - \dim \mathcal{L}(G_2). \tag{40}
\end{aligned}$$

By (38), subtracting (40) from (39) gives (37).  $\square$

## 7. Conclusion

We have shown that a quantum ramp secret sharing scheme can be constructed from any nested pair of linear codes, and also shown necessary and sufficient conditions for the qualified and the forbidden sets as Theorem 1. A construction of nested linear codes is given by the algebraic geometry in Theorem 8. The following issues are future research agenda.

What is a better construction of  $C_1 \supseteq C_2$  than Theorem 8 when  $q < n$ ? In particular, (33) should use both divisors  $G_1$  and  $G_2$  because (3) and (4) use both of nested linear codes. Also,  $J$  corresponds to a set of  $\mathbf{F}_q$ -rational points on an algebraic curve when AG codes are used, but only the size of  $J$  is taken into account in (33). The geometry of  $J$  should also be taken into account. We shall investigate them

in future.

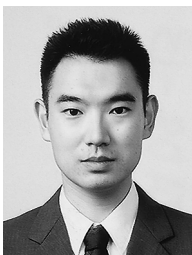
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